Morphology of Galactic Open Clusters

W. P. Chen

Institute of Astronomy and Department of Physics, National Central University, Chung-Li 32054, Taiwan

wchen@astro.ncu.edu.tw

C. W. Chen

Institute of Astronomy, National Central University, Chung-Li 32054, Taiwan

awei@outflows.astro.ncu.edu.tw

and

C. G. Shu

Shanghai Astronomical Observatory, Chinese Academy of Sciences, Shanghai 200030, and Center for Astrophysics, Shanghai Normal University, Shanghai 200234

cgshu@center.shao.ac.cn

ABSTRACT

We analyzed the shapes of Galactic open clusters by the star counting technique with the 2MASS star catalog database. Morphological parameters such as the ellipticity and size have been derived via stellar density distribution, weighed by clustering probability. We find that most star clusters are elongated, even for the youngest star clusters of a few million years old, which are located near to the Galactic disk. The shapes of young star clusters must reflect the conditions in the parental molecular clouds and during the cluster formation process. As an open cluster ages, stellar dynamics cause the inner part of the cluster to circularize, but the overall radius gets larger and the stellar density becomes sparser. We discuss how internal relaxation process competes with Galactic external perturbation during cluster evolution.

Subject headings: stellar dynamics; methods: data analysis; galaxies: star clusters

1. Morphology of Star Clusters

The way member stars distribute within a star cluster changes as the cluster evolves. The initial stellar distribution is governed by the structure of the parental molecular cloud, and by how star
formation proceeds. Sequential star formation, for example, may result in massive stars, which are responsible for inducing the formation of next-generation stars, to have a birth place markedly different from that of low-mass stars. As the star cluster evolves, the distribution is modified by internal gravitational interaction among member stars. Subsequently stellar evaporation and external disturbances—e.g., Galactic tidal force, differential rotation, and encounters with molecular clouds—would alter the spatial structure and eventually dissolve the cluster. While individual star clusters are good laboratories to study stellar dynamics (Friel 1995), they also serve as test particles to probe the local physical (e.g., gravitational field) and chemical (e.g., metallicity) conditions. Open clusters in particular, with their wide ranges of age and location distribution, would be valuable tools to study the star formation history and chemical evolution of the Galactic disk.

Stars in a globular cluster are known to concentrate progressively toward the center, more so for massive stars than for low-mass stars. The density distribution, prescribed by the King model (King 1966), is understood as a combination of an isothermal sphere (i.e., dynamically relaxed) in the inner part of the cluster, and tidal truncation by the Milky Way in the outer part. In contrast, open clusters appear irregularly shaped, with member stars sparsely distributed. The youngest clusters must still bear the imprint of their formation history, so their structure, when compared with that of molecular clouds, may shed crucial light to the fragmentation process during cloud collapse. At the other end, the oldest open clusters serve as tracers of the structural and evolutionary history of the Galactic disk (Friel 1995). Even though some open clusters are known to follow the King model (King 1962), and despite some elaborative theoretical considerations on the dynamics of open clusters (e.g., Portegies Zwart et al. 2001)), few observational studies have been done systematically until recently (e.g., Pandey, Mahra, & Sagar 1990; Nilakshi et al. 2002), perhaps due to the complexity of the problem arising from the small number of member stars, and the lack of comprehensive data of open clusters on large angular scales.

As early as almost a century ago, Jeans (1916) already considered the flattening effect of a moving cluster through a general gravitation field of stars. Oort (1979) noticed the uneven distribution of stars in Hyades by comparison of the numbers of stars in quadrants about the cluster’s apparent center. He concluded that the cluster appeared flattened with an axial ratio of about 2 aligned with the Galactic plane. Bergond, Leon, & Guibert (2001) analyzed the surface star density in 3 open clusters by the star counting method and found elongated shapes parallel to the Galactic plane. These authors furthermore employed wavelet transform to bring out possible tidal tails, presumably caused by the Galactic tidal field.

A comprehensive diagnosis of the spatial structure of open clusters appears elusive because they are loose systems with shallow gravitational potential, hence lack of organized symmetry in structure, and likewise the morphology and shape would be vulnerable to external perturbation. We would like to learn whether young clusters are mass segregated, and if so, to what extent this is due to dynamical relaxation, as opposed to relic structure in molecular clouds. To answer questions like these, it is desirable to study the spatial structure of the youngest star clusters, and
see how it evolves as a cluster ages and moves in its Galactic orbit.

