Halo Studies for the ESS and Linac4 Front-Ends

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Abstract. In high intensity proton linacs, beam halo can be observed as early as the very first stages of acceleration. The initial population of halo particles found in the output distribution of the RFQ, increases in chopper lines, which interrupt the regular focusing pattern of RFQ and DTL. The resulting emittance growth and halo development contributes to beam loss in the low-energy accelerator sections and partly defines the beam quality in the high-energy sections. Using the examples of the ESS and Linac4 (CERN) front ends (RFQ to 20 MeV) we investigate emittance growth and beam halo due to statistically distributed field and gradient errors.

INTRODUCTION

The analysis of statistical errors represents the final stage in every linac design. Emittance growth, phase & energy jitter and beam loss due to these errors define the tolerances in the machine and thus have a direct impact on the project costs.

This paper reports on the progress in simulating statistical errors in the ESS [1] and Linac4 [2] front-ends, both high-intensity H^- linacs with a beam chopper in their Medium Energy Beam Transport (MEBT) line. The simulation approach using IMPACT [3] is outlined, and results from the simulation of quadrupole gradient errors as well as RF phase and gradient errors are reported. Finally the results are interpreted with respect to the current understanding of halo development.

SIMULATION

Lattices. Both front-ends start with a simulated RFQ output distribution [4] followed by a MEBT containing a beam chopper with several focusing elements. Two Alvarez Drift Tube Linac (DTL) tanks then raise the energy from 2.5 to 20 MeV in the ESS case and from 3.0 to 25 MeV for Linac4. Long term effects will be visualized using the full Linac4 lattice up to an energy of 120 MeV. The lattice parameters are summarized in Table 1.

Conventions. In order to quantify the amount of halo in a distribution snapshot, the fraction of particles outside n times its r.m.s. emittance (in each plane) is plotted. Thus the development of beam halo can be clearly separated from $\epsilon_{r.m.s.}$ growth and/or density oscillations in

TABLE 1. Lattice parameters

	ESS	Linac4
RFQ output	2.5 MeV	3.0 MeV
peak current	57 mA	30 mA
RF frequency	280 MHz	352 MHz
MEBT	3.11 m	3.73 m
no. of chopper plates	4	2
no. of buncher cav.	6	3
no. of quadrupoles	13	11
DTL	11 m	9 m + 6.1 m*
no. of RF sources	2	2+1
no. of gaps	77	82+29
no. of quadrupoles	78	83+29
output energy	20.3 MeV	24.9 MeV
CCDTL[†]		47.6 m
no. of RF sources		10
no. of 3/4-gap tanks		37
no. of quadrupoles		37
output energy		120 MeV

* + refers to the full Linac4

[†] for full Linac4

real space. An example is shown in Fig. 1 where the two input distributions are plotted. Both distributions carry particles with core radii of up to 5 times the r.m.s. value. A uniform error distribution is used for all simulations, meaning that the r.m.s. errors can be obtained by multiplying the quoted max. error amplitudes by ≈ 0.58 .

Simulation details. A preprocessor creates m random error sets, applies them to the original lattice and stores the new input files in scratch directories. With a simple shell script the jobs are then submitted to a Linux

FIGURE 1. Transverse halo: particles outside $n \times \epsilon_{r.m.s.}$ in the input distributions for the ESS and Linac4 front-ends.

cluster. For the full Linac4, 500 error sets with 50000 particles are simulated for each error amplitude, using ≈ 14 hours of CPU time on 30 processors. For the shorter front-end simulations 300 error sets are created. The process is then repeated for each error amplitude and evaluated by a post-processor module.

STATIC & DYNAMIC ERRORS

In the following we use the classic distinction between ‘dynamic’ errors that change from pulse to pulse or within single RF pulses and ‘static’ errors that change very slowly (seasonal changes) or remain completely unchanged during operation [5], [6]. Static errors are generally a residue of the initial adjustment of the lattice elements, e.g. random (gap to gap) RF amplitude errors from the field adjustment with bead pull ($\approx 1\%$), or random quadrupole gradient errors ($\approx 1\%$). The same category applies for mismatch and alignment errors which are not covered in this paper. While the static errors usually change from element to element the dynamic errors tend to be grouped, originating from pulse to pulse variations of RF ($< 1\%$, 1°) or quadrupole ($< 0.5\%$) power supplies.

