Elastic p–3He and n–3H scattering with two- and three-body forces

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We report on a microscopic calculation of n–3H and p–3He scattering employing the Argonne v18 and v'18 nucleon-nucleon potentials with and without additional three-nucleon force. An R-matrix analysis of the p–3He and n–3H scattering data is presented. Comparisons are made for the phase shifts and a selection of measurements in both scattering systems. Differences between our calculation and the R-matrix results or the experimental data can be attributed to only two partial waves (3P0 and 3P2). We find the effect of the Urbana IX and the Texas-Los Alamos three-nucleon forces on the phase shifts to be negligible.

Introduction

It is well known that realistic nucleon-nucleon (NN) forces cannot reproduce the 3H and 4He binding energies. Three-nucleon interactions (TNIs) are added to give the necessary small corrections but they still fail to reproduce certain properties of the three nucleon system, most notably the Aγ analyzing power in Nd scattering. Yet the 30% deviation of Aγ can be resolved by tiny changes in the Nd scattering phase shifts (on the order of 0.1 degree) [2, 3, 4]. Furthermore very many operators can contribute to a TNI and the lack of stringent conditions in the three-nucleon system on the structure of the TNI makes its application to other systems desirable. In [5] it was shown that although a realistic NN force can generally reproduce the 4He system, there remain differences, most notably in the analyzing powers. The intensely studied 4He system [6] is unfortunately very difficult to describe due to the many resonances and the 4He bound state. Therefore we investigate the much simpler systems p–3He and n–3H where data exist in the energy range of interest.

We organize the paper in the following way: The next section contains a description of the R-matrix analysis of the p–3He and n–3H data. After that we discuss the Resonating Group Model (RGM) calculation and the model space used. Then we compare R-matrix and RGM results of the phase shifts for various model spaces, and demonstrate the differences for a selection of typical observables. Finally we will discuss the effect of TNIs on the scattering.

I. R-MATRIX ANALYSIS

The R-matrix analysis of the T = 1 part of the A = 4 system began many years ago with an analysis of p+3He scattering data below 20 MeV incident proton energy [7]. Only one solution described all the data included, namely the one with δ(3P1) > δ(1P1), giving a χ2 per degree-of-freedom value of about 1.23. Later, this solution was reflected to the n+3H system [8], using a simple energy shift to correct for the short-ranged Coulomb differences between 4Li and 4He. This shift of the R-matrix eigenenergies (ΔEα = −0.86 MeV) for the 4Li system was adjusted by hand to reproduce approximately the n−t total cross section measurement of Phillips, et al. [9].

More recently, we added to the p+3He analysis analyzing-power and spin-correlation data measured at energies between 4 and 10 MeV by Alley and Knutson [10]. These high-precision data made small changes in the phase shifts, but did not alter the qualitative nature of the solution, shown as the solid curves in Figs. 1-5. With these data we achieve a χ2 per degree-of-freedom of 1.27.

Finally, we combined the p+3He data with the n−t total cross sections of ref. [11] and fitted both reactions simultaneously. The single energy shift used earlier did not give a particularly good fit to the total cross sections, so we allowed the low-lying eigenenergies in the 4He system to adjust separately from those in the 4Li system, while keeping the reduced-width amplitudes the same in both systems. This resulted in a χ2 per degree-of-freedom of 1.59. Most of the increase came from the fit to the n−t total cross sections, which have uncertainties on the order of 0.2%. The χ2 per point of the fit to the p+3He data increased only from 1.22 to 1.24. The analysis is based on 1447 data points having proton energies between 1.01 and 19.7 MeV and neutron energies between 0.06 and 20.06 MeV, and allows a maximal orbital angular momentum Lmax = 4. The 4Li resonance energies still agree with those of ref. [12] which contains a brief description of the R-matrix method.

II. RGM AND MODEL SPACE

We use the Resonating Group Model [13, 14, 15] to compute the scattering in the 4H and 4Li systems using the Kohn-Hulthén variational principle [16]. The main technical problem is the evaluation of the many-body matrix elements in coordinate space. The restriction to a Gaussian basis for the radial dependencies of the wave function allows for a fast and efficient calcu-
luation of the individual matrix elements \[11\] \[13\]. How-
never, to use these techniques the potentials must also be
given in terms of Gaussians. In this work we use suitably
parametrized versions of the Argonne \[15\] and \[16\] \(NN\) potentials and the Urbana IX \[17\] and Texas - Los
Alamos \[17\] TNNs.

