G.Agakishiev\textsuperscript{1}, O.Barannikova\textsuperscript{2}, F.Ceretto\textsuperscript{3},
U.Faschingbauer\textsuperscript{3}, P.Glässel\textsuperscript{4}, E.Kolganova\textsuperscript{1}, G.Ososkov\textsuperscript{1},
Yu.Panebratsev\textsuperscript{1}, J.Rak\textsuperscript{3}, N.Saveljie\textsuperscript{5}, Th.Ullrich\textsuperscript{4},
J.P.Wurm\textsuperscript{5}

A NEW ROBUST FITTING ALGORITHM
FOR VERTEX RECONSTRUCTION
IN THE CERES EXPERIMENT

Submitted to «Nuclear Instruments and Methods»

\textsuperscript{1}JINR, 141980 Dubna, Russia
\textsuperscript{2}JINR and Ivanovo State University, Ivanovo, Russia
\textsuperscript{3}Max-Planck-Institut für Kernphysik, 69117 Heidelberg, Germany
\textsuperscript{4}Universität Heidelberg, 69120 Heidelberg, Germany
\textsuperscript{5}JINR and University of Montenegro, Podgorica, Yugoslavia
1 Introduction

The CERES experiment studies the production of low-mass electron pairs in proton-proton, proton-nucleus, and nucleus-nucleus collisions at the CERN SPS.

Vertex and track reconstruction in the experiment is based on the information of two silicon drift detectors SDD-1 and SDD-2 [1] situated about 9 cm behind the extended, segmented target (see Fig. 1). They cover the full spectrometer acceptance of 8° to 15° for all target disks. The specific target used for the 160 GeV/n Pb beam is segmented into 8 individual disks of 600 μm diameter and 25 μm thickness, equidistantly spaced along the beam direction by 2.9 mm each. This target design allows a larger interaction rate while keeping the photon conversion probability within the spectrometer acceptance low (X/X₀ = 0.37%).

We are dealing here with two sets of hits from each detector. The target and SDD doublet are located in a low magnetic field region and the particle trajectories are straight lines connecting the corresponding hits in SDD-1 and SDD-2.

![Fig. 1. Geometry of the SDD doublet and segmented target. The acceptance of the CERES spectrometer is indicated.](image)

2 Least Squares Formulation of the Problem

Let \((x_{i1}, y_{i1}), i_1 = 1, \ldots, n_1\) and \((x_{i2}, y_{i2}), i_2 = 1, \ldots, n_2\) be the measured points from SDD-1 and SDD-2, respectively, with some number of background points among them. In this case the conventional least-square method (LSM)
for estimating the vertex coordinates \( x_v, y_v \) and \( z_v \) can be based on minimizing the functions

\[
L'(x_v, y_v, z_v) = \sum w_i \epsilon_i'^2,
\]

\[
L''(x_v, y_v, z_v) = \sum w_i \epsilon_i''^2,
\]

where \( \epsilon_i' \) and \( \epsilon_i'' \) are residuals and \( w_i \) are the weights assigned. The value \( \epsilon_i' \) is a measure for the deviation of a SDD-1 hit with coordinates \( x_{i1}, y_{i1}, z_{i1} \) from the line which passes through the vertex \( (x_v, y_v, z_v) \) and its corresponding SDD-2 hit with coordinates \( x_{i2}, y_{i2}, z_{i2} \), in the SDD-1 position

\[
e_i' = \left\{ \frac{z_{i2} - z_{i1}}{z_{i2} - z_v}(x_{i2} - x_v) \right\}^2 + \left\{ \frac{z_{i2} - z_{i1}}{z_{i2} - z_v}(y_{i2} - y_v) \right\}^2 \right)^{\frac{1}{2}}.
\]

The value \( \epsilon_i'' \) is a measure for the deviation of a vertex point \( x_v, y_v, z_v \) from the straight line, given by the corresponding hits from SDD-1 and SDD-2 in the \( z_v \) position of the vertex

\[
e_i'' = \left\{ \frac{z_v - z_{i1}}{z_{i2} - z_{i1}}(x_{i2} - x_{i1}) \right\}^2 + \left\{ \frac{z_v - z_{i1}}{z_{i2} - z_{i1}}(y_{i2} - y_{i1}) \right\}^2 \right)^{\frac{1}{2}}.
\]

