Domain-Wall Fermions at Strong Coupling

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The DWF formulation becomes increasingly problematic at gauge couplings for which \( a^{-1} < 2 \) GeV, where the roughness of the gauge field leads to increased explicit chiral symmetry breaking (\( m_{\text{res}} \)). This problem becomes especially severe for sufficiently strong coupling where the underlying 4-dimensional Wilson theory is in the Aoki phase. We review our attempts to find a suitable modification of the gauge and/or the fermion action which would allow the DWF method to work reliably at stronger coupling.

1. INTRODUCTION

Domain-wall fermions have the advantage of having an exact chiral symmetry in the limit of an infinite 5th dimension on the lattice, even at a non-zero lattice spacing. In combination with the DBW2 gauge action [1] in quenched simulations it has been shown that the \( m_{\text{res}} \) is about two orders of magnitude smaller than its value for the Wilson gauge action [2]. This effect has been attributed to the properties of DBW2 to suppress lattice dislocations, which contribute directly to the \( m_{\text{res}} \), and to make perturbative corrections small [2,3]. However, in full QCD calculations at coarse lattice spacings, for the combination of DWF/DBW2, the \( m_{\text{res}} \) is still not satisfactory small. The quark-gluon system is in the Aoki phase where the DWF formulation suffers from the violations of chiral symmetry and locality. Our goal is to find a suitable modification of the gauge action and/or the fermion action which would allow us to conduct a thermodynamics calculation with DWF.

2. EFFECTS OF THE ADDITION OF AN ADJOINT TERM TO DBW2

Our studies of the plaquette distributions of dynamical DWF/DBW2 configurations [4] showed that with the increase of the coupling the tails of the distributions extend to more negative plaquette values, which are associated with the appearance of greater number of lattice dislocations and increased \( m_{\text{res}} \). These findings gave us the idea to modify the DBW2 action by adding an adjoint term to it with the expectation that the modified action will suppress more strongly the negative tails of the plaquette distributions and thus it will get rid of the lattice dislocations. The modified DBW2 action has the following form:

\[
S_{\text{dbw2+adj}} = -\frac{\beta}{N_c} \left( 1 - 8c_1 \right) \sum_{x, \mu < \nu} \text{Re} \text{Tr} P_{\mu\nu}(x) + \frac{c_0}{N_c} \sum_{x, \mu < \nu} |\text{Tr} P_{\mu\nu}(x)|^2 + c_1 \sum_{x, \mu \neq \nu} \text{Re} \text{Tr} R_{\mu\nu}(x),
\]

where we choose \( c_f + 2c_a = 1 \), to keep the normalization for this three term action the same as for the plaquette action in the continuum limit. The tail shortening effect on the plaquette distributions is an expected result of the decreasing of the effective coupling for the fundamental plaquettes in the action, when the coefficient \( c_a \) is negative.

In the quenched case, we chose to match the plaquette distributions of two runs, one with the modified DBW2 and the other with the original DBW2, and compare their \( m_{\text{res}} \) and lattice spacing. Figure 1 shows that although the plaquette distribution of the modified DBW2 run \( b \) has a shorter negative tail, its \( m_{\text{res}} = 0.0156(2) \) is actually an order of magnitude larger than the corresponding value of \( m_{\text{res}} = 0.00125(3) \) for the unmodified DBW2 run \( a \). Furthermore run \( a \) is

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out of the Aoki phase and run $b$ according to its closed spectral flow gap shown on on Figure 2 is in that phase. In conclusion, we attempted to

\[ \text{Figure 1. Plaquette distributions for quenched runs $a$ and $b$. Run $a$ has } \beta = 0.87 \text{ and } a^{-1} \approx 1.3 \text{ GeV. Run } b \text{ has } \beta = 0.40, c_f = 2.0 \text{ and } a^{-1} \approx 0.93 \text{ GeV. Both runs have volume } 16^3 \times 32 \text{ and } L_s = 12. \]

“cure” the Aoki phase by changing the action at the scale of the lattice spacing by adding an adjoint term and shortening the tails of the plaquette distributions. The negative tail of the plaquette distribution although probably contributing to $m_{\text{res}}$ out of the Aoki phase, when the system is in that phase, does not play a significant role in determining $m_{\text{res}}$. The extent of the tail of the plaquette distribution into negative values is not a primary reason for the onset of the Aoki phase and although the DBW2 with the added adjoint term does reduce it at a given lattice spacing, that does not help to push the system out of the Aoki phase.

