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Hadronic Physics with the CMS Experiment

Ferenc Siklér

KFKI Research Institute for Particle and Nuclear Physics, 1121 Budapest, Hungary

CMS collaboration

Abstract

The capabilities of the CMS detector are shown and its Heavy Ion program is outlined.

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Figure 1: Slice through CMS showing particles incident on the different sub-detectors.

1 Introduction

Experiments observing nucleus-nucleus collisions have been studying hot nuclear matter for several decades at ultra-relativistic energies. With a long history (BNL AGS, CERN SPS, BNL RHIC) the next step is again at CERN with the Large Hadron Collider (LHC).

The field has recently brought exciting results, such as anomalous suppression of heavy mesons at the SPS [1, 2], the suppression of high p_T jets [3, 4, 5] and the disappearance of back-to-back high p_T hadron correlations at RHIC [6], ultimately leading to the observation of a new kind of strongly interacting matter.

In the coming LHC era the higher available energy (5.5 TeV per nucleon pair for A+A) will provide a wider kinematic range and higher cross-sections for rare probes. Familiar as well as new observables will be available: enhancement of particles with high p_T ; the "temperature"-indicators J/ψ , ψ' and Υ family; medium effects on jets, such as their shape and fragmentation; dijets, the newly available jet- γ , jet- Z^0 correlations; the centrality dependence of these effects.

2 Heavy ion experiments at LHC

The LHC is being built in the LEP tunnel. It will provide collisions for four experiments, three of them which have a specific heavy ion program. One of them is the Compact Muon Solenoid (CMS) detector which has good acceptance, spatial and momentum resolution, and is able to clearly observe rare signals. CMS is somewhat complementary to the other heavy ion experiments. While ALICE has particle identification capability at moderate p_T , CMS is good at observing higher p_T particles, especially muons, and jets.

The CMS Heavy Ion component is an integral part of the physics program of the experiment with both a detailed A+A and p+A program [7].

3 The CMS detector

The detection of charged and neutral particles, leptons and hadrons is achieved using several detector components: namely the silicon tracker, electromagnetic and hadronic calorimeters, muon chambers (Fig. 1).

The silicon tracker includes pixels and strips, it has geometric coverage for $|\eta| < 2.5$. There are about 10 million microstrips and 40 million pixels with a size of $100 \times 150 \mu\text{m}^2$. The pixel part consists of three barrel and two forward layers on each side. The occupancy of the pixels is expected to be at the few percent level for a multiplicity of $dN/d\eta = 5000$, even for the innermost layer at a radius of 4.5 cm. The strip part consists of ten barrel and nine forward layers on each side, some of which are single-, others are double-sided. The silicon tracker has excellent

Figure 2: Efficiency (higher lines) and fake track rate (lower lines) for A+A collisions at $dN/dy = 1000$ (solid) and $dN/dy = 5000$ (dashed) as a function of p_T .

Figure 3: Comparison of the reconstructed $dN/d\eta$ distribution to the simulated primary tracks as a function of η .

reconstruction performance for $p_T > 1$ GeV/c. For lower p_T the reconstruction capabilities are limited by the high magnetic field and effects of the material in the detector.

The electromagnetic calorimeter is very compact: its lead tungstate crystals have high density, small Molière radius and short radiation length. It includes barrel and endcap parts covering $|\eta| < 3$ with a granularity of $2 \times 2 \text{ cm}^2$, the resolution is $2.7\%/\sqrt{E} \oplus 0.55\%$. The hadronic calorimeter has barrel, endcap and forward parts covering $|\eta| < 5$, with a resolution of $116\%/\sqrt{E} \oplus 5\%$. The barrel and endcap are made of copper absorber plates with scintillator sheets. The forward part is made of iron with quartz fibers for Cerenkov-light detection. A forward calorimeter (CASTOR) can extend the pseudorapidity range up to 7.

The muon system provides an important input to the heavy ion physics program. Drift tubes are used in the barrel part for $|\eta| < 1.5$, cathode strip chambers are employed in endcaps up to pseudorapidity of 2.5. Thus, the coverage is outside the fragmentation regions and complements ALICE, which detects muons in the range $2.5 < \eta < 4.5$.

In the very forward direction a zero degree calorimeter will be used to detect spectator neutrons (and forward photons). This information essential for the determination of the centrality of the A+A collisions.

The high luminosity and collision rate at the LHC requires a good event selection and powerful data acquisition system. (The expected size of Pb+Pb events in CMS is 2-3 Mbytes.) The system consists of a low level (hardware) and a high level (software) trigger: the latter can already make use of tracking information. The high level trigger system has to use around 1000 CPUs to make fast enough decisions, which enable the storage of interesting events at the rate of 40 Hz.

For offline computation enormous resources are needed: a world-wide computing grid of many thousands of CPUs is built up (LHC Computing Grid).

4 Physics studies

4.1 Soft

The tracking of charged particles using the silicon tracker is efficient and results in low fake rates for $p_T > 1$ GeV/c (Fig. 2). An efficiency of 80% at $dN/dy = 5000$ is expected even for tracks in jets. The relative resolution of p_T is about 1% for tracks in the 1-30 GeV/c range. In order to achieve these results a rethinking of the reconstruction algorithms and tuning of their parameters was necessary.

One of the first results of the heavy ion program could be the measurement of charged particle multiplicity using the silicon pixels alone (Fig. 3).

Using the electromagnetic and hadronic calorimeters, azimuthal asymmetry and flow effects can be studied.

4.2 High p_T

The extraction of the quarkonia signal is possible via the decay to $\mu^+\mu^-$: the CMS detector is well suited for detection of muons up to 2.5 in pseudorapidity. The mass resolution of dimuons in the Υ region is $50 \text{ MeV}/c^2$

Figure 4: Mass spectrum of opposite sign dimuons after background subtraction, with clear signs of J/ψ and the Υ family.

Figure 5: Reconstruction efficiency and purity of jets in the barrel.

Figure 6: Azimuthal correlation of a 30 GeV/event jet on a $dN/dy \sim 3000$ background: simulated and reconstructed signal.

Figure 7: Distribution of energy difference between the jet and its photon partner at several average energy loss values.

(Fig. 4) [8], as compared to about 100 MeV/ c^2 in ALICE. The background comes from $\pi/K \rightarrow \mu$ decay and via $c\bar{c}$ or $b\bar{b}$ production.

One of the strong points of the CMS Heavy Ion program is the detection of high energy jets, using calorimetry [9] and also tracking. Special jet finding algorithms have been developed which are able to work in the high background of a central A+A collision (subtraction of background, contribution from tracks, counting energy in a cone). High purity and reconstruction efficiency can be reached for jets with $E_T > 50$ GeV (Fig. 5). Good linearity is seen with a jet energy resolution of 16% at 100 GeV.

Correlation of jets can be studied as well: good angular resolution is expected (Fig. 6).

Due to the higher available energy, the correlation of the jet with a weakly interacting partner can be observed: jet- γ and jet- Z^0 (Fig. 7). This way the energy loss and partonic propagation effects of the individual jets through dense matter can be directly measured and compared to predictions: the back-to-back photon or Z^0 should give the unaffected p_T of the jet at the point of its initial production.

5 Summary

The CMS Heavy Ion group is a growing international collaboration with an exciting physics program. Topics to be studied include particle multiplicities, spectra and correlations, with an extensive high p_T program looking at quarkonia, medium effects on jets and jet correlations.

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