The fundamental LSM assumption is that the residuals, or the deviations from the measured point, are normally distributed. However, this is true only in the case of a clean sample which is not contaminated with background. The distribution of residuals including a background fraction \( \epsilon \) can be approximated as in the gross-error model invented by J.W. Tukey:

\[
f(\epsilon) = (1 - \epsilon)\phi(\epsilon) + \epsilon h(\epsilon)
\]

with a normal distribution \( \phi(\epsilon) = \exp(-\epsilon^2/2\sigma^2)/\sigma\sqrt{2\pi} \) and a background \( h(\epsilon) \), which is assumed to be uniform \( (h(\epsilon) = h_0 \) in some interval of the width \( \sigma \ll 1/h_0 \)). The background level \( \epsilon \) varies considerably depending on the experimental environment. It is evident that in this case the weight of distant background points in the LSM functions (1) and (2) is inappropriate and leads to unnecessary large errors in the estimated parameters \( x_v, y_v \), and
A possibility to cut-off large residuals is to introduce a new parameter and take only residuals smaller than this parameter into account. However, in such a case the obtained result would be strongly influenced by the initial values of \(x_v, y_v\) and \(z_v\).

### 3 Summed Gaussian Weights

To avoid the problems mentioned above, the LSM function can be replaced by another one, which introduces a smooth cut-off for distant hits. In particular, in [2] it was proposed to minimize the following function

\[
L(x_v, y_v, z_v) = -\sum e^{\exp[-e_i^2/2\sigma^2]}
\]  

(6)

which has been used up to now for determination of the vertex for CERES data. A suitably chosen \(\sigma\) is assumed to be constant for all data points. The reason for the choice of function (6) can be illustrated by the expansion of \(L\), which shows the similarity between this estimation and an unweighted least square method with \(w_i = (2\sigma^2)^{-1}\) for small \(e_i\),

\[
L(x_v, y_v, z_v) = -n + \sum \frac{e_i^2}{2\sigma^2} - \frac{1}{2} \sum \frac{e_i^4}{4\sigma^4} + \ldots
\]  

(7)

while for larger \(e_i\) the corresponding summands \(L_i\) decrease exponentially, suppressing the influence of strongly deviating hits. As one can see, the second term in the previous expansion corresponds to the function which follows from the unweighted least square method. Since the residuals for the hits that truly belong to the particle track are normal-distributed the obtained minimum can, to some extent, be interpreted on the basis of a \(\chi^2\) distribution. Looking for the minimum of function (6) and differentiating it with respect to each coordinate of the vertex one obtains a system of non-linear equations. The function itself cannot be linearized without loosing its properties. Therefore, a traditional function minimizing package had to be used. The initial values for \(x_v\) and \(y_v\) are set to 0, which should be the position of the center of target disks in the xy-plane. In order to obtain a starting value \(z_v^0\) for the extended target, prior to fitting, a scan is performed by stepping a \(\pm 2\) cm region (in \(z\)) around the center of the target. The found minimum of \(L\) defines the starting value for \(z\).
4 Robust Method for Vertex Reconstruction

As was mentioned before we are dealing with a contaminated data set of points in a sense that some points ("outliers") lie far from the track to be reconstructed. Such outliers can spoil the estimates of vertex coordinates $x_v$, $y_v$ and $z_v$ if their weights $w_i$ are compatible with the weights of useful points. For this case we propose the robust estimation of $x_v$, $y_v$ and $z_v$ based on the iterative reweighted least square estimation of $x_v$, $y_v$ and $z_v$.

4.1 Optimal choice of the weight function in the weighted least square

Since the residuals $e_i$ are non-Gaussian distributed we use a more general approach, i.e., the maximum likelihood (ML) method. An analogous approach was successfully used by P. Huber [3] and leads to the so-called M-estimates of the parameters in question. But we carry out our approach in a different way. Keeping in mind that the corresponding ML-functional is strongly non-linear (leading to considerably computing difficulties), we transform the functional partial derivatives in a view, which allows to reduce the problem to optimal choice of the weight function in the weighted least square sum. The logarithmic likelihood function for measured deviations $e_i$ distributed according to equation (5) is

$$l_n \prod_{i=0}^{n} f(e_i) = \sum_{i=0}^{n} l_n \left( \frac{1 - e_i}{\sigma \sqrt{2\pi}} \exp\left(-\frac{e_i^2}{2\sigma^2}\right) + c h_0 \right).$$  \hspace{1cm} (8)

