Supersymmetric Källen-Lehmann representation

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Abstract

We find the general form of supersymmetric invariant two point functions. By imposing supersymmetric positivity we obtain the general supersymmetric Källen-Lehmann representation.

1 Introduction

The Källen-Lehmann representation [1] gives the general form of the two point functions in quantum field theory in terms of a positive spectral measure. It summerizes two of the most important ingredients of quantum field theory: Poincare invariance and positivity (combined with a certain amount of singularity which is indispensable in quantum field theory see [2]). For a (free or interacting) complex scalar field \( \Phi(x) \) it is

\[
(\Phi(x)\Omega, \Phi(y)\Omega) = (\Omega, \Phi^+(x)\Phi(y)\Omega) = \int_0^\infty dm^2 \mu(m^2) \Delta_+(x - y, m^2)
\] (1.1)

where \( \Delta_+ \) is the familiar Pauli-Jordan function (for notations see [1]) and \( \mu(m^2) \) a positive measure of polynomial increase. As noted above the Källen-Lehmann representation is a consequence of Poincare invariance and positive definiteness of the distribution \( (\Phi(x)\Omega, \Phi(0)\Omega) \). A straightforward proof follows in the distribution-theoretic approach by using the general form of Lorentz invariant distributions [3]. It is surprising that a similar representation for supersymmetric theories doesn’t appear in the literature. Here we want to fill this gap by obtaining the Källen-Lehmann representation for
supersymmetric (scalar) fields. The paper is structured as follows: in the second section we find the most general supersymmetric invariant two point functions, in the third section we discuss the supersymmetric positivity and in the forth section we obtain the Kallen-Lehmann representation. We conclude the paper with a remark concerning the supersymmetric generalized free field.

2 Supersymmetric invariant two point functions

Let $F(z)$ and $F(z_1, z_2)$ be supersymmetric functions (or distributions) of variables $z(x, \theta, \bar{\theta})$ (we use the notations and conventions of [4]; the notations in [5] are the same as in [4] with the exception of the Pauli $\sigma_0$ which is minus identity in [4] and plus identity in [5]). Such functions can be expanded in powers of $\theta, \bar{\theta}$. For example

$$F(z) = f(x) + \theta \varphi(x) + \bar{\theta} \bar{\chi}(x) + \theta^2 m(x) + \bar{\theta}^2 n(x) +$$
$$\theta \sigma^i \bar{\theta} \nu_i(x) + \theta^2 \bar{\theta} \lambda(x) + \theta^2 \bar{\theta} \psi(x) + \bar{\theta}^2 \theta^2 d(x)$$

(2.1)

A similar expression can be given for $F(z_1, z_2)$. Now suppose $F(z_1, z_2)$ is invariant under simultaneous supersymmetric transformations of the variables $z_1, z_2$.

We want to find the two-point functions (i.e. functions $F(z_1, z_2)$ dependind on two variables $z_1 = (x_1, \theta_1, \bar{\theta}_1), z_2 = (x_2, \theta_2, \bar{\theta}_2)$) which are invariant under an arbitrary simultaneous supersymmetric transformation of the variables. Denoting by $Q = (Q_1), \bar{Q} = (\bar{Q}_1)$ the generators of the supersymmetric transformations (up to a factor of $i$):

$$iQ_\alpha = \partial_\alpha - i \sigma^l_{\alpha a} \bar{\theta}^a \partial_l$$

(2.2)

$$i\bar{Q}_\dot{\alpha} = -\bar{\partial}_{\dot{\alpha}} + i \theta^a \sigma^l_{a \alpha} \partial_l$$

(2.3)

simultaneous supersymmetric invariance imply

$$(Q_1 + Q_2)F(z_1, z_2) = 0$$

(2.4)

$$(\bar{Q}_1 + \bar{Q}_2)F(z_1, z_2) = 0$$

(2.5)

where $Q_1, Q_2$ and $\bar{Q}_1, \bar{Q}_2$ act on the variable $z_1$ and $z_2$ respectively.

