FURTHER THOUGHTS ON THE TRANSFER FROM BOOSTER TO CPS

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1. Introduction

The beam optics presented in the "beige report" have been calculated by A. Ball and R. Billing.

It was considered more as a plausibility study than as an optimized solution. The present paper, first discussed in December 1967, emphasises a greater simplicity and gives a basis for the optimisation which has been worked out since then.

2. New scheme for the recombination of the 4 beams

The new scheme shown in Fig. 2 meets the advantages of Bovet's first proposal and Asner's variant which has been used in the "beige report".

Comparing Fig. 1 and Fig. 2 one can list the advantages of the new scheme as follows:

- no undue vertical displacement of beam 2 and 3 in the 10 bunch mode,
- suppression of VBM2, VBM3, VSM2, VSM3,
- no physical displacement needed when changing from the 20 bunch mode to the 10 bunch mode.

3. Dispersion matching for the two-turn injection

Figures 3 and 4 show the two-turn injection with $\phi = n + \frac{\pi}{4}$ of respectively an achromatic and a matched beam.
assuming zero loss one sees that the emittance of the resulting beam is the same in both cases with a slightly higher core density in the achromatic case;

such a high core density and even a smaller emittance might be obtained with the matched beam by losing 1-2% of the beam intensity on the septum (see dashed area of Fig. 4, revolution 5).

It has therefore been agreed that the achromatic injection possibility will be discarded.

4. $\beta$-matching

The individual beams or the vertically added beams are to be matched for the injection into CPS (the wanted $\beta$ in ss 42 are: $\beta_V = 22$ m, $\beta_H = 12$ m, for single turn injection, and $\beta_H = 12/\sqrt{3}$ m, for the two-turn injection). This matching is provided by the quadrupoles of the matching section (see Fig. 12).

The mismatch resulting from the unavoidable optical action of VBM 1, 2, 3 and VSM 1, 2, 3 (see Fig. 5) may be expressed as an increase of the emittances

<table>
<thead>
<tr>
<th>VBM 1,2 + VSM 1,2 (rectangular magnets)</th>
<th>$\Delta E_{\beta_H}$</th>
<th>$\Delta E_{\beta_V}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>(wedge magnets)</td>
<td>8%</td>
<td>0</td>
</tr>
<tr>
<td>VBM 3 + VSM 3 (rectangular magnets)</td>
<td>14%</td>
<td>0</td>
</tr>
<tr>
<td>(wedge magnets)</td>
<td>7%</td>
<td>3%</td>
</tr>
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</table>

where the wedge magnets considered are between rectangular and sector magnets (their end shape angle is $\alpha/4$).
A computation showed that a perfect compensation for VBM 3 - VSM 3 action may be achieved with very small variations on $Q_1$, $Q_2$, $Q_1'$ and $Q_2'$. On the other hand a compensation for VBM 1,2 - VSM 1,2 action would require the introduction of small doublets but this does not seem worthwhile.

5. **Horizontal dispersion matching**

Ideally a transfer with dispersion matching is made as shown in Fig. 6a, but practically we have to cope with the following facts:
- an ideal dispersion matching would require all the length available for the transfer,
- recombination and $\beta$-matching also need the total length,
- there is an additional angle of 9° in the middle (see Fig. 6b).

One then has to abandon the spatial separation between $\beta$-matching and dispersion matching.

Let us suppose the dispersion vector $D$ has to be as shown in Fig. 7b. The matching is done in two steps:

i) with the lens $L$ (Fig. 8) the beam is shaped so as to give the right "dispersion invariant" after the bending $B$ (Fig. 7a);

ii) subsequent $\beta$-matching not only provides the proper $\beta$, but also the right betatronic phase $\varphi$ (see Fig. 7b).

For a complete flexibility the $\beta$-matching then needs six quadrupoles. But for the 20 bunch mode (first stage) four quadrupoles might perhaps be sufficient.

There will be no dilution due to the horizontal dispersion.

For the vertical dispersion nothing can be tuned.
6. **Kicker jitter**

The increased emittance is given by

\[ \frac{\delta \varepsilon}{\varepsilon} = \frac{\delta \alpha}{\alpha} = \frac{a}{E} \frac{\delta a}{a} \alpha \]  \hspace{1cm} (1)

where \( \alpha \) is the kicker deflexion, \( a \) and \( b \) are defined in Fig. 9a.

