General Effective Actions

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Abstract

We investigate the structure of the most general actions with symmetry group $G$, spontaneously broken down to a subgroup $H$. We show that the only possible terms in the Lagrangian density that, although not $G$-invariant, yield $G$-invariant terms in the action, are in one to one correspondence with the generators of the fifth cohomology classes. For the special case of $G = SU(N)_L \times SU(N)_R$ broken down to the diagonal subgroup $H = SU(N)_V$, there is just one such term for $N \geq 3$, which for $N = 3$ is the original Wess-Zumino-Witten term.

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Effective field theories are increasingly used to understand the dynamics of the Goldstone bosons that result from spontaneous breaking of continuous symmetries. If the action of a theory is invariant under a (compact) Lie group $G$ of global symmetries, spontaneously broken to a subgroup $H$, then the Goldstone fields $\pi^a(x)$ in the effective action parametrize the coset space $G/H$ with $a = 1, \cdots, \dim G/H$, and accordingly transform under linear representations of $H$, but under non-linear realizations of the broken symmetries of $G$. The power of effective field theories arises largely from the fact that the nonlinearly realized broken symmetry allows only a finite number of terms in the action, up to any given order in an expansion in powers of derivatives or momenta.

A general method for constructing invariant non-linear effective actions was given in Ref. [1] for $SU(2)_L \times SU(2)_R$ and was extended to the case of arbitrary $G$ and $H$ in Ref. [2]. But although this method yields the most general $G$-invariant term in the effective Lagrangian, its results are not quite complete. Wess and Zumino [3] showed that fermion loops produce a four-derivative term in the effective Lagrangian for the strong-interaction Goldstone octet that is not invariant under $SU(3) \times SU(3)$, but rather changes under $SU(3) \times SU(3)$ transformations by a total derivative, so that the action is $SU(3) \times SU(3)$ invariant. Subsequently Witten [4] was able to re-express this term as the integral over an invariant Lagrangian density in five dimensions. The WZW action has since then been generalized in Ref. [5] to $G/H$
models with arbitrary $G$ and $H$.

It is natural to ask whether there are any more possible terms in the action (not necessarily related to anomalies in the underlying theory), that, although invariant under a nonlinearly realized symmetry $G$, are not the four-dimensional integrals of $G$-invariant Lagrangian densities. This question seems to us important, as the effective field theory approach is based on our ability to catalog all invariant terms in the action with a given number of derivatives.

The first step is to show that even where the action is not the integral of a $G$-invariant Lagrangian density, its variation with respect to the Goldstone boson fields is an invariant density. The Goldstone boson fields $\pi^a(x)$ enter the action as a parameterization of a general spacetime-dependent $G$-transformation $U(\pi(x))$, so the variation of the action under an arbitrary change in $\pi$ may be written as

$$\delta S[\pi] = \int d^4x \, Tr \left\{ (U^{-1} \delta U) X J \right\} ,$$

where a subscript $X$ or $H$ will denote the terms proportional to the broken and unbroken symmetry generators $x_a$ and $t_i$, respectively, and the coefficient $J$ is a local function of the Goldstone boson fields and their derivatives. Let us work out how $J$ transforms. According to the general formalism of [2], under a global transformation $g \in G$, the Goldstone boson fields undergo the transformation $\pi \to \pi'$, with

$$g \, U(\pi) = U(\pi') \, h(\pi, g) ,$$

(2)
where $h(\pi, g)$ is some element of the unbroken subgroup $H$. Since $S[\pi] = S[\pi']$ for all $\pi$, the variational derivatives are also equal

$$\frac{\delta S[\pi']}{\delta \pi^a} = \frac{\delta S[\pi]}{\delta \pi^a}.$$  

(Note that the derivative is with respect to $\pi$, not $\pi'$, on both sides of the equation.) Using Eq. (1), this is

$$Tr \left\{ \left[ U^{-1}(\pi') \frac{\partial U(\pi')}{\partial \pi^a} \right] X J(\pi) \right\} = Tr \left\{ \left[ U^{-1}(\pi) \frac{\partial U(\pi)}{\partial \pi^a} \right] \frac{\partial \pi^a}{\partial \pi^a} \right\} J(\pi).$$  