In order to solve these equations we introduce new variables $\theta = \frac{1}{2}(\theta_1 + \theta_2)$
and \( \zeta = \theta_1 - \theta_2 \) together with their conjugates as well as (by translation invariance) the difference variable \( x = x_1 - x_2 \) (a somewhat similar argument but in a different context also appears in [6]). Note that by introducing the difference variable \( x = x_1 - x_2 \) the derivative \( \partial_1 \) in \( Q_2, \bar{Q}_2 \) taken with respect to the second variable changes sign such that in the new variables equations (2.4),(2.5) take the form

\[
\left( \frac{\partial}{\partial \theta^{\alpha}} - i \sigma^l_{\alpha \alpha} \bar{\zeta}^\alpha \partial_l \right) F = 0 \quad (2.6)
\]

\[
\left( \frac{\partial}{\partial \bar{\theta}^{\dot{\alpha}}} - i \zeta^\alpha \sigma^l_{\alpha \dot{\alpha}} \partial_l \right) F = 0 \quad (2.7)
\]

where \( F \) is a function depending on the variables \( x = x_1 - x_2, \theta, \bar{\theta}, \zeta, \bar{\zeta} \). We want to solve this system of partial differential equations in mixed commutative and non-commutative variables. A first (trivial and from the physical point of view uninteresting) solution for \( F \) is a constant. Other solutions can be obtained in a two step procedure by using in the first step the equation (2.6) to factorize from \( F \) the exponential \( \exp(i \theta \sigma^l \bar{\zeta} \partial_l) \). We write

\[
F = \exp(i \theta \sigma^l \bar{\zeta} \partial_l) D
\]

The first equation (2.6) implies \( \frac{\partial}{\partial \theta^{\alpha}} D = 0 \) which means that \( D \) is independent of \( \theta^\alpha, \alpha = 1, 2 \). In the second step we write

\[
D = \exp(-i \zeta^\alpha \sigma^l \partial_l) E
\]

and conclude as above that \( E \) is independent not only of \( \theta \) but also independent of \( \bar{\theta} \). These two exponentials cover the dependence of \( F \) from the variables \( \theta \) and \( \bar{\theta} \). The residual dependence in \( E \) is in \( \zeta, \bar{\zeta} \) and \( x = x_1 - x_2 \).

Altogether the general solution of the equations (2.6),(2.7) in the \( x, \theta, \bar{\theta}, \zeta, \bar{\zeta} \) variables is (with the exception of the constant solution) of the form

\[
F(x, \theta, \bar{\theta}, \zeta, \bar{\zeta}) = \exp[\pm i(\zeta^l \theta - \theta^l \bar{\zeta}) \partial_l] E(x, \zeta, \bar{\zeta}) \quad (2.8)
\]

where from invariance considerations \( E \) is restricted to

\[
E(x, \zeta, \bar{\zeta}) = E_1(x) + \zeta^2 E_2(x) + \bar{\zeta}^2 E_3(x) + \zeta^l \zeta \partial_l E_4(x) + \zeta^2 \bar{\zeta}^2 E_5(x) \quad (2.9)
\]
with \( E_i(x) = E_i(x_1 - x_2), i = 1 \ldots 5 \) Lorentz invariant functions (or distributions). We obtain in this way (with the exception of the trivial constant) a total of 5 linearly independent (in the Grassmann variables) invariants. Note the \( \pm \) sign in front of the exponential in (2.8). The supersymmetric invariance in (2.4),(2.5) is formulated in terms of the left multiplication. It produces the minus sign in (2.8). We could have used right multiplication in which case the equations (2.4),(2.5) are replaced by

\[
(D_1 + D_2)F(z_1, z_2) = 0 \\
(\bar{D}_1 + \bar{D}_2)F(z_1, z_2) = 0
\]

where \( D, \bar{D} \) are the covariant derivatives [4]. By solving these equations in the same manner as the equation (2.6),(2.7) we produce the plus sign in the exponential in (2.8).

