Top quark pair production cross-section results at LHC in final states with tau or no leptons

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Abstract. ATLAS and CMS measured the top quark pair production inclusive cross section using proton-proton collision data collected at 7 TeV c.m. energy at the LHC during the year 2011 in several channels. This article presents the measurements of the cross section using selections of top quark pair production events with one hadronically decaying \( \tau \) lepton or no leptons at all.

1. Introduction
Top quark pairs, \( t\bar{t} \), are produced in proton-proton collisions via the strong interaction of colliding partons. At 7 TeV c.m. energy about 80% of \( t\bar{t} \) is produced by gluon-gluon interactions. The remaining 20% is produced in quark and anti-quark annihilation processes. The theoretical cross section of the inclusive \( t\bar{t} \) production within the Standard Model (SM) is calculated at the approximate NNLO level using the HATHOR tool [1], and is \( \sigma^{\text{theory}} = 167^{+17}_{-18} \) pb at \( \sqrt{s} = 7 \) TeV and \( m_{\text{top}} = 172.5 \) GeV.

A top quark decays almost always into a \( W \) boson and \( b \) quark. Therefore, \( t\bar{t} \) events are distinguished by the decay products of the two \( W \) bosons. \( W \) bosons can decay into a quark pair or a lepton-neutrino pair. The ATLAS [2] and CMS [3] experiments measured the \( t\bar{t} \) production cross section in different final states. These final states are affected by different background processes, which makes it necessary to have unique analysis techniques developed for each particular measurement. Precise measurements of the cross section in all final states helps to test the SM against some new physics phenomena, which could influence on the measurements.

Details of cross-section measurements of both ATLAS and CMS in three final states are reviewed in this document: a) \( \tau_{\text{had}} \)+lepton: one \( W \) boson decays into a \( \tau \) lepton and a neutrino and the \( \tau \) lepton decays hadronically. The second \( W \) boson decays leptonically including the decay into \( \tau \) and neutrino, when the former decays leptonically; b) \( \tau_{\text{had}} \)+jets: same as in the previous case but the second \( W \) boson decays hadronically; c) fully hadronic: both \( W \) bosons decay hadronically.

2. Data 2011 and Monte-Carlo (MC) simulation
Both ATLAS and CMS detectors recorded a large fraction of the total integrated luminosity delivered by the LHC machine during the data taking year 2011. The ATLAS and CMS studies, which are presented in this document used up to 4.7 \( \text{fb}^{-1} \) and 3.9 \( \text{fb}^{-1} \) of data, respectively.
The ATLAS signal sample was generated with MC@NLO [4] and showered and hadronised with HERWIG [5]. The single top quark background sample was also produced with these generators. In contrast, CMS used MADGRAPH [6] together with PYTHIA [7] to generate the signal as well as $W/Z+\text{jets}$. The single top quark background was generated with POWHEG [8] together with PYTHIA. PYTHIA was used by CMS also to generate di-bosons and QCD multi-jet samples. The ATLAS studies used di-boson samples generated with HERWIG while the QCD multi-jet and $W/Z+\text{jets}$ background was generated with ALPGEN [9] and HERWIG. Both experiments used TAUOLA [10] to simulate $\tau$ decays. All top quark samples were generated assuming $m_{\text{top}} = 172.5$ GeV.

3. Three cross-section measurements

3.1. $\tau_{\text{had}}+\text{lepton}$ final state

The total branching ratio of all hadronic decays of $\tau$ is about 65%. The branching ratio of the $\tau_{\text{had}}+\text{lepton}$ final state is about 3.7% of the inclusive $t\bar{t}$ production. The background is any process with a $\tau_{\text{had}}+\text{lepton}$ or a $\tau_{\text{had}}+\text{jets}$ final state, where one jet is wrongly identified as $\tau_{\text{had}}$. The main background processes are $t\bar{t}$ with different $W$ boson decay modes and $W/Z+\text{jets}$.