To put this in a useful form, take the derivative of Eq. (2) with respect to $\pi^a$, and multiply on the left with $U(\pi')^{-1}$ and on the right with $h^{-1}(\pi, g)$:

$$U^{-1}(\pi') \frac{\partial U(\pi')}{\partial \pi^a} = h(\pi, g) U^{-1}(\pi) \frac{\partial U(\pi)}{\partial \pi^a} h^{-1}(\pi, g) - \frac{\partial h(\pi, g)}{\partial \pi^a} h^{-1}(\pi, g)$$

and so

$$\left[ U^{-1}(\pi') \frac{\partial U(\pi')}{\partial \pi^a} \right] X = h(\pi, g) \left[ U^{-1}(\pi) \frac{\partial U(\pi)}{\partial \pi^a} \right] h^{-1}(\pi, g).$$  

Eq. (3) then becomes

$$Tr \left\{ \left[ U^{-1}(\pi) \frac{\partial U(\pi)}{\partial \pi^a} \right] X \left[ h^{-1}(\pi, g) J(\pi) h(\pi, g) \right] \right\} = Tr \left\{ \left[ U^{-1}(\pi) \frac{\partial U(\pi)}{\partial \pi^a} \right] \frac{\partial \pi^a}{\partial \pi^a} \right\} J(\pi).$$  

From linear combinations of the quantities $[U^{-1}(\pi) \frac{\partial U(\pi)}{\partial \pi^a}] X$ we can form arbitrary linear combinations [6] of the broken symmetry generators $x_a$, so (5) yields the transformation rule for $J$:

$$J(\pi') = h(\pi, g) J(\pi) h^{-1}(\pi, g).$$  

(6)
Following the same arguments that led to (4), we easily see that also

\[
\left[ U^{-1}(\pi') \delta U(\pi') \right] \chi = h(\pi, g) \left[ U^{-1}(\pi) \delta U(\pi) \right] \chi h^{-1}(\pi, g),
\]

so \( Tr \{ (U^{-1} \delta U)_\chi J \} \) is invariant under \( G \).

This result leads to a natural five-dimensional formulation of the theory. As usual, we compactify spacetime to a four-sphere \( M_4 \) by requiring that all fields approach definite limits as \( x^\mu \to \infty \). The operator \( U(\pi(x)) \) therefore traces out a four-sphere in the manifold of \( G/H \) as \( x^\mu \) varies over \( M_4 \). If the homotopy group \( \pi_4(G/H) \) is trivial (as is the case for \( SU(N) \times SU(N) \) spontaneously broken to \( SU(N) \) with \( N \geq 3 \)), or if \( U(\pi(x)) \) belongs to the trivial element of \( \pi_4(G/H) \), then we may introduce a smooth function \( \tilde{\pi}^a(x, t_1) \), such that \( \tilde{\pi}^a(x, 1) = \pi^a(x) \), and \( \tilde{\pi}^a(x, 0) = 0 \). In this way spacetime is extended to a five-ball \( B_5 \) with boundary \( M_4 \) and coordinates \( x^\mu \) and \( t_1 \).

The action may then be written in the five-dimensional form

\[
S[\pi] = \int_{B_5} d^4x \, dt_1 \, \mathcal{L}_1,
\]

where \( \mathcal{L}_1 \) is the \( G \)-invariant density \( Tr \{ (U^{-1} \partial U/\partial t_1)_\chi J \} \). (When \( \pi_4(G/H) \neq 0 \), we may interpolate between \( \pi^a(x) \) and a fixed representative \( \pi_0^a(x) \) of the homotopy class of \( \pi^a(x) \). The difference \( S[\pi] - S[\pi_0] \) is given by the integral over the cylinder \( M_4 \times [0, 1] \) of the same density \( \mathcal{L}_1 \) as in (8) and the arguments to be presented below still hold. In some cases, \( G/H \) may be naturally embedded into a larger space with vanishing fourth homotopy group, as is the case for \( SU(2) \) embedded in \( SU(3) \), considered in [4].)
We next show that this is the integral of a $G$-invariant 5-form on $G/H$.