After differentiating eq. (8) with respect to $x_v$, $y_v$ and $z_v$ and dividing by

$$(1 - e_i c \sigma^{-1}(2\pi)^{-\frac{1}{2}} \exp(-e_i^2/2\sigma^2)),
$$

we obtain a conventional LSM system of normal equations

$$\left\{ \begin{array}{l}
\sum w_{w} \frac{1}{\sigma^2} \frac{\partial f}{\partial x_v} = 0 \\
\sum w_{w} \frac{1}{\sigma^2} \frac{\partial f}{\partial y_v} = 0 \\
\sum w_{w} \frac{1}{\sigma^2} \frac{\partial f}{\partial z_v} = 0
\end{array} \right. \hspace{1cm} (9)$$

with the optimal weight function

$$w_{opt}(i) = \frac{1 + e_i}{1 + c \cdot \exp\left(\frac{e_i}{2}\right)}, \hspace{1cm} (10)$$
where with \( t \) we denote the ratio \( \epsilon / \sigma \). The only parameter

\[
\epsilon = (1 - t)^{-1} \sigma h_0 \sqrt{2\pi}
\]

is the ratio of the mean number of noise hits to the mean number of useful hits within an interval \( \sigma \sqrt{2 \pi} \). Thus it is determined by the contamination of data not in the whole range of the sample but within its essential part where all useful observations are concentrated. Like in our case, the value of \( \epsilon \) is often roughly known in experimental models. The factor \((1 + \epsilon)\) is introduced to fulfill \( w(0) = 1 \). The approximation of \( w_{opt}(t) \) by a fourth order polynomial leads to the famous Tukey's bi-squared weights \[4\]

\[
w(t) = \begin{cases} 
(1 - (t/c_T)^2)^2 & \text{if } |t| < c_T \\
0 & \text{otherwise} 
\end{cases}
\]  

(11)

Here, with \( c_T \) we denote Tukey's constant.

### 4.2 Calculation of vertex coordinates with robust approach

To find the vertex coordinates \( x_v, y_v \) and \( z_v \) we need to solve the LSM system (9) for either of the residuals \( \epsilon'_i \) or \( \epsilon''_i \) given with eq. (3) and (4), respectively.

Dealing with function \( L'(x_v, y_v, z_v) \), i.e., with \( \epsilon'_i \) residuals we get a system of equations which is not simple to solve because they contain a term \( (z_{i2} - z_v) \) in the denominator. To avoid this problem function \( L'(x_v, y_v, z_v) \) can be linearized by multiplying it with an approximately constant term \( (z_{i2} - z_v) \), so we can deal with the function

\[
L'(x_v, y_v, z_v) = \sum w_i \tilde{e}_i^2,
\]

(12)

where \( \tilde{e}_i = (z_{i2} - z_v) \epsilon'_i \).

In this case, evaluating \( \epsilon_i \) from equation system (9) with \( \tilde{e}_i \), one has to solve the corresponding system of linear equations:

\[
\begin{cases}
A'x_v + B'z_v &= C' \\
A'y_v + E'z_v &= F' \\
B'x_v + F'y_v + G'z_v &= H'.
\end{cases}
\]

(13)
where we denote

\[
A' = \sum w_i(z_{i2} - z_{i1})^2; \quad B' = \sum w_i(x_{i1} - x_{i2})(z_{i2} - z_{i1})
\]

\[
C'' = \sum w_i(z_{i2}(x_{i1} - x_{i2})(z_{i2} - z_{i1}) + x_{i2}(z_{i2} - z_{i1})^2)
\]

\[
E'' = \sum w_i(y_{i1} - y_{i2})(z_{i2} - z_{i1})
\]

\[
F'' = \sum w_i(z_{i2}(y_{i1} - y_{i2})(z_{i2} - z_{i1}) + y_{i2}(z_{i2} - z_{i1})^2)
\]

\[
G'' = \sum w_i[(x_{i1} - x_{i2})^2 + (y_{i1} - y_{i2})^2]
\]

\[
H'' = \sum w_i(z_{i2}(x_{i1} - x_{i2})^2 + x_{i2}(z_{i2} - z_{i1})(x_{i1} - x_{i2}) + z_{i2}(y_{i1} - y_{i2})^2 + y_{i2}(z_{i2} - z_{i1})(y_{i1} - y_{i2})).
\]