CMS measured the cross section separately in $t\bar{t} \rightarrow \tau_{\text{had}} + e + X$ and $t\bar{t} \rightarrow \tau_{\text{had}} + \mu + X$ final states using data samples of $2.0 \, \text{fb}^{-1}$ and $2.2 \, \text{fb}^{-1}$ respectively [11]. Signal candidate events were selected using the combined lepton (electron or muon) + dijets + $H_T^{\text{miss}}$ trigger, where $H_T^{\text{miss}}$ is the absolute value of the vectorial sum of all jets momenta. The further (offline) selection of events required an isolated electron (muon) with a transverse momentum, $p_T > 35 \, (30) \, \text{GeV}$, and an absolute pseudo-rapidity, $|\eta| < 2.5 \, (2.1)$. A large missing transverse energy, $E_T^{\text{miss}} > 45 \, (40) \, \text{GeV}$, and at least two jets with $p_T > 35 \, (30) \, \text{GeV}$ and $|\eta| < 2.4$ were required. At least one of the jets was required to be $b$-tagged. One isolated $\tau_{\text{had}}$ candidate jet was required with $p_T > 20 \, \text{GeV}$, $|\eta| < 2.4$ and a charge sign opposite to the selected electron (muon). The $\tau_{\text{had}}$ algorithm used neutral and charged hadrons in a jet as inputs in order to discriminate $\tau_{\text{had}}$ from hadronic jets. The background was estimated from MC simulation except for the misidentified $\tau_{\text{had}}+\text{lepton}$ background, which was not well modeled in simulation due to the complexity of a precise description of numerous detector effects. A data-driven method was used for the estimation of the misidentified $\tau_{\text{had}}+\text{lepton}$ background. The method assigns a variable probability to every jet to be misidentified. This probability map, which depends on $p_T$, $\eta$ and a width in $\eta \times \phi$ of jets, where $\phi$ is the azimuthal angle around the beam axis, was obtained from a data sideband region enriched with QCD multi-jet events. Since one knows the total rate of the processes (e.g., $W(\to \ell\nu)+\text{jets}$), which contribute into the misidentified background one estimates this contribution by weighting jets. The cross section was measured using the “cut and count” method: the expected number of background events was subtracted from the number of selected events in data and the result was normalized to the inverse of the product of the signal acceptance (obtained from MC simulation) and the integrated luminosity of data. Systematic uncertainties on the expected number of background events, the signal acceptance and the integrated luminosity of data were considered. The results are $\sigma_{t\bar{t}}(t\bar{t} \rightarrow \tau_{\text{had}} + e + X) = 136 \pm 23^{(\text{stat.})} \pm 23^{(\text{syst.})} \pm 3^{(\text{lumi.})} \, \text{pb}$ and $\sigma_{t\bar{t}}(t\bar{t} \rightarrow \tau_{\text{had}} + \mu + X) = 147 \pm 18^{(\text{stat.})} \pm 22^{(\text{syst.})} \pm 3^{(\text{lumi.})} \, \text{pb}$ [11]. These results were combined using the Best Linear Unbiased Estimator (BLUE) method [12]. The result of the combination is $\sigma_{t\bar{t}} = 143 \pm 14^{(\text{stat.})} \pm 22^{(\text{syst.})} \pm 3^{(\text{lumi.})} \, \text{pb}$ [11].

The ATLAS measurement used $2.05 \, \text{fb}^{-1}$ of data [13]. Single lepton triggers were used for the event selection. The offline selection required a single isolated electron (muon) with $E_T > 25 \, \text{GeV}$ ($p_T > 20 \, \text{GeV}$) and $|\eta| < 2.47 \, (2.5)$ excluding the calorimeter gap regions $1.37 < |\eta| < 1.52$. Presence of a $\tau_{\text{had}}$ candidate with $20 < E_T < 100 \, \text{GeV}$ and $|\eta| < 2.3$ and not overlapping with the selected lepton, $\Delta R(\text{lepton}, \tau_{\text{had}}) > 0.4$, was required. In addition, $E_T^{\text{miss}} > 30 \, \text{GeV}$ and $E_T^{\text{miss}} + \sum |p_T| > 200 \, \text{GeV}$, where the sum is over the lepton and hadronic
jets, requirements were applied. Presence of at least two hadronic jets with $p_T > 25$ GeV and $|\eta| < 2.5$ and not overlapping with the $\tau_{\text{had}}$ candidate, $\Delta R(\text{jet}, \tau_{\text{had}}) > 0.4$, was required.

