The R Coronae Borealis stars - carbon abundances from forbidden carbon lines

Gajendra Pandey\textsuperscript{1}, David L. Lambert\textsuperscript{1}, N. Kameswara Rao\textsuperscript{2}, Bengt Gustafsson\textsuperscript{3}, Nils Ryde\textsuperscript{3}, David Yong\textsuperscript{1}

\textsuperscript{1}Department of Astronomy, University of Texas, Austin, TX 78712-1083, USA
\textsuperscript{2}Indian Institute of Astrophysics, Bangalore 560034, India
\textsuperscript{3}Astronomiska observatoriet, Uppsala, Sweden S-751-20

ABSTRACT

Spectra of several R Coronae Borealis (RCB) stars at maximum light were examined for the $[\text{C}\, \text{i}]$ 9850Å and 8727Å absorption lines. The 9850Å line is variously blended with a Fe\,ii and CN lines but positive identifications of the $[\text{C}\, \text{i}]$ line are made for R\,CrB and SU\,Tau. The 8727Å line is detected in the spectrum of the five stars observed in this wavelength region. Carbon abundances are derived from the $[\text{C}\, \text{i}]$ lines using the model atmospheres and atmospheric parameters used by Asplund et al. (2000).

Although the observed strength of a $[\text{C}\, \text{i}]$ line is constant from cool to hot RCB stars, the strength is weaker than predicted by an amount equivalent to a factor of four reduction of a line’s $gf$-value. Asplund et al. dubbed this ‘the carbon problem’ and discussed possible solutions.

The $[\text{C}\, \text{i}]$ 9850Å line seen clearly in R\,CrB and SU\,Tau confirms the magnitude of the carbon problem revealed by the $[\text{C}\, \text{i}]$ lines. The $[\text{C}\, \text{i}]$ 8727Å line measured in five stars shows an enhanced carbon problem. The $gf$-value required to fit the observed $[\text{C}\, \text{i}]$ 8727Å line is a factor of 15 less than the well-determined theoretical $gf$-value. We suggest that the carbon problem for all lines may be alleviated to some extent by a chromospheric-like temperature rise in these stars. The rise far exceeds that predicted by our non-LTE calculations, and requires a substantial deposition of mechanical energy.

Key words: stars: abundances – stars: atmospheres – stars: variables – stars: R Coronae Borealis.

1 INTRODUCTION

The rare class of R Coronae Borealis (RCB) variables continues to present fascinating puzzles. Two defining characteristics of these supergiants are the unpredictable fading of the star and the presence of helium, not hydrogen, as the most common constituent of the stellar atmosphere. The fadeings are generally assumed to occur when a cloud of carbon soot forms on the line of sight to the star, but the formation process, and initial location and evolution of the cloud encompass unanswered questions. How a normal star or a pair of stars evolves to possess a H-deficient atmosphere has challenged theorists and observers alike. A leading idea is that an RCB star is the result of a merger of a He white dwarf with a C-O white dwarf. To develop observational tests of this and other ideas, one requires accurate estimates of the elemental abundances for a representative sample of the RCB stars.

In attempting to supply this data, Asplund et al. (2000) uncovered yet another interesting puzzle which is discussed below. In RCB stars, as in normal stars, the strength of an absorption line is influenced by two principal factors: the ratio of the line to continuous opacity and the temperature gradient through the region of the atmosphere in which the line is formed. It has been suggested that photoionization of neutral carbon from high-lying levels appears to be the dominant source of continuous absorption across the visible spectrum (Searle 1961). A characteristic of RCB star spectra is a rich collection of absorption lines of neutral carbon from levels only slightly less excited than those contributing to the continuous opacity. Since the line and continuous opacity...
originate from similar excited levels of the neutral carbon atoms, their opacity ratio is quite insensitive to fundamental parameters of the stars such as effective temperature, surface gravity, and the fractional carbon abundance (effectively, the C/He ratio once it exceeds 1% – see also Schönberner 1975; Cottrell & Lambert 1982). For weak lines, this insensitivity extends to differences in the turbulent velocities.

This is observed in the RCB stars. A C1 line retains its equivalent width even as ‘metal’ lines may vary considerably from star to star (see Figure 1 in Rao & Lambert 1996). However, the predicted equivalent widths of the C1 lines exceed the observed equivalent widths by a considerable factor. Expressed as the change in a line’s gf-value needed to bring the predicted in tune with the observed equivalent width, the disagreement amounts to a factor of four. Asplund et al. dubbed this ‘the carbon problem’.

The principal purpose of this paper is to add a new clue to the solution to the carbon problem. That dimension is offered by the [C1] 8727 Å and 9850 Å lines and the estimate of the carbon abundance that they provide. We propose to investigate this aspect in more stars and to see whether a solution can be obtained for this puzzle. The present paper provides the results of these observations which, to our surprise, show an enhanced carbon problem.