From the equation system (13) we have

\[
\begin{align*}
x_v &= (C'' - B'z_v)/A' ; \\
y_v &= (F'' - E'z_v)/A';
\end{align*}
\]

where

\[
z_v = (A'H'' - E'F'' - B'C'')/(A'G'' - B''^2 - E''^2).
\]

Dealing with \(e''\), residuals we go straight to the solution. In this case one should solve the following system of linear equations

\[
\begin{align*}
A''x_v + B''z_v &= P'' \\
A''y_v + D''z_v &= Q'' \\
B''x_v + D''y_v + F''z_v &= R'',
\end{align*}
\]

where we denote

\[
A'' = \sum w_i; \quad B'' = \sum w_i x_{i1} - x_{i2}/z_{i2} - z_{i1}; \quad C'' = \sum w_ix_{i1};
\]

\[
D'' = \sum w_i y_{i1} - y_{i2}/z_{i2} - z_{i1}; \quad E'' = \sum w_i y_{i1};
\]

\[
F'' = \sum w_i \left[ \left( \frac{x_{i2} - x_{i1}}{z_{i2} - z_{i1}} \right)^2 + \left( \frac{y_{i2} - y_{i1}}{z_{i2} - z_{i1}} \right)^2 \right];
\]

\[
G'' = \sum w_i \left[ x_{i1} \left( \frac{x_{i1} - x_{i2}}{z_{i2} - z_{i1}} \right) + y_{i1} \left( \frac{y_{i1} - y_{i2}}{z_{i2} - z_{i1}} \right) \right];
\]
\[ P'' = z_1 B'' + C''; \quad Q'' = z_1 D'' + E''; \quad R'' = z_1 F'' + G''. \]

\( z_1 \) is the \( z \)-position of the SDD-1.

From the equation system (15) we have

\[
\begin{align*}
  x_v &= [B''(z_1 - z_v) + C'']/A''; \\
  y_v &= [D''(z_1 - z_v) + E'']/A''.
\end{align*}
\]  \hspace{1cm} (16)

where

\[
  z_v = z_1 + (A''G'' - D''E'' - B''C'')/(A''F'' - B''^2 - D''^2).
\]

The weights in the above expressions are computed iteratively using Tukey’s weight formula

\[
w_i^{(k)} = \begin{cases} 
  (1 - (e_i^{(k)}/(c_T + \hat{\sigma}^{(k-1)}))^2) & \text{if } |e_i^{(k)}| \leq c_T \cdot \hat{\sigma}^{(k-1)} \\
  0 & \text{otherwise}
\end{cases}
\]

where \( e_i^{(k)} \) is the residual of either of the deviations \( e_i' \) or \( e_i'' \); obtained at the \( k \)-th iteration, and \( \hat{\sigma}^{(k-1)} \) is the estimate of variance which is evaluated as [5]

\[
\hat{\sigma}^{(k)} = \frac{\sum w_i^{(k)}(e_i^{(k)})^2}{\sum w_i^{(k)}}.
\]

For our calculations we have varied the constant \( c_T \) and obtained the best resolutions of vertex coordinates for \( c_T \simeq 3 \). Instead of scanning the 8 targets to determine the initial parameters, as it was done for the Gaussian summed weights (SGW), the center of the target region with \( x_v^{(0)} = y_v^{(0)} = 0 \) and \( z_v^{(0)} \) was used as the starting value for the first iteration.

5. Calculation Results

In this section we compare the results for the vertex reconstruction obtained with the Summed Gaussian Weights (SGW) approach and the robust approach. The underlying sample consists of 4900 Pb + Au events. In the following, the results for the SGW method were obtained by using \( e_i' \) residuals in function (6), the results for the robust approach were obtained by using \( e_i'' \) residuals in equation system (9). It should be noted that the usage of \( e_i' \) residuals leads essentially to the same results.
Fig. 2. Reconstructed $z$-coordinates of the vertices fitted with eight Gaussians corresponding to eight target disks obtained with the SGW approach.

Fig. 3. Reconstructed $z$-coordinate of the vertices fitted with eight Gaussians corresponding to eight target disks obtained with the robust approach.

Figures 2 and 3 show the vertex $z$-coordinate distribution for the case of the SGW and the robust method, respectively. As one can see from the resulting histograms, both distributions reflect nicely the target region. Each of the disks is clearly seen as a peak in the distribution. All peaks have Gaussian form, which is illustrated by the fitted Gaussians.