Selected events were split in eight categories depending on the decay type of the $\tau_{\text{had}}$ candidate: with one charged track ("1 prong") and with three charged tracks ("3 prong"); the sign of the isolated lepton and $\tau_{\text{had}}$ electric charges: opposite-signed (OS) and same-signed (SS); the number of $b$-jets: $\geq 1$ and 0 $b$-jets. A track-based algorithm was used to identify "1 prong" and "3 prong" $\tau_{\text{had}}$ candidates. A Boosted Decision Tree discriminator, BDT$_{\ell}$ was used to reject the background from misidentified electrons. Another BDT$_{j}$ discriminator was used for the further selection of events. The input templates for BDT$_{j}$ were obtained by subtracting the corresponding distributions, OS-SS, that canceled out the symmetric background in these two selections. Data-driven templates of the fake $\tau_{\text{had}}$ background due to jets were obtained in the 0 $b$-jet events and corrected for the kinematics of $\geq 1$ $b$-jets events using weights calculated from MC simulation.

Figure 1 shows the BDT$_{j}$ output for the "1 prong" $\tau_{\text{had}}$+lepton events with $\geq 1$ $b$-jets after the template fit. Using the template fit of the BDT$_{j}$ output for the "1 prong" and "3 prong" $\tau_{\text{had}}$+lepton events with $\geq 1$ $b$-jets the cross section was measured to be $\sigma_{tt}(tt \rightarrow \tau_{\text{had}} + e + X) = 186 \pm 15(\text{stat.}) \pm 20(\text{syst.}) \pm 7(\text{lumi.})$pb and $\sigma_{tt}(tt \rightarrow \tau_{\text{had}} + \mu + X) = 187 \pm 18(\text{stat.}) \pm 20(\text{syst.}) \pm 7(\text{lumi.})$pb [13]. These results were combined using the BLUE method, $\sigma_{tt} = 186 \pm 13(\text{stat.}) \pm 20(\text{syst.}) \pm 7(\text{lumi.})$pb [13].

3.2. $\tau_{\text{had}}$+jets final state

The branching ratio of the $\tau_{\text{had}}$+jets final state is about 9.8%. QCD multi-jet events are one of the main backgrounds due to fake $\tau_{\text{had}}$ leptons. The combinatorial background, where in a signal event the $\tau_{\text{had}}$ candidate is misidentified as a hadronic jet, is important, too.

The ATLAS measurement used 1.67 fb$^{-1}$ of data [14]. The event selection was done with a multi-jet trigger. The offline selection required the presence of at least five jets with $p_T > 20$ GeV, $|\eta| < 2.5$ and at least two of them tagged as $b$-jets. One of the untagged jets had to be identified as $\tau_{\text{had}}$. A veto on isolated leptons was applied. To reduce the QCD multi-jet background large $E_{\text{T}}^{\text{miss}}$ significance, $E_{\text{T}}^{\text{miss}}/(0.5 \times \sqrt{\text{GeV}} \times \sqrt{\sum E_{\text{T}}}) > 8$, was required, where $\sum E_{\text{T}}$ is the sum of transverse energies of all objects contributing to the calculation of $E_{\text{T}}^{\text{miss}}$. The $\tau_{\text{had}}$ candidate jet was selected by excluding other jets using the kinematics and $b$-tagging information. A template fit method was used to measure the signal cross section, exploring the number of tracks, $n_{\text{track}}$, of the $\tau_{\text{had}}$ candidate jet. MC simulation predicted a template for misidentified electrons as $\tau_{\text{had}}$ candidates similar to the signal template. Therefore, the two were joined in a single template. A multi-jet template was derived using a sideband region while a template for the combinatorial background was obtained from MC simulation. In the template fit the number of events of the combinatorial background was constrained to the number of the signal events since they both originate from $tt$ production. Figure 2 shows...
the fit result. Contamination of the fitted signal+electron template by the single top quark and +jets background was subtracted using the predictions from MC simulation. The known (MC simulation) fraction of +jets background was used to estimate the number of τ̄τ̄ events in the fitted signal+electron template. The result is σττ = 200 ± 19(stat.) ± 43(syst. + lumi.)pb [14].