Consider a general deformation $\pi(x) \to \tilde{\pi}(x; t)$, where $t^i$ are a set of $\dim(G/H) - 4$ free parameters, that along with the $x^\mu$ provide a set of coordinates for $G/H$. The coordinate $t_1$ in (8) can be chosen to be any one of these parameters. We have shown that

$$\frac{\partial S[\tilde{\pi}]}{\partial t_i} = \int_{M_4} d^4x \, \mathcal{L}_i ,$$

where $\mathcal{L}_i \equiv \text{Tr} \{ (\tilde{U}^{-1} \partial \tilde{U} / \partial t^i)_x J \}$ are $G$-invariant functions of $\tilde{\pi}^a$ and its derivatives. The general rules of [2] would allow a wide variety of terms in $\mathcal{L}_i$, but these are limited by integrability conditions. From (9) we see that

$$\int_{M_4} d^4x \left( \frac{\partial \mathcal{L}_i}{\partial t^i} - \frac{\partial \mathcal{L}_j}{\partial t^j} \right) = 0 .$$

Since this integral vanishes for all $\tilde{\pi}(x)$, its integrand must be an $x$-derivative [7]:

$$\frac{\partial \mathcal{L}_i}{\partial t^i} - \frac{\partial \mathcal{L}_j}{\partial t^j} = -\partial_\mu \mathcal{L}^\mu_{ij} .$$

This can be written in the language of differential forms, as $d_t F_1 = -d_x F_2$, where

$$d_t \equiv dt^i \partial_i \quad \text{and} \quad d_x \equiv dx^\mu \partial_\mu$$

and $F_1$ and $F_2$ are the five-forms

$$F_1 \equiv \frac{1}{24} \epsilon_{\mu\nu\rho\sigma} \mathcal{L}_i \ dt^i dx^\mu dx^\nu dx^\rho dx^\sigma \quad \text{and} \quad F_2 \equiv \frac{1}{12} \epsilon_{\mu\nu\rho\sigma} \mathcal{L}^\mu_{ij} dt^i dt^j dx^\nu dx^\rho dx^\sigma .$$

It follows that $0 = d_t^2 F_1 = d_x (d_t F_2)$, so by an extension of Poincaré’s Lemma, in any simply connected patch we will have $d_t F_2 = -d_x F_3$, where $F_3$ is a
five-form $\epsilon_{\mu
u\rho} L_{ijk}^{\mu} dt^i dt^j dt^k dx^\rho$. Continuing in this way, we can construct five-forms $F_4$ and $F_5$ proportional respectively to four and five $dt$ factors, with $d_t F_3 = -d_x F_4$, $d_t F_4 = -d_x F_5$, and $d_t F_5 = 0$. Hence $F \equiv \sum_{N=1}^5 F_N$ is a closed five-form on $G/H$:

$$dF = 0 \quad \quad d \equiv dx + dt.$$  \hspace{1cm} (10)

Also, because $B_5$ has $t_2, t_3$, etc., all constant, Eq. (8) may be written

$$S[\pi] = \int_{B_5} F.$$  \hspace{1cm} (11)

So far, only the term $F_1$ has been shown to be $G$-invariant. The group $G$ acts transitively on the manifold $G/H$, so a $G$ transform of a form is always continuously connected to the original form. Thus the two forms are homotopic and define the same de Rham cohomology class. One can construct a $G$-invariant form in this cohomology class by integrating the form over the group $G$ with the invariant Haar measure [8,9]. This has no effect on (11), since the integral depends only on $F_1$, which is already invariant. Also, one can similarly show that any two invariant $p$-forms in the same cohomology class differ not only by an exterior derivative, but specifically by the exterior derivative of an invariant $(p-1)$-form. Such an exterior derivative term in the five-form $F$ would yield a term in (8) that can be written as the four-dimensional integral of a $G$ invariant density, so the classification of terms in $S[\pi]$ that cannot be so written is now reduced to the problem of finding the fifth de Rham cohomology group $H^5(G/H; R)$ of
the manifold $G/H$ [10].

The fifth de Rham cohomology group is well known where $G/H$ is itself a simple Lie group. For $G = SU(N)$ with $N \geq 3$ (including the case $SO(6) \sim SU(4)$), $H^5(G; R)$ has a single generator

$$\Omega_5 = \frac{-i}{240\pi^2} Tr(U^{-1} dU)^5. \quad (12)$$

(Here and henceforth, we suppress wedges in the exterior product of differential forms, reserving them for the products of cohomology groups.) This is in particular the case for $SU(N) \times SU(N)$ spontaneously broken to $SU(N)$ with $N \geq 3$, where $G/H$ is itself just $SU(N)$. Eq. (12) is the original Wess-Zumino-Witten term, which we now see is indeed unique. All other simple (or $U(1)$) Lie groups have trivial fifth cohomology groups. For the original case [1] of $SU(2) \times SU(2)$ spontaneously broken to $SU(2)$ the cohomology is trivial, so all invariant actions are the integrals of invariant Lagrangian densities.