The different resolutions obtained by fitting each of the peaks individually are shown in Table 1 for both cases. The robust approach gives a slightly better resolution for each disk. This was confirmed by tests of both methods with Monte-Carlo generated data, which also resulted in a better resolution of all vertex coordinates, obtained with the robust approach.
Fig. 4. Reconstructed x coordinate of the vertices obtained with the SGW and the robust approach fitted by Gaussian

Fig. 5. Reconstructed y coordinates of the vertices obtained with the SGW and the robust approach fitted by Gaussian

The distributions of the $x$ and $y$ coordinates of the vertex obtained by each algorithm are shown in Fig. 4 and Fig. 5 respectively. Here, the two methods give very similar results both in the mean and in the width of the obtained distributions.

In practice, a reasonable accuracy of the geometric position of the vertex, obtained by the robust approach, is already achieved after several iterations. Fig. 6 shows the distribution of number of iterations per event for the robust weights approach. As one can see, 5 iterations on average are enough to find the minimum. It means that the robust iterative procedure for vertex reconstruction is about an order of magnitude faster than standard general purpose packages for minimization (for example MINUIT[6]).
Fig. 6. Number of iterations per event with the robust approach

Table 1.

$z$-resolutions obtained with SGW and robust fit

<table>
<thead>
<tr>
<th>disk No.</th>
<th>rms (SGW) [(\mu\text{m})]</th>
<th>rms (robust) [(\mu\text{m})]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>307</td>
<td>264</td>
</tr>
<tr>
<td>2</td>
<td>322</td>
<td>287</td>
</tr>
<tr>
<td>3</td>
<td>361</td>
<td>289</td>
</tr>
<tr>
<td>4</td>
<td>327</td>
<td>329</td>
</tr>
<tr>
<td>5</td>
<td>386</td>
<td>344</td>
</tr>
<tr>
<td>6</td>
<td>396</td>
<td>337</td>
</tr>
<tr>
<td>7</td>
<td>409</td>
<td>372</td>
</tr>
<tr>
<td>8</td>
<td>461</td>
<td>395</td>
</tr>
</tbody>
</table>

Fig. 7. Local track accuracy of the silicon drift detectors, radial and azimuthal residuals in SDD-1 for tracks defined by the vertex and a hit in SDD-2 (SGW: solid line, robust: dashed line)

10
Local track accuracy of the silicon drift chambers, radial and azimuthal residuals in SDD-1 for tracks defined by the interaction vertex and a hit in SDD-2, are shown in Fig. 7. Results obtained with both algorithms are shown. As can be seen from the figure, both approaches lead to almost the same distributions of $\Delta r$ and $\Delta \varphi$ residuals ($\sigma_\varphi \approx 6$ mrad, $\sigma_r \approx 100$ $\mu$m). The track resolution results from the combined effect of the intrinsic resolution of the chambers, the vertex resolution and the multiple scattering.

6 Conclusion

We presented results on vertex reconstruction for CERES data obtained with SGW approach and robust approach. Both algorithms give good results clearly reflecting the target region profile. The vertex x-y coordinate resolutions obtained by both methods are almost the same, the z-resolution is somewhat improved by robust method.

The advantage of the robust approach, as an iterative method, is its insensitivity to the choice of initial values for the parameters in question. Starting from the middle of the segmented target we come to the right position after $\sim 5$ iterations. The robust fitting approach allows to reconstruct the vertex coordinates without using standard general purpose packages for minimization. This results in a considerable increase in speed, a very important factor for the time consuming mass-production stage of the analysis of huge data samples.

References

[6] MINUIT - Function Minimization and Error Analysis, CERN Program Library entry D506

Received by Publishing Department on December 31, 1996.
## SUBJECT CATEGORIES
### OF THE JINR PUBLICATIONS