As a pilot study on how a star cluster shapes out of the molecular cloud, we (Chen & Chen 2002) investigated the radial star density profiles of 7 open clusters with ages ranging from a few million years to a few billion years, based on the 2MASS (Two-Micron All-Sky Survey) Second Increment Release star catalog. The infrared data enabled us to probe the distribution of member stars of clusters that are very young and still embedded in molecular clouds. Our study indicated that stars, regardless of their masses, tend to concentrate progressively toward the center of a cluster, and the degree of concentration is higher for luminous (presumably more massive, or as binary) stars than for fainter members. Such a segregation structure appears to exist in even the youngest star clusters. The relaxation time \( \tau_{\text{relax}} \approx (0.1N/\ln N) \cdot \tau_{\text{cross}} \), where \( N \) is the number of stars and \( \tau_{\text{cross}} \) is the crossing time of the system (Binney & Tremaine 1987). For a typical open cluster \( N \sim 10^3 \), and \( \tau_{\text{cross}} = R/V \) where the size of the cluster \( R \sim 2 \) pc and the velocity dispersion \( V \sim 1 \) km s\(^{-1}\), yielding \( \tau_{\text{relax}} \approx 3 \times 10^7 \) yrs. The youngest open clusters (ages of a few million years) therefore have not had time for dynamical relaxation to take place efficiently. The spatial distribution of member stars in the youngest star clusters hence is much relevant to the structure in the parental cloud out of which the cluster was formed, and to subsequent mass redistribution during star (cluster) formation process.

Nilakshi et al. (2002) analyzed the radial star density profiles of some 38 open clusters based on the U.S. Naval Observatory (USNO)-A v2.0 star catalog, which in turns was derived from the Digital Sky Survey (DSS) plates. Because no membership information is available, the extent, size, shape, or any spatial structure of a particular star cluster is estimated in a statistical sense against adjacent Galactic star fields. One of the advantages of working with a homogeneous sky-survey database, such as the 2MASS or USNO catalog, in addition to the convenient availability, is the extended angular coverage to encompass not only the entirety of a cluster itself (core, envelope, or possible tails), but also sufficient field regions for comparison. A fair assessment of the stellar density fluctuations in comparison fields is crucial in the star counting technique.

Parameterization of the spatial structure of an open cluster by its radial distribution alone obviously is not adequate, as already pointed out by Nilakshi et al. (2002). Many open clusters have highly irregular shapes, often even with no clear centers, so circular symmetry cannot be readily assumed. Dense molecular cloud cores are shaped on average as a prolate Gaussian with an intrinsic axial ratio of 0.54, as a part of an evolutionary sequence from filamentary molecular cloud complexes to roughly spherically condensed cores (Curry 2002). It is therefore desirable to represent the stellar distribution by a more sophisticated method than a one-dimensional analysis. This paper summarizes the result of our attempt to analyze the morphological shapes of open clusters, and how the shaping would evolve in Galactic environments.

We take a probabilistic approach to estimate the stellar surface density of stars in the 2MASS point source catalog, and represent a star cluster with an ellipsoid. This allows us to investigate not only the structure (concentration, segregation) but also the shape (elongation, orientation) of
a cluster, and hence by a sample of clusters of different ages and environments, to delineate the morphological evolution influenced by Galactic dynamics. The 2MASS infrared data are free from much of interstellar extinction, and so would reveal the true shape of a star cluster more readily than in optical wavelengths. Some youngest star clusters may be seen only in infrared wavelengths (Lada & Lada 2003). We describe our 2MASS sample of open clusters in Section 2, and present the methodology to analyze the morphology of open clusters in Section 3. The results and discussions are summarized as the last section.

2. Data of Open Clusters

Dias et al. (2002) compiled the latest catalog of open clusters, which is based on the previous work of Lyngå (1987) and of Mermilliod¹, with some updated data on radial velocity, proper motion and metallicity. Of the total of more than 1600 entries, about 38% have distance, age and color-excess determinations. We note that while such a compilation is useful for information retrieval, one should exercise caution when deriving statistics from the dataset, because the catalog is far from completion—some of the entries may not be bona fide stellar groups at all, and perhaps a lot more open clusters are yet to be discovered. Severe extinction by dust near the Galactic plane makes it difficult at wavelengths shortward of infrared to recover the true shapes, or even their bare existence, of young star clusters, which as recent studies show may outnumber optically visible open clusters by a factor of an order (Lada & Lada 2003).

We selected among the first 800 entries in the Dias et al. (2002) catalog open clusters, roughly from RA 0ʰ to 12ʰ which suffer less extinction and source crowdedness toward the Galactic center, have distance and age determinations available and have angular sizes between 3′ and 40′. The choice of the angular range, somewhat arbitrary, is a convenient compromise between spatial resolution and the practical limit of maximal data (1 degree field of view) downloadable from the 2MASS web interface. We selected those clusters that are rich in density enhancement (by eye) and with as complete data coverage in the 2MASS star catalog (All-Sky Data Release) as possible, i.e., no nearby bright stars so as to contaminate the field. In addition to bright field stars, young open clusters often contain hot, luminous members. A real bright star would leave a blank pattern, rendering an incomplete listing in the 2MASS database. Even a moderately bright star would cause unreliable astrometric and photometric determinations on neighboring stars. By working with the 2MASS data, our sample suffers less brightness contrast between the hottest stars and faint members than in the visible wavelengths. Interpolation of stellar density is possible in most cases as long as the contamination is not overwhelming (e.g., too bright or too close within the cluster boundary). In the study we report here, none of the sample suffers bright-star contamination, and a total of 31 open clusters were selected in the morphological analysis.