Where $G/H$ is a product space, we use the Künneth formula [8,9]:

$$H^k(K_1 \times K_2; R) = \sum_{k_1+k_2=k} H^{k_1}(K_1; R) \wedge H^{k_2}(K_2; R), \quad (13)$$

which gives $H^5(G/H; R)$ in terms of the cohomologies of its factors up to degree 5. For this purpose, we need to know that [11,12,13] for all simple Lie groups $G$, $H^k(G; R)$ vanishes for $k = 1, 2, 4$ while $H^3(G; R)$ has a single generator (corresponding to the Goldstone-Wilczek topologically conserved
current \textsuperscript{[14]})
\[ \Omega_3 = \frac{i}{12\pi} Tr (U^{-1} d U)^3. \quad (14) \]
Also \( H^k(U(1); R) \) vanishes for \( k > 1 \), while for \( k = 1 \) it has a single generator
\[ \Omega_1 = -i Tr (U^{-1} d U). \quad (15) \]
Finally, \( H^4(K; R) = R^c \), where \( c \) is the number of connected components of \( K \); for our purposes this just means that if \( H^5(K; R) \) for some space \( K \) has a generator \( \Omega_5 \), then \( H^5(K \times K'; R) \) has the same generator for any \( K' \).

To each generator of \( H^5(G/H; R) \), there corresponds a WZW-like term in the five-dimensional Lagrangian, and an independent coupling constant. In particular, if \( G \) is semi-simple, with precisely \( p \) factors \( SU(N_i) \) with \( N_i \geq 3 \) and all other simple factors with \( H^5 = 0 \), then we have \( p \) different terms of the Wess-Zumino-Witten type, each of which has an independent coupling constant in the action. This result is of course expected for a product of groups, and is known to appear explicitly in the low energy effective action when massive fermions are integrated out of the path integral \textsuperscript{[15]}.

When \( G/H \) is not itself a Lie group, the fifth cohomology group of \( G/H \) may still be obtained from that of \( G \). For any simple group \( G \) and subgroup \( H \), we may construct a ‘projected’ five-form on \( G/H \) that is invariant under local \( H \) transformations \textsuperscript{[5,15,16,17]}, and is given by:
\[ \Omega_5(U; V) = \frac{-i}{240\pi^2} \{ Tr (U^{-1} DU)^5 - 5 Tr W(U^{-1} DU)^3 + 10 Tr W^2 U^{-1} DU \}, \quad (16) \]
where $V$ is the $H$-connection $V = (U^{-1}d U)_{\mu}$, $DU$ is the $H$-covariant derivative $DU = dU - UV$, and the trace is evaluated in any convenient representation of $G$, usually taken as the defining representation. In general, $\Omega_5(U; V)$ is neither closed nor simply related to the generator $\Omega_5(U; 0)$ of $H^5(G; R)$. Rather,

$$d\Omega_5(U; V) = \frac{i}{24\pi^2} d_{rst} W^r W^s W^t$$

(17)

and

$$\Omega_5(U; V) = \Omega_5(U; 0) + \Omega_5(1; V) + d \gamma(U; V),$$

(18)

where $W$ is the field strength $W = dV + V^2$, and $d_{rst}$ is the trace of the symmetrized product of generators $\rho^r$ of $H$, $2d_{rst} = Tr \rho^r \{ \rho^s, \rho^t \}$, which plays a key role in the study of the chiral anomaly in four dimensions [18].

But if $d_{rst} = 0$, then the five form $\Omega_5(U; V)$ is closed, and also each term in the 5-dimensional Chern-Simons term $\Omega_5(1; V)$ for the $H$-valued gauge field $V$ vanishes. The form $\Omega_5(U; V)$ then belongs to the same cohomology class as $\Omega_5(U; 0)$ and it can be shown that there is a one to one correspondence between the fifth cohomology generators of $G/H$ and those of $G$. On the other hand, if $d_{rst} \neq 0$, then the projected form of (16) is not closed and it can be shown that the fifth cohomology is trivial in this case. For example, any coset space of the type $SU(n)/H$ with $n \geq 3$, where $H$ is embedded in $G$ in such a way that $d_{rst} = 0$, has one cohomology generator, given in (16).