Although the problem seems to be solved it is clear that we didn’t obtain yet the result we wish to obtain. Indeed there are a bunch of invariants easy to construct by hand and which are intensively used in physics. We will give them below. We want to recognize them among the solutions of the partial differential equations above and to investigate if there are some others invariants too. We start with the ”supersymmetric Pauli-Jordan” function

\[
K_0 = \delta(\theta_1 - \theta_2)\delta(\bar{\theta}_1 - \bar{\theta}_2)\Delta_+(x_1 - x_2)
\]

It is supersymmetric invariant and it can be located among the invariants deduced as solutions of the partial differential equations as the one containing the \( \delta \)– function product \( \zeta^2 \bar{\zeta}^2 \). Certainly the exponential in (2.8) disappears because of the \( \delta \)–function. In order to see directly that \( K_0 \) is supersymmetric invariant we can apply on both variables \( z_1, z_2 \) the supersymmetric transformations

\[
x'^m = x^m + i\theta^m \bar{\epsilon} - i\epsilon^m \bar{\theta} \\
\theta' = \theta + \epsilon, \bar{\theta}' = \bar{\theta} + \bar{\epsilon}
\]

with increments \( \epsilon_1, \epsilon_2 \). The invariance follows after disregarding third powers of \( \theta_1 - \theta_2 \) and \( \bar{\theta}_1 - \bar{\theta}_2 \) variables. This means that the even Grassmann part in (2.10) has no influence on \( \Delta_+(x_1 - x_2) \) which immediately implies the invariance of \( K_0 \).

We continue by giving other invariants used in physics which can be constructed from \( K_0 \). For simplifying the notations in the rest of the paper let us denote
\[
\Delta = \Delta(\theta_1 - \theta_2, \bar{\theta}_1 - \bar{\theta}_2) = \delta(\theta_1 - \theta_2)\delta(\bar{\theta}_1 - \bar{\theta}_2) = (\theta_1 - \theta_2)^2(\bar{\theta}_1 - \bar{\theta}_2)^2
\]

(2.12)

First of all note that by replacing in the supersymmetric Pauli-Jordan function \(K_0\) the function \(\Delta_+\) by an arbitrary Poincare invariant function (or distribution) \(G_0\) we obtain other supersymmetric invariants. All of them belong to the same family whose elements will be denoted by \(\Delta G_0 = P_0 \Delta G_0\) where for notational reasons to be seen later we define \(P_0 = 1\). Beside \(K_0\) other five (families of) invariants can be obtained from \(\Delta G_0\) (with different \(G_0\)) by applying to it the covariant and (invariant) derivatives \(D^2, \bar{D}^2, D^2 \bar{D}^2, \bar{D}^2 D^2\) and \(D^a \bar{D}^\alpha D_\alpha\). Because \(\Delta G_0\) depends on the variables \(x_1, \theta_1, \bar{\theta}_1\) and \(x_2, \theta_2, \bar{\theta}_2\) we have to specify on which variables we apply the covariant derivatives. It turns out that this is a minor problem. We will let them act on the first variables \(x_1, \theta_1, \bar{\theta}_1\) and take the plus sign in the exponential (2.8). Other choices give no new results.

Whereas as discussed above it is relatively easy to see that \(\Delta G_0\) is of the form (2.8),(2.9) it is not totally obvious that the members of the five derived families above are of the form (2.8),(2.9) First note that a compact way to write them down is by using the formal projection operators

\[
P_c = \frac{1}{16} \frac{D^2 \bar{D}^2}{\Box} \quad (2.13)
\]

\[
P_a = \frac{1}{16} \frac{\bar{D}^2 D^2}{\Box} \quad (2.14)
\]

\[
P_T = -\frac{1}{8\Box} D^a \bar{D}^\alpha D_\alpha = -\frac{1}{8\Box} \bar{D}_\alpha D^2 \bar{D}^\alpha \quad (2.15)
\]

and two more operators

\[
P_+ = \frac{D^2}{4\Box^{1/2}} \quad (2.16)
\]

\[
P_- = \frac{\bar{D}^2}{4\Box^{1/2}} \quad (2.17)
\]

We do not comment on some technical aspects related to the d’alembertian in the denominator.