The CMS measurement used 3.9 fb⁻¹ of data [15]. A multi-jet trigger was used for the event selection. A jet identified as a τ̄had candidate with pT > 45 GeV and |η| < 2.3 was required to match to one of the trigger jets. The event selection required the presence of at least three jets with pT > 45 GeV and |η| < 2.4 among which exactly one jet is tagged as b-jet. Additional jets were counted if pT > 20 GeV. A veto on isolated leptons was also applied. Finally, it was required that EmissT > 20 GeV. A Neural-Network discriminator output was built using the MultiLayerPerceptron (MLP) algorithm. The input distributions were derived from MC simulation except for the multi-jet background. For the latter, the same signal selection of events was used but with a 0 b-jet requirement. Correction factors as derived from MC simulation were applied to the distributions in order to take the difference between the kinematics of the 0 and 1 b-jet events into account. A result of the template fit of the MLP output is shown in Figure 3. The final event selection was done by imposing a cut, NNOutput > 0.5, on these distributions after the fit. The measured cross section is σττ = 156 ± 12(stat.) ± 33(syst.) ± 3(lumi.)pb [15].

3.3. Fully hadronic final state

The fully hadronic final state has the largest branching ratio, 46%. A cross-section measurement suffers from the overwhelming QCD multi-jet background. The ATLAS and CMS measurements were based on kinematic fit methods described below.

The CMS measurement used 1.09 fb⁻¹ of data [16]. The event selection was done using a multi-jet trigger. The offline selection of events required at least six jets with |η| < 2.4 among which four leading jets had to fulfill a pT > 60 GeV requirement. The fifth and sixth jets were required to have pT > 50 GeV and pT > 40 GeV, respectively. Additional jets with pT > 30 GeV were also counted. At least two jets were required to be b-tagged. The jets were
used to reconstruct the $t\bar{t}$ system. All possible permutations of the jets were used to reconstruct two $W$ bosons and two top quarks but only the $b$-tagged jets were used as candidates for the true $b$ quark jets. For each permutation a least-square fit was done. Gaussian resolutions as derived from MC simulation were used for the energies of the jets. The true masses of the $W$ bosons were constrained to the known value, $m_W = 80.4$ GeV, while both top quark true masses were the free parameters of the fit, but constrained to be equal to each other. The permutation leading to the smallest $\chi^2$ value was taken and if its probability was greater than 0.01, the event was selected. The fitted mass in the selected events was used to create templates for a likelihood fit with the signal fraction, $f_{\text{sig}}$, as a free parameter. The QCD multi-jet template was derived from a sideband selection of events, which was similar to the signal selection but requiring no $b$-jets. A two dimensional map ($\eta \times p_T$) of weights, aimed at taking into account the difference in kinematics of light jets in events with 0 and $\geq 2$ $b$-jets, was derived using the whole data sample (after trigger selection). Since this sample is dominated by events with only light jets, the impact of the true $b$-jets on the weights is negligible. The likelihood fit result is shown in Figure 4. The fitted value of the signal fraction was converted to a cross section, $\sigma_{t\bar{t}} = 136 \pm 20(\text{stat.}) \pm 40(\text{syst.}) \pm 8(\text{lumi.})\text{pb}$ [16].