¹http://obswww.unige.ch/webda
Our sample has no obvious additional selection effects, other than to avoid the Galactic center where the open cluster catalog (Dias et al. 2002) itself may be highly incomplete. Figure 1 shows the pole-on and edge-on Galactic distributions of our star cluster sample. Clusters younger and older than 800 Myrs are marked differently. It is noted that the majority (29) of our sample are in the direction of Galactic anti-center, a consequence of our sample selection from half of the Dias et al. (2002) catalog. It is seen that old open clusters have a larger average scale height above the Galactic plane than young star clusters do.

3. Methodology

3.1. Statistical Membership

In the absence of membership information on individual stars, we estimate the structure of a star cluster by a probabilistic star counting technique, i.e., the boundary, shape, size are all determined in a statistical sense. In essence, the degree of clustering of neighbors around any star gives a measure of the likelihood of that star being in a cluster environment. One defines for each individual star \( i \) the clustering parameter,

\[
P_i = \frac{N_t - N_f}{N_t} = 1 - \frac{N_f}{N_t},
\]

for which \( N_t \) is the total number of neighboring stars within a specified angular size (a “neighborhood aperture” centering on the \( i \)th star) and \( N_f \) is the average number of field stars within the same aperture. The number of field stars can be estimated by \( N_f = \Sigma_f \times \pi r_p^2 \) where \( \Sigma_f \) is the surface number density of field stars, estimated in regions away from the apparent star cluster and \( r_p \) is the radius of the neighborhood aperture.

We see that the clustering parameter \( P_i \) gives a measure of local enhancement of stellar density, whose value ranges from \( P_i \sim 0 \) in a field region (for which \( N_t \sim N_f \)) to \( P_i \sim 1 \) near a rich cluster (\( N_t \gg N_f \)). In other words, \( P_i \) behaves very much like a probability of cluster membership. Our probabilistic method to analyze membership in a star cluster is similar to that used by Danilov & Seleznev (1990).

Note that the choice of \( r_p \) should not be arbitrary. If \( r_p \) is too small, the uncertainty in \( P_i \) will be large due to small-number statistics. On the other hand, if \( r_p \) is too large, the intrinsic structure of the cluster will be smoothed out and detailed structure information is no longer available.

We describe below our procedure to select an appropriate aperture size \( r_p \) from the surface density of field stars \( \Sigma_f \). Once both \( r_p \) and \( \Sigma_f \) for a cluster are determined, we can estimate the membership probability for each star by Eq. 1.
3.2. Surface Density of Field Stars

The cluster membership probability defined above is based on a measure of enhancement of local stellar density compared to field stars. A fair estimate of the density of field stars therefore is prerequisite. For each cluster, we use regions away ($\gtrsim 20'$) from the central cluster to estimate the mean number density of field stars.

Since our analysis of the cluster parameters is based on the statistics (star density, aperture size) in the surrounding field regions, obviously a homogeneous distribution of field stars must be assumed in order for our technique to work. We note actually in almost every case we encountered, the distribution of field stars is not homogeneous, due to the general stellar gradient vertical to the Galactic disk. Figure 2 shows such an example for which a density gradient is seen in both the RA and the DEC projections, with the overall gradient vector pointing toward the Galactic disk, as expected. We also find that, not surprisingly, the density gradient is higher for a line of sight with a lower Galactic latitude. In our analysis we removed a flat surface density as background for individual star clusters, but empirically found this kind of density gradient—though potentially useful for study of Galactic disk stellar populations—to have little effect on our morphological results because the density differs no more than a few percents across the fields of our star clusters.

The surface number density of field stars $\Sigma_f$ is computed by counting the number of stars in a certain sky area. Due to the discrete nature of individual stars, even for a uniform star field, $\Sigma_f$ would approach a constant only when the sampling sky area is large enough to include a sufficient number of stars. Otherwise, when the sampling size is smaller than about the average angular separation of stars, large fluctuations result. As the minimum, one would demand a sky area to have a signal-to-noise of $\sqrt{N_f} \gtrsim 3$ against Poisson fluctuations, or $N_f \gtrsim 10$ field stars to determine accordingly the optimal neighborhood aperture size $r_p$. In our analysis we take $N_f = 50$ (i.e., $S/N \sim 7$) and select the corresponding $r_p$ for each cluster. In general $r_p$ is on order of a few arcminutes.

3.3. Cluster Shape and Morphology

With each star now being represented by a membership probability, the surface density of cluster member stars is then the sum of the clustering parameter of every star within each sky-coordinate grid (with area $\Delta S$); that is, the effective number density of member stars is,

$$\Sigma_s = \sum P_i/\Delta S$$

(2)

Obviously in field regions, $\Sigma_s \to 0$. The morphology of a cluster is prescribed by density contours, at both the one-third level of the maximal density ($\Sigma_{\text{max}}/3$), and at the boundary, defined as where the density drops to 2–3 times the background fluctuations. In the extreme case, e.g., in
NGC 2567, for which the star cluster is very sparse, a mere 1-σ outer boundary was used. Effectively the inner and outer ellipses trace, respectively, the core and the halo (or the corona as termed by some researchers) of a star cluster. The density contours are least-squares fitted with ellipses to obtain the eccentricities of the inner and outer ellipses, and the corresponding sizes (average of the semimajor and semiminor axes). We use the flattening parameter to quantify the shape of an ellipse. The flattening parameter (or oblateness\(^2\)) is defined as \( f = 1 - (b/a) \), where \( a \) and \( b \) are respectively the semimajor and semiminor axes of an ellipse, so \( f \) is related to the eccentricity \( e \) by \( e^2 = 1 - (b/a)^2 = 1 - (1 - f)^2 \). For \( b/a \) very close to 0 (highly flattened) or 1 (our case), the flattening parameter \( f \) is more discriminative than the eccentricity itself.