The inclusion of an additional TNI requires an order of
magnitude more computing power than the realistic \(NN\)
forces alone. It is therefore very fortunate that enough
data exists at low energies for the comparatively simple
isospin \(T = 1\) systems \(^4\)H and \(^3\)Li to allow for a com-
parison between calculation and the experimental data
or the \(R\)-matrix analysis thereof.

In the \(^3\)Li system we use a model space with three
two-fragment channels, namely the \(p - 3\)He, the \(3\)H–(pp)
and the \(^2\)H(S = 0) – (pp) channels. The last two are an
approximation to the three- and four-body breakup
channels that cannot in practice be treated within the
RGM. The \(^4\)Li is treated as four clusters in the framework
of the RGM to allow for the required internal orbital
angular momenta of \(^3\)He or \(^2\)H.

For the scattering calculation we include the \(S\), \(P\)
and some of the \(D\) wave contributions to the \(J^\pi = 0^+, 1^+, 0^-, 1^\) − and \(2^\) − channels. From the \(R\)-matrix
analysis these channels are known to give essentially the
experimental data. The full wave function for these chan-
nels contains over 100 different spin and orbital angular
momentum configurations, hence it is too complicated to
be given in detail. To give an impression of the model
space we will describe the important structures of \(^3\)He as
used in the present work.

The dominant spin \(S = 1/2\) configurations of \(^3\)He in-
clude those without angular momentum and with two \(D\) waves
coupled to total angular momentum \(L = 1\).

\(S = 3/2\) in turn occurs together with a single \(D\) wave on
each of the Jacobi coordinates or two \(D\) waves coupled to
\(L = 2\). Each of these configurations uses a set of one to
three Gaussians whose width parameters were obtained
by a non-linear optimization using a genetic algorithm
\[18\]. In this small, 29-dimensional model space we still
achieve -6.37 MeV binding energy, an \(r_{rms}\) radius of 1.78
fm and a \(D\) state probability of 7.7% for the \(^3\)He using
\(A_{18}\). This must be compared to -6.92MeV known from
Faddeev calculations \[19\] whereas we find -6.88MeV in
a rather large model space.

This representation of \(^3\)He, together with a \(S = 1/2\)
and \(L = 2\) configuration (excited state) and the \((pn)\) –
(pp) fragmentations, form the basis of our \(p - 3\)He scat-
ering calculation. Once the fragment wave functions are
fixed the scattering problem is solved with our RGM code
relying on the Kohn-Hulthén variational principle \[14\]:

\[
\delta \left( \langle \Psi | H - E | \Psi \rangle - \frac{1}{2} a_{ij} \right) = 0,
\]

where \(a_{ij}\) denotes the reactance matrix.

The model space described above (consisting of two
to four physical scattering channels for each \(J^\pi\)) is by
no means sufficient to find reasonable results. So-called
distortion or pseudo-inelastic channels \[13\] have to be
added to improve the description of the wave function
within the interaction region. Accordingly, the distortion
channels have no asymptotic part.

For practical purposes it is obvious to reuse some of
the already calculated matrix elements as additional dis-

tortion channels. In that way we include all the pos-
itive parity states of the three-nucleon subsystem with
\(J_3^\pi \leq 5/2^+\) in our calculation. However, it was recently
pointed out by A. Fonseca \[20\] that states having a nega-
tive parity \(J_3^\pi\) in the three-nucleon fragment increase the
\(n - 3\)H cross section noticeably. Therefore we also added
the appropriate distortion channels in a similar complex-
ity as in the \(J_3^\pi\) case to our calculation, thereby doubling
the size of the model space.

### III. PARTIAL WAVE ANALYSIS

The \(R\)-matrix analysis of the \(T = 1\) part of the \(4N\)
system described above represents the currently available
\(p - 3\)He and \(n - t\) data very well, as the solid lines in figs.
\[2\] \[3\] and \[4\] show. In addition it provides elastic scattering
phase shifts (solid line in figs. \[1\] and \[1\]) that in turn suffice
to describe the experimental data. Therefore we take
these phase shifts as benchmarks against which to test
our calculations.