It is noteworthy that the simple groups $SU(2); Sp(2N); SO(N), N \geq 7; E_6, E_7, E_8; F_4, G_2$ that have zero fifth cohomology are also those that
have vanishing $d$ symbols. We now see that for such groups, the coset spaces $G/H$ have $H^5(G/H; R) = 0$ for all subgroups $H$. These properties are easily verified for the special case of compact symmetric spaces [12,13]. Also, when rank$(G)=$rank$(H)$, a classic theorem [13] states that all odd cohomology classes vanish. An example of a general class of manifolds $G/H$ with rank$(H) \leq$ rank$(G)$ for which $H^5(G; R) \neq 0$ and $H^5(G/H; R) = 0$ is provided [12] by

$$SU(n)/S(U(k_1) \times \cdots \times U(k_q)) \quad k = \sum_{\alpha=1}^{q} k_{\alpha} \geq 3,$$

with the $U(k_\alpha)$ embedded in $SU(n)$ in such a way that the defining representation of $SU(n)$ transforms also as the defining representation of $S(U(k_1) \times \cdots \times U(k_q))$.

Finally, if $G$ is not simple, and $H$ is a non-trivial subgroup, the cohomology problem can be solved by analyzing the restriction of the $d$-symbols of $G$ to the subgroup $H$ [19]. If $G$ is semi-simple (and $H$ is connected), two types of cohomology generators arise [20]. First, the projected form of (16) is now obtained as a linear combination of $\Omega_5(U; V)$ on each simple component of $G$ with non-vanishing fifth cohomology. Linear combinations for which $d_{rst}$ on the subgroup $H$ vanishes, yield generators of $H^5(G/H; R)$. Second, there may be generators that are linear combinations of products of cohomology generators on $G/H$ of degrees 2 and 3. Generators of degree 2 correspond to the field strength associated with generators of invariant Abelian subgroups of $H$ (i.e. $U(1)$ factors). Generators of degree 3 correspond to the Goldstone-
Wilczek current of (14), projected to $G/H$. When $G$ is not semi-simple and contains extra $U(1)$ factors, there are also linear combinations of products of generators of degree 1 with generators of degrees 1, 2, 3 and 4.

We conclude with a brief discussion of global quantization conditions. Different interpolating maps are generally topologically inequivalent (their equivalence classes being given by $\pi_3(G/H)$), and there is no natural way of choosing one interpolation above another. Witten has argued that the quantum action can be allowed to be multiple valued, provided the action changes additively by integer multiples of $2\pi$. The dependence of interpolation becomes invisible in the quantum theory provided the coupling constants multiplying $\Omega_5$ as normalized in (12) are integers. In the present case, this quantization condition must be enforced on every independent coupling constant multiplying each non-trivial WZW term normalized as in (12).

A slight refinement of this quantization condition is required when $\pi_4(G) = 0$ and $\pi_4(H) \neq 0$. For all simple groups $H$ we have $\pi_4(H) = 0$, except when $H$ is a symplectic group, for which $\pi_4(Sp(2n)) = Z_2$. Whenever $\pi_4(H) = Z_2$, $H$ has a discrete anomaly [21], even though its $d$-symbols vanish identically, and it can be shown that the coupling constant of the corresponding term of $H^5(G/H; R)$ must be quantized in terms of even integers to obtain a single-valued path integral [22].

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References


6. With the exponential parameterization, $U^{-1}(\pi) \partial U(\pi)/\partial \pi^a \rightarrow ix_a$ for small $\pi$, so the quantities $[U^{-1}(\pi) \partial U(\pi)/\partial \pi^a]_X$ span the same space as the $x_a$ for $\pi$ in at least a finite neighborhood of the origin.

7. Here we are using a general theorem, that if the integral over a closed manifold of a local function of a field and its derivatives vanishes for all such fields, then the integrand must be a derivative of another local function of the field and field derivatives. Since we do not know where
this theorem is to be found in the mathematical literature, we have proven it by direct construction of the latter function.


10. The above arguments may easily be carried over to the construction of invariant actions in space-times of dimension $d$, where the allowed non-invariant Lagrangian densities are in one to one correspondence with the generators of $H^{d+1}(G/H; R)$.


20. A discussion of these results will be presented elsewhere.


22. This includes the cases $SU(2) \sim SO(3) \sim Sp(2)$ and $SO(5) = Sp(4)$.