The uncertainties in the determination of the flattening are estimated by Monte Carlo simulations of star clusters of different shapes. For typical parameters of the open clusters in our sample, i.e., with \( \sim 3 \) times enrichment of stellar number density of cluster members with respect to the field, the uncertainties in the flattening for the outer boundary and for the core (1/3 maximum) are \( \delta f_{\text{out}} \sim 0.18 \) and \( \delta f_{1/3} \sim 0.12 \) for a spherical cluster (i.e., \( b/a = 1 \)), and are \( \delta f_{\text{out}} \sim 0.12 \) and \( \delta f_{1/3} \sim 0.06 \) for an elongated cluster with \( b/a = 0.5 \). Obviously, the richer the cluster, the smaller the uncertainties. For instance, for a spherical globular cluster with 100 times density enhancement, \( \delta f_{\text{out}} \sim 0.06 \), and \( \delta f_{1/3} \sim 0.02 \).

In addition to the shape (flattening) and size, other parameters, such as the position angle of the ellipse, can also be obtained. We illustrate in Fig. 3 examples of two clusters with different morphology, one of relatively round shape (NGC 2414) and the other of elongated shape (NGC 1893). These two clusters will be compared in details in the next section.

We list the results of individual clusters in our sample in Table 1, where the first column gives the name of the cluster, the second column is the Galactic coordinates in degrees, and the next three columns are the heliocentric distance, height from the Galactic disk, and the age of each star cluster, with the distance and age taken from Dias et al. (2002). Columns 6 to 11 list, respectively, the flattening parameters and corresponding sizes (average of semimajor and semiminor axes) derived from the density contours of member stars. The last column gives the total number of member stars, by summing the total membership probability of a cluster.

4. Morphological Evolution of Open Clusters

Two kinds of dynamical effects act on, and influence the morphology of, an open cluster. The first is internal interaction of two-body relaxation due to encounters among member stars. This leads to stellar distribution in spherical shape, ever denser toward the cluster center. At the same time, low-mass members may gain enough kinetic energy through the encounters and get thrown out of the system (i.e., stellar “evaporation”). The other dynamical process is external interaction,
including tidal force due to Galactic disk or giant molecular clouds, and differential rotation especially for a cluster located in the inner disk region. These disruptive effects tend to make a star cluster elongated in shape, with the outer parts particularly vulnerable.

Tidal disturbance on a star cluster is stronger when closer to the Galactic disk plane. Fig. 5 shows the vertical heights from the Galactic disk versus the ages of the open clusters in the Dias et al. (2002) catalog, of which our sample, separately marked, is a subset. As can be seen, young open clusters tend to reside close to the disk where molecular clouds, from which the star clusters were formed, are distributed (Tadross et al. 2002; Chen, Hou, & Wang 2003). If open clusters are separated into two age groups, old and young, with the dividing age, somewhat arbitrary, of that of the Hyades, 0.8 Gyr (Phelps et al. 1994; Chen, et al. 2003), the scale heights are 354 pc (old) and 57 pc (young), respectively (Chen, et al. 2003), based on the data compiled by Dias et al. (2002)). Tadross (2003) included additional consideration of the galacto-distances of the clusters and the results support the previous assertion (Lyngå & Palous 1987) that old clusters seem to distribute at larger scale-heights in the outer parts of the Galaxy than in the inner parts. Obviously, only star clusters away from the inner disk regions—where the tidal force from giant molecular clouds plays a major disruptive role in the structure or even the existence of a star cluster—would have survived on Galactic dynamical time scales (Janes & Phelps 1994).

Comparison between the two open clusters in Fig. 3 is informative. NGC 2414 is relatively poorly studied, perhaps because of its paucity and small angular extent. Much of the literature about this open cluster can be traced back to the photometric measurements in Vogt & Moffat (1972), based on 10 member stars. In comparison, NGC 1893 is prominently stretched toward the disk plane. The cluster is associated with bright nebulosity and dark clouds, but does not seem to show positional variation of color excess $E(B-V)$ across the cluster (Yadav & Sagar 2001). The extinction effect would certainly be even smaller in the 2MASS 2 μm data. Its elongated shape (or two subgroups) thus should be inherent to the stellar distribution within the cluster, rather than caused by extinction variation.