Figure 2 shows the average $\epsilon_{r.m.s.}$ growth rates with respect to maximum transverse and longitudinal error amplitudes for grouped or completely random errors. In the transverse plane the grouped and ungrouped errors result in the same average $\epsilon_{r.m.s.}$ growth rates, cutting in half the number of simulations needed to evaluate both error sources. In the longitudinal plane, however, the grouped errors yield larger $\epsilon_{r.m.s.}$ growth (and larger phase & energy jitter) than the ungrouped errors. Although modern RF power sources with temperature controlled cables provide a margin of 0.5% field error and 0.5° phase error during the pulse, the error margins for the first bunches of a pulse may be considerably higher [5] and thus become a major concern.

FIGURE 2. Upper graph: average $\epsilon_{r.m.s.}$ increase for dynamic (grouped) and static (random) quadrupole gradient errors in the full Linac4; lower graph: dynamic and static RF field and phase errors

The average $\epsilon_{r.m.s.}$ growth for realistic dynamic and static error margins seems to be very small and therefore of no concern for low-loss machine operation. Nevertheless, a combination of all static and dynamic errors including alignment errors and mismatch may easily yield unacceptable losses and has to be evaluated for each particular machine.

In both planes, the maximum $\epsilon_{r.m.s.}$ growth rates for the worst cases are ≈ 3 to 4 times higher than for the plotted average cases. For dynamic errors the worst cases are of little importance since normal machine operation will automatically yield the average. The same reasoning applies for certain static errors (e.g. quadrupole gradient errors) which can be reduced with automated tuning systems. Some static errors, however, cannot be compensated (e.g. residual field adjustment errors in multi-gap RF tanks) and for those the worst case results should be taken into account.

PHASE & ENERGY JITTER

The dynamic errors in the longitudinal plane not only result in $\epsilon_{r.m.s.}$ growth but also in phase and energy deviation (jitter) from the nominal trajectory. If the linac beam is injected into a subsequent accelerator or a transport line with RF cavities, the phase and energy jitter has

to be limited to fit into the RF bucket of the following system.

We find that the largest phase & energy deviations usually originate from transitions between sections that are powered by different amplifiers and/or have different focusing structures. Opposite error amplitudes at these transitions can provide large kicks which are then amplified in subsequent drifts. An example is shown in Fig. 3, where the four worst case energy deviations are plotted for the two front-ends.

FIGURE 3. The four worst case energy deviations for a maximum of 1% field error and 1° phase error in the ESS (upper graph) and Linac4 (lower graph) front-end.

The average and maximum phase and energy jitter at the end of Linac4 are between 15 and 30% higher than for the ESS. This can probably be explained by the more compact design of the ESS MEBT (see Table 1) and the ‘missing gap’ at the transition of DTL1 and DTL2 in Linac4. In the ESS front-end, the two DTL tanks have a ‘seamless’ transition.

Considering Fig. 2 and 3 it is obvious that the partitioning of RF power supplies and cavities plays a crucial role in the development of phase & energy jitter. In the following example the jitter at the end of Linac4 is plotted for its original partitioning of power supplies and RF cavities, and for a partitioning with a reduced number of amplifiers (5 instead of 16, Fig. 4).

Contrary to the results with grouped and ungrouped RF errors in Fig. 2, we find that the lower number of power supplies results in a much lower phase & energy jitter and a much reduced average longitudinal $\epsilon_{r.m.s.}$ growth (2.6% instead of 5.8%). This suggests that the power

FIGURE 4. Phase & energy jitter at the Linac4 output for a maximum of 1% field error and 1° phase error using 16 (upper graph) or 5 (lower graph) RF amplifiers.

splitting should be carefully optimized for each machine and that the optimum number of RF gaps per power source may be different for different lattice types and different energies.