![Figure 4](image_url)

**Figure 4.** The result of the likelihood fit of the $t\bar{t}$ signal (solid red curve) and QCD multi-jet background (dashed blue curve) templates of the top quark mass (from the kinematic fit) to data. The uncertainty stated on the signal fraction, $f_{\text{sig}}$, is only statistical [16].

The ATLAS measurement used 4.7 fb$^{-1}$ of data [17]. Events were selected with a multijet trigger. The offline selection required at least six jets with $|\eta| < 2.5$ and among them at least five jets with $p_T > 55$ GeV and a sixth jet with $p_T > 30$ GeV. At least two out of the leading five jets were required to be $b$-tagged. The jets were required to not overlap, $\Delta R(\text{jet}_i, \text{jet}_k) > 0.6$, where $i$ and $k$ are the jet indices but $i \neq k$. A more stringent cut was applied to the two leading $b$-jets, $\Delta R(b_{j1}, b_{j2}) > 1.2$. It was required that the $E_T^{\text{miss}}$ significance be $S_T = E_T^{\text{miss}} / (0.5 \sqrt{\text{GeV}^2} \times \sqrt{H_T}) < 6$, where $H_T$ is the scalar sum of $p_T$ of all jets. A veto on isolated leptons was applied. All possible permutations of at least six and up to ten jets were used to construct a likelihood function for every selected event. The free parameters

![Figure 5](image_url)

**Figure 5.** The result of the unbinned likelihood fit of the signal (black histogram) and background (green filled histogram) templates of the top quark mass, $m_t$, to data, $m_t$ is calculated in the kinematic fit (see text) [17].
of the likelihood function were the top quark true mass and the six quark energies. Breit-Wigner constraints were introduced in the function for the masses of the two reconstructed W bosons with the fixed true mass and width. Breit-Wigner constraints were also used for the reconstructed top quark masses. Both true top quark masses (free parameters) were required to be equal to each other. The true width of top quarks was constrained in the fit to vary according to expectation. The likelihood function included a product of six transfer functions of quark energies to energies of jets as obtained from MC simulation. Also, the likelihood function had factors, which accounted for the $b$-tagging efficiency and the mistagging rate of jets depending on a particular permutation of these jets. These additional factors helped to give higher probability to those permutations, where $b$-tagged jets were used as true $b$-jets. An event was selected if the permutation with a maximum likelihood value, $\mathcal{L}^\text{max}$, fulfilled the requirement, $\mathcal{L}^\text{max} / \sum \mathcal{L} > 0.8$, where the sum is over all permutations, and the fitted mass, $m_t$, was greater than 125 GeV. Using the best permutation of jets and fitted values of free parameters, a $\chi^2$ value was calculated for the reconstructed $t\bar{t}$ system in a given event. The event was finally selected if $\chi^2 < 30$. The distributions of the fitted top quark mass in the selected events were used in a template fit in order to measure the signal fraction in data. The fit result is shown in Figure 5. The QCD multi-jet template was derived using the same signal selection but requiring no $b$-tagged jet. Differences in kinematics of jets in the events with 0 and $\geq 2$ $b$-tagged jets were studied in a MC simulation sample of QCD multi-jet production and taken into account as a systematic uncertainty. The measured signal fraction was converted into a measurement of the cross section, $\sigma_{t\bar{t}} = 168 \pm 12(\text{stat.})^{+60}_{-57}(\text{syst.}) \pm 7(\text{lumi.}) \text{pb}$ [17].

4. Summary
The $t\bar{t}$ production cross section was measured by the ATLAS and CMS experiments in three different final states independently: $\tau_{\text{had}} + \text{lepton}$, $\tau_{\text{had}} + \text{jets}$ and fully hadronic. The obtained results are consistent with the Standard Model prediction for the $t\bar{t}$ production cross section.

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References
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