These two clusters share some similarities, namely both being relatively young, log (age/yr) $\sim$ 7, and close to the disk plane, with NGC 2414 at $\ell$ $\sim$ 231°, $b$ $\sim$ +2°, and NGC 1893 at $\ell$ $\sim$ 174°, $b$ $\sim$ −2°. They however contrast greatly in shape, size, and apparent richness. NGC 2414 is round (flattering $\sim$ 0.1 throughout the entire cluster), small, and contains some 74 member stars within its derived radius $\sim$ 3′. NGC 1893 on the other hand is oval, twice as large in angular extent, and encompasses 645 member stars. NGC 1893 therefore has a much stronger gravitational binding than NGC 2414 against external disruption. What we see now in NGC 2414 may well be its remnant cluster core.

Such a stripping off of cluster halos is not uncommon. For example, Berkeley 17 (=C0517+305), the oldest open cluster known (Salarism, Weiss, & Percival 2004), has a protrusion manifestly pointing toward the Galactic disk (Fig. 6), with a projected extent comparable to the cluster’s radius, $\sim$ 6–7′, or about 5 pc in projected length assuming a 2.7 kpc distance (Dias et al. 2002).
The enhancement ("corona") has already been inferred from radial stellar density distribution, and from the color-magnitude diagram of stars away from the nominal cluster region (Kahúñy 1994). Our analysis brings up clearly the tail and its geometry, and further hints on an associated antitail, which typifies tidal distortion.

The old, metal-poor, and large-height cluster NGC 2420 (Friel et al. 2002), serves as a good showcase for interplay between internal stellar dynamics and external disturbances. This cluster is nominally listed in Dias et al. (2002) to have an angular diameter of 5′, at a distance of 3085 pc, and a logarithmic age of 9.048 years, apparently taken from WEBDA. We adopt this distance and those of other star clusters in our studies from the Dias et al. (2002) catalog as a homogeneous source of input data, but note that this cluster may be considerably closer, at 2.28 kpc (Janes & Phelps 1994). Errors in distance estimation—which often turn out to be quite significant among open clusters—would not affect the shape determination, but obviously would influence the results of any statistical analysis. Using star counting on a Palomar Sky Survey plate, Leonard (1988) already noticed a much larger extent, to at least 20′, from the apparent center of NGC 2420. Our analysis of the 2MASS data shows a negligible gradient in the field star density toward NGC 2420 (Fig. 4). The cluster itself is determined to have an inner (1/3 maximum) and outer (3-σ sky) ellipses with flattening of $f_{1/3} \sim 0.06$, and $f_{\text{out}} \sim 0.12$, respectively. Average sizes of $r_{1/3} \sim 3.7′$ and $r_{\text{out}} \sim 11.6′$ were obtained. For NGC 2420, an aperture size $r_p = 3.6′$ has been used to estimate the field star density, $\Sigma_f \sim 1.2 \pm 0.1$ arcmin$^{-2}$, from which the surface density of member stars is then derived according to Eq. 2. At the 2MASS limit of $K_s=15.6$ mag, by summing up all $P_i$s (Eq. 1), there are a total of 468±22 member stars. This is to be compared with 685±27 within a radius of 20′ up to the completeness of photographic $\sim 19.5$ mag studied by Leonard (1988), which covers a larger sky area with deeper stellar photometry than the 2MASS data we have used. We may be witnessing the disintegrating process of the outer part of NGC 2420.

Globular clusters are known to have elliptical shapes (White & Shawl 1987), likely due mainly to their overall rotation, rather than to Galactic tidal interactions (King 1961; Meyan & Mayor 1986). Kontizas et al. (1990) compared globular clusters in the Large and Small Magellanic Clouds, and found that in virtually all cases, the inner parts are more elliptical than the outer parts. This is understood as the diminishing effect of rotation from the inner to the outer parts of a globular cluster. These authors also found that the SMC globular clusters are more elliptical than those of the LMC, which in turn are more elliptical than those in our galaxy, and that the outer shapes are somewhat flatter for younger systems (Kontizas et al. 1991).

Open clusters are also recognized to have elongated shapes but the orientation of the flattening, however, cannot be accounted for by Galactic tides alone (Jefferys 1976). This may be because a giant molecular cloud, local to a particular open cluster, plays a more influential role in shaping the cluster than a general Galactic disk potential (Danilov 1996). Our analysis shows that even the youngest open clusters (several Myrs old) have very elongated shapes, with $f \sim 0.5$. This is to be compared with the inferred mean intrinsic axial ratio of 0.54 for the observed shapes of molecular cloud cores (Curry 2002). The youngest star clusters hence appear by and large to have
inherited the morphological shapes from the prenatal molecular clouds. Because of the low volume stellar number density of an open cluster, internal dynamics never becomes dominant so as to sphericalize the system (Portegies Zwart et al. 2001), in contrast to the case in globular clusters.