2 OBSERVATIONS

High-resolution optical spectra of eight RCB stars (Table 1) were obtained at the W. J. McDonald Observatory’s 2.7-m telescope with the coudé cross-dispersed echelle spectrograph (Tull et al. 1995). In addition, γ Cyg, a normal F-type supergiant was also observed (Table 1) at R = 120,000. The telluric absorption lines from the spectra of the programme stars were removed interactively using early-type rapidly rotating stars. We have used the Image Reduction and Analysis Facility (IRAF) software packages to reduce the spectra, and the task telluric within IRAF to remove the telluric absorption lines.

The Th-Ar hollow cathode lamp fails to provide lines in the 9850Å region for wavelength calibration. Plenty of Th-Ar lines are available shortward of 9670Å to provide a wavelength solution to an accuracy of one tenth of a pixel. We apply this solution for the wavelengths longer to 9670Å. To check the accuracy of the wavelength calibration in the 9850Å region, the measured wavelengths of the atmospheric H2O lines are compared with those measured from the solar spectrum (Swensson, Benedict, Delbouille & Roland 1970). We measured the wavelengths of 11 atmospheric H2O lines finding agreement to 0.003Å ± 0.02Å with those measured from the solar spectrum.

Sample portions of R = 120,000 and, R = 60,000 spectra are shown in Figures 1, and 2. Note that, the spectra of the RCB stars are fully resolved at R = 60,000. All spectra are aligned to the rest wavelengths of N1 lines which fall in the wavelength regions. Stellar lines were identified using the Revised Multiplet Table (Moore 1972), tables of spectra of H, C, N, and O (Moore 1993), Kurucz’s list*, the Vienna Atomic Line Database†, the new Fe1 multiplet table (Nave et al. 1994), and a spectrum of γ Cyg, a normal F-type supergiant, as a reference.

All stars except V482 Cyg were observed at maximum light. Our spectrum of V482 Cyg was taken when the star was in a shallow minimum in July 1996, about 2 magnitudes below maximum light. It took the star about one year to return to maximum light.

3 LINE DATA AND IDENTIFICATIONS

3.1 The [C1] lines

The standard reference for the C1 spectrum (Moore 1993) gives a predicted wavelength of 8727.126 Å for the [C1] 2p2 1D2 − 2p2 1S0 transition and 9850.264 Å for the [C1] 2p2 3P2 − 2p2 1D2 transition. Wavelengths of the forbidden carbon lines have not been measured directly from laboratory sources.

The predicted wavelength of the 8727 Å line is confirmed by the wavelength of the absorption line in the solar spectrum (Allende Prieto, Lambert, & Asplund 2002) and the emission line in planetary nebulae (Liu et al. 1995). For the 8727 Å line, we adopt log gf = −8.14 (Galavis, Mendoza & Zeippen 1997). The lower excitation potential is 1.264 eV.

There is supporting evidence from our spectra for this wavelength. The 2002 July spectra of RY Sgr show the 9850Å line, as the rest wavelength for the latter emission and assuming the two lines have the same velocity, we find the astrophysical wavelength of the 9850 Å emission is close to the predicted wavelength. Liu et al. (1995) report, however, a rest wavelength of 9850.36 ± 0.01Å from detections of the line in four planetary nebulae. The 0.1Å or 3km s−1 difference between the wavelength predicted from energy levels and that observed from planetary nebulae does not seriously impact derivation of the carbon abundance from the line in the RCB star spectra. The transition probability for the 9850 Å line is taken from Galavis, Mendoza & Zeippen (1997) who give A = 2.23 × 10−4 s−1 or log gf = −10.79 with an uncertainty of about 0.04 dex. The lower excitation potential is a mere 0.005 eV.

Two additional lines complete multiplet 1F. The stronger line with a gf-value about 3 times smaller than for the 9850 Å line is at 9824.30 ± 0.05 Å (Liu et al. 1995) and irretrievably blended in the spectra of cooler RCB stars (GU Sgr, V482 Cyg, SU Tau, and R CrB). In the spectra of VZ Sgr, UV Cas, and XX Cam, the 9824.30 Å line is absent, but this is as expected from the S/N and the estimated strength of the observed 9850 Å line. The third [C1] line at 9808.3 Å is expected to be 4000 times weaker than the 9850 Å line, and undetectable. (The multiplet 2F 2p2 3P1 − 2p2 1S0 transition at 4621.57 Å with log gf = −11.12 is almost

* http://kurucz.harvard.edu
† http://www.astro.univie.ac.at/~vald
Table 1. Table of observations; the stars are ordered by their effective temperatures decreasing from top to bottom