The operators \(P_i, i = c, a, T, +, -\) form a closed algebra under multiplication. We recognize the invariants which we have constructed by hand to be of the form \(P_i \Delta G_i, i = 0, c, a, T, +, -\) where \(G_i\) are Poincare invariants. The operators \(P_1, P_2, P_3\) satisfy the relations [4]
\[ P_c + P_a + P_T = 1 \quad (2.18) \]

\[ (P_c + P_a - P_T)\delta(\theta_1 - \theta_2)\delta(\bar{\theta}_1 - \bar{\theta}_2) = \]

\[ = 4\exp[i(\theta_2\sigma^l\bar{\theta}_1 - \theta_1\sigma^l\bar{\theta}_2)\partial_l] \quad (2.19) \]

Now we come back to (2.8), (2.9). By promoting the term \( \zeta \sigma^l \bar{\theta}_1 \) in (2.9) to the exponential and adjusting in the linear hull of \( E_1, \zeta^2 E_2, \zeta^2 E_3, \zeta^2 \bar{\zeta}^2 E_5 \) we finally obtain the following result:

**Theorem 1** The two-point supersymmetric invariant functions (distributions) are linear combinations of \( P_i G_i(x_1 - x_2) \), where \( G_i(x), i = c, a, T, +, - \) are Lorentz-invariant functions (distributions).

In order to prove this assertion we need the following identities [4]

\[
D^2_1(\theta_1 - \theta_2)^2 = -4\exp[-i(\theta_1 - \theta_2)\sigma^l\bar{\theta}_1\partial_l^1] \\
\bar{D}^2_1(\bar{\theta}_1 - \bar{\theta}_2)^2 = -4\exp[i\theta_1\sigma^l(\bar{\theta}_1 - \bar{\theta}_2)\partial_l^1] \\
\bar{D}^2_1D^2_1(\theta_1 - \theta_2)^2(\bar{\theta}_1 - \bar{\theta}_2)^2 = 16\exp[i(\theta_1\sigma^l\bar{\theta}_1 + \theta_2\sigma^l\bar{\theta}_2 - 2\theta_1\sigma^l\bar{\theta}_2)\partial_l^1] \\
D^2_1\bar{D}^2_1(\theta_1 - \theta_2)^2(\bar{\theta}_1 - \bar{\theta}_2)^2 = 16\exp[-i(\theta_1\sigma^l\bar{\theta}_1 + \theta_2\sigma^l\bar{\theta}_2 - 2\theta_2\sigma^l\bar{\theta}_1)\partial_l^1] 
\]

where for the convenience of the reader we provided the derivatives with the index of the acted on variable (the first one). These relations exactly provide the connection between the invariants obtained by solving the partial differential equations (2.6), (2.7) and the invariants given with the help of the covariant derivatives. In particular from the \( \zeta^2, \bar{\zeta}^2 \)-terms in (2.9) we obtain the \( P_x \)-invariants and from the rest of the terms in (2.9) we obtain the \( P_i, i = c, a, T \)-invariants.

In other words we have rigorously proved that the expressions \( P_i \Delta G_i \) are supersymmetric invariants (a fact which was clear from the very construction of these expressions) and, more important, that there are no other two-point supersymmetric invariants. Stated in a clear way the statement shows that besides the supersymmetric invariant two-point functions which already appear in the physical literature (and are easy to construct by hand) there are no others.
3 Supersymmetric positivity

Now after we obtained the invariant two point functions $F(z_1, z_2)$ we have to ask the important question which ones of them are positive definite in superspace i.e. for which kernels $F(z_1, z_2)$ the (hermitean) sesquilinear form

$$\int d^8z_1 d^8z_2 \overline{f(z_1)} F(z_1, z_2) g(z_2)$$

is positive for $f = g$. The bar includes besides the usual complex conjugation also the Grassmann conjugation. Here $d^8z = d^4xd^2\theta d^2\bar{\theta}$.

First of all let us remark that for $F = F(z_1, z_2) = P_\pm \Delta(\theta_1 - \theta_2, \bar{\theta}_1 - \bar{\theta}_2) G_\pm$ the integral (3.1) with $f = g$ does not have definite sign even if $G_\pm$ is a positive definite distribution. This can be proven by decomposing $f$ with the help of the projection identity $P_c + P_a + P_T = 1$ and using the properties of the covariant derivatives $D, \bar{D}$. It means that from the five invariants associated to $P_i, i = c, a, T, +, -$ only three of them (those associated to $P_i, i = c, a, T$) have to be taken into consideration. It was shown in [7] that all of them can give rise to positive definite sesquilinear forms (scalar products). Indeed $P_i \Delta G_i(z_1, z_2), i = c, a$ as well as $- P_T \Delta G_T$ are positive definite supersymmetric kernels if $G_i(x_1 - x_2)$ are positive definite in the usual sense. Note the minus sign in front of $P_T$ which seems to be at odd with the projection identity $P_c + P_a + P_T = 1$ and plays a considerable role because it induces a natural, inherent Krein structure of supersymmetries. The proof of these assertions [7] is mainly by computation using the positivity of $-p^2$ and $\sigma p$ for the momentum $p$ in the forward light cone and, as far as the $P_T \Delta G_T$ is concerned, by also using a rather old argument which appeared in the study of the Gupta-Bleuler quantisation [8].