Fig. 7 and Fig. 8 summarize our analysis of the shapes of the open cluster sample. In Fig. 7 the shapes of the cluster core (diagnosed with $f_{1/3}$) and of the outer boundary ($f_{\text{out}}$) are shown as a function of age. One sees that most open clusters are elongated in shape. Because what we see is the projected shape, the clusters can only actually be flatter. The number of our sample of open clusters is small (31), yet it is tantalizing to note that while the shape of the outer boundary remains unconstrained with age, the core tends to circularize as a star cluster ages. If the sample is further divided to two groups according to the height from the disk plane $|H|$ by the median value of the sample $\sim 170$ pc, one sees clearly that the large-height group (denoted by open circles in Fig. 7) evinces a noticeable tendency of rapid spheridalization both in the core and in the halo. Low-lying clusters remain elongated from birth to date, prominently so for the outer halos, which are the most vulnerable to disruption. Fig. 8 shows how the flattening of star clusters varies with the height from the disk plane, for young (filled circles) and old (open circles) systems. The youngest star clusters close to the disk clearly manifest the most elongated shapes. The tendency toward spherical shaping away from the disk, in particular in the cores of old systems, demonstrates again the working of internal dynamical relaxation. The flattening shapes of young clusters are primarily inherited from cluster formation process, but after some $\sim 10^8$ yr or so, internal stellar dynamics become effective in shaping the core of a cluster.

Internal stellar dynamics affect not only the shape of a cluster, but also its size. Stellar evaporation results in a reduction of number of member stars, and the cluster responds by expansion in radius and hence a decrease in stellar density. This is clearly shown in Fig. 9 for which older clusters tend to be larger in size and less dense in stellar density. Young open clusters are born to be rich in members, amounting to more than $10^2$ stars per cubic parsec. Subsequent dynamical evolution apparently causes them to "loose up" and become dispersed, with the stellar number density dropping to that comparable to the field in the solar neighborhood, or even lower (for large-height clusters). Because old clusters on average have large scale heights, it is difficult to distinguish from our data unambiguously the distortion effect by aging from that by Galactic tides. Fig. 10 plots the same radius and stellar density of our sample versus the height. The stellar density remains more or less the same for clusters within the disk (height less than $\sim 100$ pc), yet diminishes markedly with height. The age seems to play a more definite role in the cluster dynamics than Galactic tidal force.

This paper presents our first attempt to delineate the structural evolution as a tool to probe the mass distribution and perhaps the dynamical evolution of our Galaxy. Our sample is too small for firm quantitative inference, but some preliminary conclusions can be drawn. Our study indicates that the shape or morphological structure of a young open cluster is dictated by the initial conditions in the parental molecular cloud and, as the cluster evolves, by both the internal gravitational interaction and external tidal perturbations. Only the initially massive and compact
star clusters would have strong enough self-gravitational binding to endure the continuing destructive effects, which intensify near the disk plane and toward the inner parts of the Galaxy. Statistics based on a larger sample of open clusters are obviously needed to quantify, e.g., the time scales, the intricate interplay between cluster evolution (age) and its Galactic location (galacto-centric distance, and height from the Galactic disk).

This work makes use of data products from the Two-Micron All-Sky Survey (2MASS), which is a joint project of the University of Massachusetts and the Infrared Processing and Analysis Center/California Institute of Technology, funded by the National Aeronautics and Space Administration and the National Science Foundation. We thank the referee for very constructive suggestions that greatly improve the quality of the paper. CGS expresses his gratitude for the hospitality during his visit at NCU. WPC and CWC acknowledges financial support of the grant NSC92-2112-M-008-048 from National Science Council of Taiwan.

REFERENCES

Danilov, V. M., 1996, IAUS, 174, 389
King, I., 1961 AJ, 66, 68
King, I., 1966, AJ, 71, 64
Tadross, A. L., 2003, New Astro., 8, 737

This preprint was prepared with the AAS L\TeX macros v5.2.
Table 1. Morphological Parameters of Open Clusters