In a calculation of elastic \(p - 3\)He scattering using the
full model space described above (including both the \(J_3^\pi\)
and \(J_3^\) distortion channels) we find for both the Ar-
gonne \(v_{18}\) and \(v_6^*\) (not shown) interactions in general
phase shifts very similar to those given by the \(R\)-matrix
analysis (see figs. \[1\] and \[1\]). The \(S\)-waves are negative
due to the underlying Pauli-forbidden states, whereas
all the \(P\)-waves are positive. The \(J_3^\pi\) part increases the
model space and therefore reduces the repulsion in the
\(S\)-waves and increases the (attractive) interaction in the
\(P\)-waves. Hence, all phase shifts have to become more
positive. The \(^1\)S\(_0\) and \(^3\)S\(_1\) phase shifts (see fig. \[1\])
depend only weakly on the \(J_3^\pi\) part of the model space
(long dashed line as compared to dashed line), because
the central terms of the \(NN\)-interaction can only con-
nect to a few states of the increased model space. The
\(^3\)P\(_0\) and \(^3\)P\(_2\) phase shifts deviate markedly from the \(R\)-

matrix results (see fig. \[1\]). Without the \(J_3^\) components
(long dashed line) the major difference between our cal-
culation and the \(R\)-matrix results was the \(^3\)P\(_2\) matrix
element being too small, whereas the \(^3\)P\(_0\) results were
very close. The additional \(J_3^\) distortion channels reduce
by half the difference for \(^3\)P\(_2\), but also raise the \(^3\)P\(_0\) phase
considerably beyond its \(R\)-matrix values.

It has already been pointed out previously \[2\] that the
calculated \(^3\)P\(_2\) matrix element in the \(^4\)He system is too
small (then using the realistic Bonn \(NN\) interaction \[21\]).
It was shown that the manual change of \(^3\)P\(_2\) to its \(R\)-
matrix value removed most of the discrepancies in the
description of the analyzing power. The \(J_3^\) part of the
\(p - 3\)He model space in the present calculation acts mainly
FIG. 1: $p-^3\text{He}$ $P$ wave scattering phase shifts. RGM calculations ($A_{v18}$, dashed line, restricted to positive parity in the three-nucleon subsystem, long dashed, and $A_{v18}+\text{Urbana IX}$, dotted) compared to the $R$-matrix analysis (solid line).

IV. COMPARISON WITH DATA

The increase of the $^3P_2$ matrix element due to the $J^-$ components suffices to describe the total and differential cross sections (figures 2 and 5) in both the $n-^3\text{H}$ and $p-^3\text{He}$ scattering. Again, the improvement is due to the larger $^3P_2$ matrix element, the effect of $^3P_0$ being too large on the cross sections is negligible (on the order of 1%) because of the small statistical weight.

There remain deviations of the differential cross section at small and large scattering angles as well as in the low energy behavior of the total cross section. In the first case, we attribute the differences at large scattering angles to be an effect of the model space lacking higher partial waves, and our calculation agrees with the results obtained by restricting the $R$-matrix partial waves to those included in the RGM. Even compared to these results, our calculation differs in the region of the interference between the Coulomb and the strong interaction at small scattering angles. This discrepancy cannot be attributed to a single partial wave, but originates from
similar to the findings in [20]. We know from [3] that the analyzing power is sensitive to the $^3P_2$ matrix element and it was surprising that even with the $J_y$ distortion channels increasing the $^3P_2$ phase shift, we cannot improve the description of $A_y$. The reason is that most of the expected improvement due to the larger $^3P_2$ matrix element is canceled by the increased $^3P_0$ phase shift that decreases $A_y$, as can be seen when we lower $^3P_0$ to its $R$-matrix value (fig. 4, where the values of $A_y$ are displayed close to the maximum). Increasing $^3P_2$ further to the $R$-matrix value will then raise $A_y$ again, but still not to the experimental data. Only a small part of the remaining difference is due to higher partial waves that are not included in the model space. Only a small part of the remaining comes from small differences in the other partial waves.