<table>
<thead>
<tr>
<th>Index</th>
<th>Subject</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.</td>
<td>High energy experimental physics</td>
</tr>
<tr>
<td>2.</td>
<td>High energy theoretical physics</td>
</tr>
<tr>
<td>3.</td>
<td>Low energy experimental physics</td>
</tr>
<tr>
<td>4.</td>
<td>Low energy theoretical physics</td>
</tr>
<tr>
<td>5.</td>
<td>Mathematics</td>
</tr>
<tr>
<td>6.</td>
<td>Nuclear spectroscopy and radiochemistry</td>
</tr>
<tr>
<td>7.</td>
<td>Heavy ion physics</td>
</tr>
<tr>
<td>8.</td>
<td>Cryogenics</td>
</tr>
<tr>
<td>9.</td>
<td>Accelerators</td>
</tr>
<tr>
<td>10.</td>
<td>Automatization of data processing</td>
</tr>
<tr>
<td>11.</td>
<td>Computing mathematics and technique</td>
</tr>
<tr>
<td>12.</td>
<td>Chemistry</td>
</tr>
<tr>
<td>13.</td>
<td>Experimental techniques and methods</td>
</tr>
<tr>
<td>14.</td>
<td>Solid state physics. Liquids</td>
</tr>
<tr>
<td>15.</td>
<td>Experimental physics of nuclear reactions at low energies</td>
</tr>
<tr>
<td>16.</td>
<td>Health physics. Shieldings</td>
</tr>
<tr>
<td>17.</td>
<td>Theory of condensed matter</td>
</tr>
<tr>
<td>18.</td>
<td>Applied researches</td>
</tr>
<tr>
<td>19.</td>
<td>Biophysics</td>
</tr>
</tbody>
</table>
Агакишин Г. и др.
Новый метод, основанный на алгоритме робастного фитирования для реконструкции вершин взаимодействия в эксперименте CERES

Представлены результаты робастного подхода для нахождения вершины взаимодействия, основанного на методе наименьших квадратов с перенастройными весами, для эксперимента CERES/NA45 на ускорителе SPS в ЦЕРН по изучению столкновений 160 GeV/u Pb + Au и проведено сравнение нового алгоритма с существующим, основанным на минимизации суммарных гауссовых весов.

В эксперименте CERES/NA45 дуплет цилиндрических кремниевых дрейфовых детекторов (SDD) использовался для восстановления вершины и внешнего трекинга, что позволяет существенно улучшить эффективность нахождения черенковских колец в RICH-детекторах, а также для подавления фона.

Предложенный помехоустойчивый алгоритм решает проблему восстановления вершины при больших загрузках детекторов фонаевыми срабатываниями. Более того, в случаях ультрарелятивистских столкновений тяжелых ионов, где множественность достигает $dN_{ch} / d\eta \simeq 500$, высока вероятность перекрытия сигналов (в детекторах) от близких или пересекающихся треков. Предложенный алгоритм решает указанные выше проблемы, связанные с большими загрузками при высоком быстродействии.

Робастная итеративная процедура, по крайней мере, в несколько раз быстрее, чем подходы, использующие дополнительные пакеты программ для минимизации, обеспечивая такую же точность в восстановлении вершины взаимодействия.

Работа выполнена в Лаборатории высоких энергий ОИЯИ.

Препринт Объединенного института ядерных исследований. Дубна, 1996

---

Agakishiev G. et al.
A New Robust Fitting Algorithm for Vertex Reconstruction in the CERES Experiment

We report on the application of a robust approach based on a reweighted least squares method for vertex reconstruction for the CERES/NA45 experiment at CERN SPS for the case of 160 GeV/u Pb + Au collisions and compare its performance with the presently used algorithm based on the minimization of summed Gaussian weights.

In the CERES/NA45 experiment a doublet of cylindrical silicon drift detectors (SDD) is used for vertex reconstruction and tracking. The tracks identified in the SDD doublet are projected into the first RICH and serve as ring candidates, significantly increasing the ring finding efficiency in the RICH detector. Close background pairs from photon conversion and close Dalitz pairs can be rejected by resolving the tracks in the SDD's or by $dE / dx$.

The vertex reconstruction algorithm has to cope with a background-contaminated environment. Furthermore, in the case of ultrarelativistic heavy-ion collisions where high event multiplicities occur ($dN_{ch} / d\eta \simeq 500$), the probability of confusing hits from close or intersecting tracks is high. The algorithm needs to provide the precise vertex location even at high occupancies, at a reasonable speed. Our findings are that the robust iterative procedure presented here is about an order of magnitude faster than approaches using standard minimization packages.

The investigation has been performed at the Laboratory of High Energies, JINR.

Preprint of the Joint Institute for Nuclear Research. Dubna, 1996
Редактор Э.В.Ивашикевич. Макет Р.Д.Фоминой

Подписано в печать 05.02.97
Формат 60 × 90/16. Офсетная печать. Уч.-изд.листов 1,35
Тираж 380. Заказ 49694. Цена 1620 р.

Издательский отдел Объединенного института ядерных исследований
Дубна Московской области