<table>
<thead>
<tr>
<th>Name</th>
<th>(l, b) (°, °)</th>
<th>d (pc)</th>
<th>H (pc)</th>
<th>log age (yr)</th>
<th>f_{out}</th>
<th>f_{1/3}</th>
<th>r_{out} (arcmin)</th>
<th>r_{1/3} (arcmin)</th>
<th>N*</th>
</tr>
</thead>
<tbody>
<tr>
<td>NGC 2420</td>
<td>(198.11, +19.63)</td>
<td>3085</td>
<td>1036</td>
<td>9.05</td>
<td>0.12</td>
<td>0.06</td>
<td>11.60</td>
<td>3.72</td>
<td>468</td>
</tr>
<tr>
<td>NGC 2506</td>
<td>(230.56, +09.93)</td>
<td>3460</td>
<td>597</td>
<td>9.05</td>
<td>0.26</td>
<td>0.15</td>
<td>10.42</td>
<td>3.48</td>
<td>1091</td>
</tr>
<tr>
<td>NGC 1893</td>
<td>(173.59, -01.68)</td>
<td>3280</td>
<td>-96</td>
<td>6.75</td>
<td>0.32</td>
<td>0.51</td>
<td>6.13</td>
<td>4.11</td>
<td>645</td>
</tr>
<tr>
<td>King 5</td>
<td>(143.74, -04.27)</td>
<td>1900</td>
<td>-141</td>
<td>9.00</td>
<td>0.43</td>
<td>0.21</td>
<td>4.86</td>
<td>3.90</td>
<td>292</td>
</tr>
<tr>
<td>NGC 6791</td>
<td>(69.96, +10.90)</td>
<td>5853</td>
<td>1107</td>
<td>9.64</td>
<td>0.20</td>
<td>0.09</td>
<td>8.60</td>
<td>3.90</td>
<td>1180</td>
</tr>
<tr>
<td>Berkeley 17</td>
<td>(175.65, -03.65)</td>
<td>2700</td>
<td>-172</td>
<td>10.08</td>
<td>0.51</td>
<td>0.21</td>
<td>8.19</td>
<td>4.23</td>
<td>373</td>
</tr>
<tr>
<td>Melotte 71</td>
<td>(228.95, +04.50)</td>
<td>3154</td>
<td>247</td>
<td>8.37</td>
<td>0.51</td>
<td>0.12</td>
<td>10.13</td>
<td>3.83</td>
<td>659</td>
</tr>
<tr>
<td>NGC 1245</td>
<td>(146.65, -08.93)</td>
<td>2876</td>
<td>-446</td>
<td>8.70</td>
<td>0.24</td>
<td>0.02</td>
<td>10.21</td>
<td>4.40</td>
<td>629</td>
</tr>
<tr>
<td>Berkeley 69</td>
<td>(174.44, -01.79)</td>
<td>2860</td>
<td>-89</td>
<td>8.95</td>
<td>0.30</td>
<td>0.28</td>
<td>3.79</td>
<td>3.08</td>
<td>110</td>
</tr>
<tr>
<td>NGC 1960</td>
<td>(174.53, +01.07)</td>
<td>1318</td>
<td>25</td>
<td>7.47</td>
<td>0.32</td>
<td>0.23</td>
<td>7.60</td>
<td>5.65</td>
<td>607</td>
</tr>
<tr>
<td>King 8</td>
<td>(176.39, +03.12)</td>
<td>6403</td>
<td>348</td>
<td>8.62</td>
<td>0.22</td>
<td>0.28</td>
<td>2.54</td>
<td>1.69</td>
<td>190</td>
</tr>
<tr>
<td>Berkeley 21</td>
<td>(186.84, -02.51)</td>
<td>5000</td>
<td>-219</td>
<td>9.34</td>
<td>0.06</td>
<td>0.14</td>
<td>3.84</td>
<td>2.55</td>
<td>288</td>
</tr>
<tr>
<td>NGC 2414</td>
<td>(231.41, +01.95)</td>
<td>3455</td>
<td>118</td>
<td>6.98</td>
<td>0.11</td>
<td>0.11</td>
<td>3.14</td>
<td>2.81</td>
<td>74</td>
</tr>
<tr>
<td>Trumpler 7</td>
<td>(238.21, -03.33)</td>
<td>1474</td>
<td>-86</td>
<td>7.43</td>
<td>0.40</td>
<td>0.61</td>
<td>8.60</td>
<td>4.56</td>
<td>178</td>
</tr>
<tr>
<td>NGC 1907</td>
<td>(172.62, +00.31)</td>
<td>1556</td>
<td>8</td>
<td>8.57</td>
<td>0.12</td>
<td>0.24</td>
<td>4.19</td>
<td>3.33</td>
<td>55</td>
</tr>
<tr>
<td>NGC 2421</td>
<td>(236.27, +00.07)</td>
<td>2181</td>
<td>3</td>
<td>7.37</td>
<td>0.14</td>
<td>0.12</td>
<td>4.16</td>
<td>3.70</td>
<td>355</td>
</tr>
<tr>
<td>NGC 1817</td>
<td>(186.20, -13.10)</td>
<td>1972</td>
<td>-447</td>
<td>8.61</td>
<td>0.43</td>
<td>0.37</td>
<td>13.92</td>
<td>9.92</td>
<td>239</td>
</tr>
<tr>
<td>NGC 2567</td>
<td>(249.80, +02.96)</td>
<td>1677</td>
<td>87</td>
<td>8.47</td>
<td>0.44</td>
<td>0.18</td>
<td>7.51</td>
<td>4.02</td>
<td>290</td>
</tr>
<tr>
<td>Ruprecht 18</td>
<td>(239.93, -04.94)</td>
<td>1056</td>
<td>-91</td>
<td>7.65</td>
<td>0.27</td>
<td>0.22</td>
<td>6.84</td>
<td>6.29</td>
<td>411</td>
</tr>
<tr>
<td>NGC 2354</td>
<td>(238.37, -06.79)</td>
<td>4085</td>
<td>-483</td>
<td>8.13</td>
<td>0.28</td>
<td>0.23</td>
<td>3.35</td>
<td>1.95</td>
<td>51</td>
</tr>
<tr>
<td>Berkeley 39</td>
<td>(223.46, +10.09)</td>
<td>4780</td>
<td>837</td>
<td>9.90</td>
<td>0.18</td>
<td>0.04</td>
<td>7.76</td>
<td>3.55</td>
<td>424</td>
</tr>
<tr>
<td>NGC 2425</td>
<td>(231.50, +03.30)</td>
<td>4053</td>
<td>235</td>
<td>9.20</td>
<td>0.37</td>
<td>0.20</td>
<td>3.23</td>
<td>2.59</td>
<td>372</td>
</tr>
<tr>
<td>NGC 2383</td>
<td>(235.27, -02.46)</td>
<td>1655</td>