3. EFFECTS OF A LARGE RECTANGLE TERM IN THE GAUGE ACTION

The rectangle term in the gauge action reduces the effective coupling for the plaquette term and “smoothes” the gauge field on the lattice by introducing some amount of non-locality at the scale of two lattice spacings. Since we have concluded from the previous section that changing the action at the scale of one plaquette does not influence the Aoki phase, which is a more long-range global phenomenon, we want to investigate the effects of adding more of the less-local rectangle term. In practice we change the coefficient $c_1 = -1.4069$ in the DBW2 action with a larger negative value.

We performed a quenched calculation with a DBW2-style action with $\beta = 0.53$ and $c_1 = -2.3$, which has $a^{-1} \approx 1.0$ GeV and $m_{\text{res}} = 0.0035(1)$. The spectral flow for this run on Figure 2 shows a gap and some crossings which we believe means that this run is out of the Aoki phase even at that coarse $a$. However to use such a large rectangle term for physics calculations could have undesired consequences, since the short distance physics is distorted by the introduced non-locality.

4. TWISTED MASS TERM IN THE DWF DIRAC OPERATOR

As another attempt to solve the problem posed by the Aoki phase we would like to add an irrelevant to the physics term in the DWF action, which would regularize the approach to the Aoki phase. We want to experiment with the term $im_{\gamma_5, s'} \delta_{s, s'}$, called a “twisted mass” term. As an analogue of $M_5$, the twisted mass term provides another parameter, $m_\gamma$, which could be

\[ \text{Figure 2. Spectral flow of } \gamma_5 D_W \text{ for run $a$ and $b$ from Figure 1.} \]
independently tuned to hopefully achieve a better localization of the bound to the domain-walls light chiral states, which, if possible, could help with the Aoki phase problem. We have investigated the free case analytically to make sure that the twisted mass term addition does not break down the most important properties of the DWF formulation, namely, the existence of light bound to the domain walls states. We tested the effects of the twisted mass term on the eigenvalues and eigenmodes of $D_+ D_+^*$, by running the Ritz algorithm on one chosen lattice from the dynamical DWF/DBW2 $\beta = 0.75$ run. The graph in the middle of Figure 4 shows the dependence of the lowest eigenvalue of $D_+ D_+^*$ on $m_\tau$. We see that for the range $m_\tau < 0.8$, the eigenvalue is not strongly affected by the value of the twisted mass, a behavior very similar to the one observed in the free theory (shown on the same graph). This is encouraging since it shows that $m_\tau$, as we hoped by its design, is an irrelevant parameter for a certain range of values, in the sense that it does not act as an additional fermion mass in the problem. The dependence on $m_\tau$ of the density in the fifth dimension of the eigenvector corresponding to the lowest eigenvalue can be studied from the bottom graph on Figure 4. Obviously the chirality and the localization of this mode changes with $m_\tau$ and we expect that this is true for the rest of the eigenmodes. Whether we can manipulate the system by varying $m_\tau$ in a way that globally the zero-eigenmode condensation in the Aoki phase is suppressed, is still to be determined and will be a subject of future investigations.

5. CONCLUSIONS

The addition of an adjoint term to the DBW2 action in order to reduce further $m_{\text{res}}$ at strong couplings did not give a significant improvement. Increasing the amount of the rectangle term in the gauge action gives an improvement in $m_{\text{res}}$ at coarse scales, but at the possible expense of the distortion of short distance physics.

Finally, we showed that the twisted mass term in the DWF Dirac operator does affect the localization of the eigenvectors, which we hope is an evidence that by tuning $m_\tau$ we can achieve better localization of the eigenstates and influence the onset of the Aoki phase.

REFERENCES