V. THREE-NUCLEON FORCE EFFECTS

Both the $\text{Av}_6$ and $\text{Av}_{18}$ $NN$ forces yield essentially the same phase shifts, and we convinced ourselves that the Gaussian parametrization of the Bonn interaction [21] also agrees with these results. In the previous section we showed that we cannot reproduce the $p-^3\text{He}$ analyzing power $A_y$ mainly due to the too small splitting of $^3P_2$ and $^3P_0$ (fig. 7). Therefore we included an additional TNI in our calculations. For $\text{Av}_{18}$ we chose the $2\pi$-exchange model implemented as the Urbana IX force [14] and for $\text{Av}_6$ the Texas-Los Alamos (TLA, [17]) three-nucleon force. The latter includes only short range operators and is claimed to resolve the $A = 3 A_y$ problem for a certain choice of its operator strength ($c_1 = 3$, [17]). In our calculation we chose this factor to be $c_1 = 1.5$ to give reasonable $^3\text{He}$ binding energies in the model space used.

In the case of the $\text{Av}_{18}$ and Urbana IX interaction, the additional TNI leads to the expected improvement of the binding energy. For the small $^3\text{He}$ wave function we find $-6.88$ MeV instead of the $-6.37$ MeV for $\text{Av}_{18}$ alone and the TNI increases the $D$ state probability $P_D$ somewhat (from 7.7% to 8.2%). In a rather large model space we find a binding energy of $-7.72$ MeV and kinetic energy of $-50.05$ MeV, close to the values given in [19].

Both TNIs are unfortunately very time-consuming to evaluate and neither of them contributes significantly to the phase shifts (fig. 4). Especially the Urbana IX interaction does not affect the phase shifts at all while still improving the binding energy of $^3\text{He}$. The $^3P_2 - ^3P_0$ splitting is somewhat improved by the TLA force but still far from its $R$-matrix values (fig. 4). Yet this improvement can already be seen in the analyzing power (fig. 5), where the calculation including the TLA force gives significantly better results. Since the Urbana force left the phase shifts almost unchanged it affects the analyzing power (fig. 5) only marginally. These effects can be more clearly seen in fig. 6, where the analyzing power close to the maximum is shown for the different models.

Since the $2\pi$ exchange operators leave the scattering...
phase shifts unchanged, we considered the $V_3^*$ operators proposed in [24] in addition to the Urbana IX interaction. The $V_3^*$ uses a pion-exchange between the third particle and the center of mass of the $NN$ system, which in turn uses the full $NN$ interaction, but only for $T = 1$ and $S = 1$. The $V_3^*$ therefore acts predominantly in the $P$-waves. In order to study the effects of this interaction within a relatively simple operator structure, we approximated the full $t$-matrix by using the central part of the Argonne $v_{18}$ potential as the $NN$-force part in $V_3^*$, and chose its strength to leave the $^3$H binding energy unchanged while reproducing the $^3P_2$ matrix element at $E_{cm} = 2.4$ MeV.

The full model of the $V_3^*$ NNN interaction has only a small influence on the binding energy of $^3$H [28], which also holds for $^4$He with the $V_3^*$ potential in our calculation. However, these operators have a large effect on the $P$-wave phase shifts. If we choose the strength of $V_3^*$ so that the $^3P_2$ matrix element is reproduced and the $^3$H binding remains unchanged, the $^3P_0$ phase shift...
unfortunately increases again, and therefore part of the improvement due to the larger $^3P_2$ matrix element is cancelled. Nevertheless, Urbana IX and $V_3^*$ together achieve an $A_y$ of 0.35 at $E_{cm} = 4.1$ MeV, a much larger effect than the Urbana IX force alone achieved.

VI. CONCLUSIONS

We have discussed a new $R$-matrix analysis of the currently available experimental data. The phase shifts calculated in this analysis were compared to an RGM calculation of $p-^3$He scattering. We showed that realistic $NN$ interactions describe most of the phase shifts quite well but fail to reproduce the $^3P_2$ and $^3P_0$ phase shifts. The calculated splitting between these two channels is much too small, and neither the Urbana IX nor the TLA three-nucleon force is able to improve the splitting significantly. In fact, it is more important to include in the calculation negative parity states of the three-nucleon subsystem than one of these two TNIs. These findings show that new contributions to the $NNN$ force acting on the $P$-waves should be considered, like an LS type TNI, as proposed in [29] for the $N-d$ analyzing powers, or the $V_3^*$ operators proposed in [27].

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