<td>-71</td>
<td>7.17</td>
<td>0.47</td>
<td>0.28</td>
<td>4.58</td>
<td>2.95</td>
<td>258</td>
</tr>
<tr>
<td>NGC 2355</td>
<td>(203.39, +11.80)</td>
<td>2200</td>
<td>450</td>
<td>8.85</td>
<td>0.13</td>
<td>0.21</td>
<td>7.47</td>
<td>4.04</td>
<td>343</td>
</tr>
<tr>
<td>NGC 2158</td>
<td>(186.63, +01.78)</td>
<td>5071</td>
<td>158</td>
<td>9.02</td>
<td>0.11</td>
<td>0.07</td>
<td>7.43</td>
<td>2.89</td>
<td>1455</td>
</tr>
<tr>
<td>NGC 2194</td>
<td>(197.25, -02.35)</td>
<td>3781</td>
<td>-155</td>
<td>8.52</td>
<td>0.10</td>
<td>0.21</td>
<td>7.35</td>
<td>4.14</td>
<td>925</td>
</tr>
<tr>
<td>NGC 2204</td>
<td>(226.01, -16.11)</td>
<td>2629</td>
<td>-730</td>
<td>8.90</td>
<td>0.36</td>
<td>0.26</td>
<td>8.27</td>
<td>4.16</td>
<td>311</td>
</tr>
<tr>
<td>NGC 2304</td>
<td>(197.20, +08.90)</td>
<td>3991</td>
<td>67</td>
<td>8.90</td>
<td>0.17</td>
<td>0.12</td>
<td>3.16</td>
<td>2.20</td>
<td>186</td>
</tr>
<tr>
<td>Tombaugh 2</td>
<td>(232.83, -06.88)</td>
<td>13260</td>
<td>-1588</td>
<td>9.01</td>
<td>0.11</td>
<td>0.13</td>
<td>2.90</td>
<td>2.46</td>
<td>95</td>
</tr>
<tr>
<td>IC 348</td>
<td>(160.40, -17.72)</td>
<td>320</td>
<td>-97</td>
<td>6.80</td>
<td>0.47</td>
<td>0.10</td>
<td>8.53</td>
<td>3.87</td>
<td>269</td>
</tr>
<tr>
<td>NGC 1931</td>
<td>(173.90, +00.28)</td>
<td>3086</td>
<td>15</td>
<td>7.00</td>
<td>0.15</td>
<td>0.44</td>
<td>4.13</td>
<td>2.71</td>
<td>316</td>
</tr>
</tbody>
</table>
Fig. 1.— (a) Face-on view and (b) edge-on view of sample open clusters. The sun and Galactic
Fig. 2.— The number density of field stars exhibits gradients, often pointing toward the Galactic plane. In the example shown here for NGC 1893, the projected gradient across a 1 deg squared field, which amounts to $\sim 2 \text{ sq. arcmin}^{-2}$ per degree, is significantly smaller in the Galactic longitude than along the latitude direction.
Fig. 3.— Example of (top) a spherically shaped open cluster, NGC 2414 and (bottom) an elongated open cluster, NGC 1893. In each case, all stars brighter than 15.6 mag, the $\sim 3 - \sigma$ limit of 2MASS Ks band, are plotted on the left panel and the right panel shows the isodensity contours of cluster member stars, superimposed with the ellipses fitted to the core (1/3 maximum number density) and to the outer boundary of the cluster (see text.)
Fig. 4.— NGC 2420, being high above the Galactic plane, has a relatively uniform background stellar density in both the longitudinal and latitudinal directions, as seen in the top panel. The lower panel has the same format as that in Fig. 3 and shows a rather spherical core and a loosely defined outer boundary for this old cluster.
Fig. 5.— The age versus height from Galactic disk for open clusters in the Dias et al. (2002) catalog. Filled circles are the 31 open clusters studied in this paper.
Fig. 6.— Morphology of the old open cluster Berkeley 17 ($\ell \sim 176^\circ$, $b \sim -4^\circ$) shows an upward tail toward Galactic disk. The outer boundary is duly fitted with an ellipse, obviously not appropriate here.
Fig. 7.— The flattening versus age for (top) the outer boundary and (bottom) the inner part of a cluster. Filled circles denote half of the sample of the clusters which are relatively close (less than \( \sim 170 \) pc) to the Galactic plane.
Fig. 8.— The flattening versus the height from the Galactic disk for (top) the outer boundary and (bottom) the inner part of a cluster. Filled circles are for younger star clusters in the sample.
Fig. 9.— The physical size (radius in parsec) and richness ($N_\star$ within the volume, in pc$^{-3}$) of open clusters versus the age. Older clusters tend to be larger in size and less dense in stellar density. Filled symbols are for clusters closer to the Galactic plane (cf. Fig. 7)
Fig. 10.— The physical size (radius in parsec) and richness ($N_*$ within the volume, in pc$^{-3}$) of open clusters versus the height from Galactic disk.