We have obtained our second result:

**Theorem 2** The two point supersymmetric positive invariant functions (distributions) are precisely the elements of the positive cone generated by $P_c \Delta G_c, P_a \Delta G_a$ and $- P_T \Delta G_T$.

We are ready now to present in the next section the Källen-Lehmann representation for free or interacting supersymmetric quantum field theory.

4 Källen-Lehmann representation

Let $\Phi(z)$ be a (scalar possible complex) supersymmetric quantum field. It is understood that it acts in a Hilbert space of supersymmetric states with
a vacuum $\Omega$. In order to cope with the singularities of $\Phi(z)$ in $x$-space we perform the smearing with a supersymmetric function $\varphi(z)$ obtaining

$$
\Phi(\varphi) = \int d^8z \Phi(z) \varphi(z)
$$

(4.1)

Denoting by $(.,.)$ the Hilbert space scalar product the two-point function of the quantum field $\Phi(z)$ is defined to be

$$(\Omega, \Phi^+(z_1)\Phi(z_2)\Omega)$$

where $\Phi^+$ is the Hilbert space adjoint of $\Phi$. It follows that, as in the usual quantum field theory too, the two-point function $(\Omega, \Phi^+(z_1)\Phi(z_2)\Omega)$ has to be positive definite (in the distribution theoretic sense). By making contact to the result in section 3 we obtain

**Theorem 3 (Källen-Lehmenn supersymmetric representation)** The two point function (distribution) in a supersymmetric quantum field theory has the form

$$(\Omega, \Phi^+(z_1)\Phi(z_2)\Omega) = \lambda_c P_c \Delta G_c + \lambda_a P_a \Delta G_a - \lambda_T P_T \Delta G_T$$

(4.2)

where $\lambda_i$ are positive constants and $G_i$ have (spectral) representations of the form

$$G_i = G_i(x_1 - x_2) = \int_0^\infty d\mu_i(m^2) \Delta_+(x_1 - x_2)$$

(4.3)

for $i = c, a, T$. If the spectral measures coincide then the representation reads

$$(\Omega, \Phi^+(z_1)\Phi(z_2)\Omega) = (\lambda_c P_c + \lambda_a P_a - \lambda_T P_T) \Delta G$$

(4.4)

where

$$G = G(x_1 - x_2) = \int_0^\infty d\mu(m^2) \Delta_+(x_1 - x_2)$$

(4.5)

We remind the reader that $\Delta = \delta(\theta_1 - \theta_2)\delta(\bar{\theta}_1 - \bar{\theta}_2)$ and $\Delta_+ = \Delta_+(x_1 - x_2)$ is the faminiar Pauli-Jordan function.

Note that the above result is valied for real or complex one component fields.
(an example is the free supersymmetric (real) massive vector field in the unitary gauge; for the arbitrary gauge fixing therm there is still a representation of type (4.2) but because the lack of positivity, the constants $\lambda_i$ are not all of them positive [4]). If a quantum field theory has several fields (like for instance the Wess-Zumino model, which involves chiral as well as antichiral fields, or fields with a Lorentz spin) the result may be different. For instance the positive definite kernel which appear in the Källen-Lehmann representation of the Wess-Zumino model is a two by two matrix kernel with nonvanishing entries (see [7]).

5 Conclusions

We obtained the most general form of supersymmetric invariants in two variables. Combining this result with supersymmetric positivity we studied the supersymmetric Källen-Lehmann representation. Starting with our results one can construct the generalized supersymmetric field by applying the same construction as usual [9].

Acknowledgement. We thank M. Schork for discussions.

References