As an illustration of this effect let us consider a beam where \( E = 40 \pi \mu \text{ rad.m}, \ a = 40 \text{ mm}, \ \alpha = 10 \text{ mrad} \) and \( \delta \alpha/\alpha = \pm 1\%: \Delta \varepsilon/E = 10\% \).

Therefore one has to minimize rather carefully the product \( a \cdot \alpha \) for the beam in the kickers.

7. **Vertical addition**

The formulae giving the circumscribing ellipse in the general case of Fig. 9b have been derived.

Describing the incoming ellipse by \( \varepsilon_1, L_1, S_1 \) (see Fig. 9a), the distance between the centres of the beams \( 2d_f \) (see Fig. 9b), the circumscribing ellipse parameters \( \varepsilon_2, L_2, S_2 \) are as follows:

\[
\begin{align*}
\varepsilon_2 &= R_1 R_2^2 \varepsilon_1 \\
L_2 &= \frac{R_1 L_1 (L_1^2 + S_1^2)}{R_1^2 S_1^2 + L_1^2} \\
S_2 &= \frac{R_2 S_1 (L_1^2 + S_1^2)}{R_1^2 S_1^2 + L_1^2}
\end{align*}
\]  \hspace{1cm} (2)

where

\[
R_1 = \sqrt{(3f + \sqrt{f^2 + 8})/(-f + \sqrt{f^2 + 8})}
\]

\[
R_2 = \frac{1}{2\sqrt{4 - f^2 + \sqrt{f^2 + 8}}}
\]
Two curves giving the values of $R_1 \cdot R_2$ and $R_2$ are drawn in Fig. 10.

Then the required angle $\alpha$, between the two beams at the addition point is

$$\alpha = f \cdot d \cdot s_4 / [L_4 (H + \sqrt{H^2 - 1}) - 1],$$

(3)

where

$$H = (L_4^2 + S_4^2 + 1)/2 L_4.$$ 

Other calculations were made under the assumption of a uniform density in the phase plane, which give:

1) the increase in emittance $E/E_0$,
2) the loss of particles on the septum $p$, and
3) the loss of density (after filamentation) $D/D_0 = (1 - p)E_0/E$, as functions of the separation parameter $f$ and the septum thickness parameter $e$.

The curves are shown in Fig. 11.

8. Beam layout

Figure 12 shows the arrangement of the elements. Major changes are:

- dropping the 3,791° bend on the Linac line
- reducing from 4° to 3° the ejection angle
- dropping the 3,052° horizontal bending
- joining the bendings for beam stopper and dispersion matching, presently + 25°, - 10°
- replacing the first focusing triplet by a doublet
- other changes have been mentioned in Chapter 2.
9. **Aims for optimization of the transfer**

- Rather large $V$-size at DSM
- Minimum $a \cdot a$ in kickers $K_1,2$ (for 20 bunch mode only)
- Small enough $H$-size in $K_1,2$
- Small $V$-size at VSM 3 (especially for the 10 bunch mode)
- Minimum $a \cdot a$ in kicker $K_3$
- $H$-dispersion matching at HBM.

**ACKNOWLEDGEMENTS**

I would like to thank Mr G. Guignard who performed the calculations with very great care.
REFERENCES


2) A.E. Ball, R. Billinge, Beam Optics for the Transfer from the Booster to the CERN PS, MPS/Int. MA/B 67-24, 29.12.1967.


20-bunch mode

10-bunch mode

Fig. 2. New scheme
Fig. 3. Achromatic beam injection
<table>
<thead>
<tr>
<th>Number of revolutions</th>
<th>Injection channel</th>
<th>Accelerator aperture</th>
<th>Closed orbit</th>
</tr>
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<td>5</td>
<td><img src="image13" alt="Diagram" /></td>
<td><img src="image14" alt="Diagram" /></td>
<td><img src="image15" alt="Diagram" /></td>
</tr>
</tbody>
</table>

\( S = \alpha \rho \Delta \phi \) closed orbits

\[ \text{Fig. 4. Matched beam injection} \]
Fig. 5. Optical mismatch compensation
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Fig. 6a. Ideal case

Fig. 6b. Actual situation
Fig. 7. Dispersion matching diagrams

Fig. 8. Layout for dispersion matching
Fig. 9. Phase plane diagram for the vertical addition
Fig. 10. Parameters $R_1$ and $R_2$ for the vertical addition.
Fig. 11. Vertical addition chart.