HALO FROM STATISTICAL ERRORS?

Parametric beam halo develops for integer ratios (usually 2:1) between the oscillations of a mismatched beam core and the oscillations of single particles. This type of coherent core oscillation can be excited by initial mismatch and can remain remarkably stable throughout an entire linac. Figure 5 shows one such example for a fast-mode

FIGURE 5. 30% fast mode excitation in the superconducting section of the SPL.

excitation in the 600 m long superconducting section of the SPL [7]. Energy transfer from the core oscillations to the single particle orbits eventually damps the oscillations of the core and thus creates halo.

Similar oscillations, though much more irregular, can be observed in simulations with statistical errors, e.g. in Fig. 6 where the radial deviations for an average case of quadrupole errors (max. $\pm 1\%$) reach amplitudes between 5 and 10%. For the worst case they increase to 40%, reaching similar levels to those generated by strong initial mismatch. Figure 6 also shows that the worst ra-

FIGURE 6. Upper graph: average case of radial deviations for $\pm 1\%$ quadrupole gradient errors in the ESS front-end ($\approx 5\% \epsilon_{r.m.s.}$ growth), lower graph: worst case for Linac4 ($\approx 100\% \epsilon_{r.m.s.}$ growth).

dial deviations occur after the MEBT, in the first two DTL tanks. Unless these tanks are designed with large bore radii they are likely to suffer from high beam loss. The corresponding beam halo for the ESS front-end (Fig. 7) shows an almost invisible increase for the average case and roughly a doubling of large amplitude particles for the worst case. For the full Linac4, Fig. 8 shows a more visible effect for the average error case. Despite the difficult interpretation of these curves (beam halo already for the matched case, particle loss, low number of particles ≤ 50000 , etc.) it seems as if statistical errors might have the potential to create beam halo. Fig-

FIGURE 7. Transverse halo due to ungrouped 1% (max.) quadrupole gradient errors in the ESS front-end.

ures 7 and 8 both show an increased number of particles beyond the core emittance ($> 5 \cdot \epsilon_{r.m.s.}$) even though the amount of the observed halo is less dramatic than for initial mismatch.

FIGURE 8. Transverse halo due to ungrouped 1% (max.) quadrupole gradient errors in Linac4.

SUMMARY AND OUTLOOK

First results of IMPACT simulations of dynamic and static gradient & phase errors were presented for Linac4 and the ESS front-end. Quadrupole gradient variations for grouped and ungrouped errors resulted in equal $\epsilon_{r.m.s.}$ growth rates and seem to be of little concern. Longitudinally the partitioning of power sources and RF gaps had a significant influence on the resulting $\epsilon_{r.m.s.}$ growth and energy & phase jitter. Additional concerns are the possibility of an unfortunate combination of residual field adjustment errors, and the enhancement of RF error bars for the first bunches of an RF pulse. The simulations indicate the possibility of halo development due to statistical errors, although more systematic work is needed to understand the process fully. A combination of all error sources and an evaluation of alignment errors has yet to be done and might add significance to the effects of single error sources.

REFERENCES

- Gerigk, F., Revised ESS Front-End (2.5 - 20 MeV), Tech.Rep. (2003).
- Gerigk, F., and Vretenar, M., "Design of a 120 MeV H^- Linac for CERN High-Intensity Applications," in *Proceedings LINAC02*, Kyongju, Korea, 2002.
- Qiang, J., Ryne, R. D., Habib, S., and Decyk, V., *Journal of Computational Physics*, **163**, 1-18 (2000).
- The ESS Project Volume III, Technical Report*, ISBN 3-89336-303-3, 2002.
- Vretenar, M., private communications (2003).
- Findlay, D., private communications (2003).
- Vretenar, M., editor, *Conceptual Design of the SPL, a High-Power Superconducting H^- Linac at CERN*, CERN 2000-012, 2000.