M. J. GAREYTE

III

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COMPARISON OF DELTA FUNCTION AND GAUSSIAN WAKES
FOR LONGITUDINAL TRACKING

by D. Brandt

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On the other hand, even for strong deviations of the kurtosis, there seems to be little influence on the results. For these reasons we expect a model approximating the bunch distribution as a function of both its r.m.s. standard deviation and its asymmetry to be sufficient for most longitudinal tracking simulations. Such a model has the advantages of being faster than the delta-function approach and at the same time requiring less space allocation than an expansion of the particle distribution into many orthogonal polynomials.

Fig.1 Bunchlength versus number of turns (Gauss)

Fig.2 Comparison of the skewness (full line = Gauss; dotted line = δ -wake)

Fig.3 Comparison of the kurtosis (full line = Gauss; dotted line = δ -wake)

Fig.4 Bunchlength versus number of turns (δ -wake formalism)

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WIGGLERS FOR CONTROL OF BEAM CHARACTERISTICS IN LEP

by

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WIGGLERS FOR CONTROL OF BEAM CHARACTERISTICS IN LEP

J.M. Jowett and T.M. Taylor
LEP Division
CERN
1211 Geneva 23, Switzerland

Summary

Eight three-block dipole wigglers will be used in LEP to maximize luminosity and control other beam parameters such as emittance, damping rates at injection, energy spread and polarization. The excitation law for constant emittance (chosen, for example, to fill the dynamic aperture) and damping partition numbers is given. In contrast to conventional wigglers where orbit perturbation is avoided by means of a series of equal fields of alternate polarity, the LEP wigglers will use unequal fields and pole lengths. This allows one to retain a high potential asymptotic polarisation level together with an increased polarisation rate. The choice of the degree of field asymmetry is imposed by limitations on the transverse dimensions of the magnets and field levels, by polarisation requirements and by cost. A practical design of the proposed wiggler is presented.

1. Emittance Control

In an electron-positron storage ring with fixed optics, the emittances and energy spread of the beam are uniquely determined by the operating energy. Putting in a dipole wiggler makes it possible to have larger energy spreads and emittances than these "natural" values¹. When the possibility of increasing the coupled horizontal emittance, ϵ_{xc} , by reducing the damping partition number, J_x , is exhausted, a wiggler allows the storage of larger currents before reaching the beam-beam limit (assumed to correspond to a definite value, e.g. $\xi = 0.03$, of the linear tune-shift parameter). Together with the requirement of keeping the occupied region of phase space sufficiently within the dynamic (or physical) aperture, this leads to the prescription of constant $\epsilon_{xc} = \hat{\epsilon}_{xc}$ for maximum stored current at energy $E < \hat{E}$. This is the energy at which the dynamic aperture constraints are just fulfilled with the lowest tolerable value of J_x (0.5, for example) and the wiggler unexcited. In the high-luminosity configurations of machines like LEP, the dynamic aperture is generally smaller than the physical aperture for $E < \hat{E}$. With standard notations:

$$\epsilon_{xc} = \frac{1}{1 + \kappa^2} \frac{55}{32 \sqrt{3}} \left(\frac{\hbar}{mc} \right) \left(\frac{E}{mc^2} \right)^2 \frac{I_5^0 + I_5^w}{J_x (I_2^0 + I_2^w)},$$

where I_i^0 and I_i^w ($i = 1, \dots, 5$) are the contributions of the lattice and the wiggler to the synchrotron integrals, respectively.

Fig. 1 Model used for calculations.

With N_s of the wiggler installation shown schematically in Fig. 1, and setting $r = L_-/L_+$:

$$I_2^w = N_s L_- (1+r) G_-^2, \quad \text{and}$$

$$I_5^w = N_s L_- (1+r^2) \langle W \rangle_w |G_-|^3,$$

where $G_- = G_+/r$ is the inverse of the bending radius in the outer '-' blocks of the wiggler. The function

$$W = \left[D_x^2 + \left(\beta_x D_x' - \frac{1}{2} \beta_x' D_x \right)^2 \right] / \beta_x$$

is changed by the excitation of the wiggler but this can be neglected in the cases of interest.

Neglecting the energy lost by synchrotron radiation in the wiggler in comparison with that in the rest of the machine allows one to solve the equation for constant emittance and gives:

$$B_-(E) \approx \pm \frac{1}{ec} \left[\frac{I_5^0}{N_s (1+r^2) L_- \langle W \rangle_w} \right]^{1/3} \tilde{E}, \quad B_+ = -r B_-,$$

where $\tilde{E}$ is the geometric mean of $\hat{E} - E$, E and $\hat{E} + E$. This excitation provides maximum luminosity with colliding beams and need not be followed during ramping. Injection can also be facilitated by using a high field level to increase damping rates.

The solid curve in Fig. 2 shows how the required wiggler field varies as a function of energy. It is interesting to note that the highest wiggler field occurs at $E = \hat{E}/\sqrt{3}$.

Fig. 2 Typical curves of required wiggler excitation and corresponding polarization

Fig. 4 Computed field over half the length of the magnet.

4. Optical Properties

The trajectory of the beam through the wiggler is such that edge focusing will occur at the passages from - to + and from + to - fields. Whereas the effects are self-compensating in the horizontal plane, in the vertical plane they add, and the wiggler acts as a thin focusing lens. In the worst case, i.e. at injection energy (20 GeV) and with the wiggler fully excited, the focal length of this lens is 4300 m: this gives rise to a slight mismatch of the vertical beta function and a total vertical Q-shift of about 0.01 for the ensemble of eight magnets. It is foreseen to compensate for this either by applying a general correction, or by matching, using nearby quadrupoles. The latter method is facilitated by the fact that, as the wigglers will be installed at the lattice end of the dispersion suppressor, the closest quadrupoles are individually controlled, independent of the lattice strings.

The maximum sagitta in the wiggler, occurring at injection, is estimated to be 4.7 mm. In order to minimize the required horizontal aperture, the wiggler mag-

Fig. 5 End coil vs. centre coil excitation for no net effect on the closed orbit.

nets will be located near to vertically focusing quadrupoles, i.e. where B_x is small. Such a location of the magnet is also favourable for spreading the synchrotron radiation in the vertical plane.

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