<table>
<thead>
<tr>
<th>Star</th>
<th>Date</th>
<th>R = λ/∆λ at 8727 Å</th>
<th>S/N at 8727 Å</th>
<th>S/N at 9850 Å</th>
<th>V_{rad}</th>
<th>km s^{-1}</th>
</tr>
</thead>
<tbody>
<tr>
<td>RCB stars:</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>XX Cam</td>
<td>9 Oct 1997</td>
<td>60,000</td>
<td>...</td>
<td>185</td>
<td>11</td>
<td></td>
</tr>
<tr>
<td>XX Cam</td>
<td>26 Jan 1998</td>
<td>60,000</td>
<td>...</td>
<td>120</td>
<td>13</td>
<td></td>
</tr>
<tr>
<td>XX Cam</td>
<td>17 Nov 2002</td>
<td>60,000</td>
<td>...</td>
<td>186</td>
<td>11</td>
<td></td>
</tr>
<tr>
<td>XX Cam</td>
<td>17 Nov 2002</td>
<td>60,000</td>
<td>278</td>
<td>...</td>
<td>...</td>
<td></td>
</tr>
<tr>
<td>RY Sgr</td>
<td>22 June 1997</td>
<td>60,000</td>
<td>...</td>
<td>170</td>
<td>−30</td>
<td></td>
</tr>
<tr>
<td>RY Sgr</td>
<td>31 July 2002</td>
<td>120,000</td>
<td>300</td>
<td>200</td>
<td>−12</td>
<td></td>
</tr>
<tr>
<td>UV Cas</td>
<td>24 July 1996</td>
<td>60,000</td>
<td>...</td>
<td>140</td>
<td>−32</td>
<td></td>
</tr>
<tr>
<td>UV Cas</td>
<td>15 Nov 2002</td>
<td>60,000</td>
<td>...</td>
<td>120</td>
<td>−32</td>
<td></td>
</tr>
<tr>
<td>UV Cas</td>
<td>15 Nov 2002</td>
<td>60,000</td>
<td>178</td>
<td>...</td>
<td>...</td>
<td></td>
</tr>
<tr>
<td>VZ Sgr</td>
<td>18 Aug 1999</td>
<td>60,000</td>
<td>...</td>
<td>80</td>
<td>245</td>
<td></td>
</tr>
<tr>
<td>R CrB</td>
<td>17 June 1995</td>
<td>60,000</td>
<td>340</td>
<td>180</td>
<td>20</td>
<td></td>
</tr>
<tr>
<td>R CrB</td>
<td>7 Aug 1995</td>
<td>60,000</td>
<td>...</td>
<td>110</td>
<td>22</td>
<td></td>
</tr>
<tr>
<td>R CrB</td>
<td>20 June 1997</td>
<td>60,000</td>
<td>...</td>
<td>185</td>
<td>21</td>
<td></td>
</tr>
<tr>
<td>R CrB</td>
<td>31 July 2002</td>
<td>120,000</td>
<td>400</td>
<td>200</td>
<td>19</td>
<td></td>
</tr>
<tr>
<td>SU Tau</td>
<td>15 Nov 2002</td>
<td>60,000</td>
<td>...</td>
<td>100</td>
<td>43</td>
<td></td>
</tr>
<tr>
<td>SU Tau</td>
<td>15 Nov 2002</td>
<td>60,000</td>
<td>180</td>
<td>...</td>
<td>...</td>
<td></td>
</tr>
<tr>
<td>V482 Cyg</td>
<td>24 July 1996</td>
<td>60,000</td>
<td>...</td>
<td>90</td>
<td>−40</td>
<td></td>
</tr>
<tr>
<td>GU Sgr</td>
<td>22 June 1997</td>
<td>60,000</td>
<td>...</td>
<td>95</td>
<td>−45</td>
<td></td>
</tr>
<tr>
<td>Standard star:</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>γ Cyg</td>
<td>31 July 2002</td>
<td>120,000</td>
<td>520</td>
<td>230</td>
<td>...</td>
<td></td>
</tr>
</tbody>
</table>

certainly blended as the blue spectral region of an RCB star is rich in strong lines.)

3.2 Overview of the Spectra

The region around 9850 Å is shown in Figure 1 where the stars are ordered by increasing temperature from bottom to top. The spectrum of the normal yellow supergiant γ Cyg is shown at the bottom. In the spectra of the hottest stars, the strongest lines are from C i and N i transitions with weaker lines of Si i, Fe i, and Fe ii. The [C i] line falls in the red wing of a Fe ii line. Spectra of the coolest stars GU Sgr and V482 Cyg show additional lines which we identify as high rotational lines of the 1-0 band of the CN molecule’s Red System (Davis & Phillips 1963). These CN lines are prominent in the spectra of GU Sgr and V482 Cyg and traceable in SU Tau, R CrB and RY Sgr. Two CN lines bracket the [C i] line.

Inspection of Figure 1 shows several additional features of interest. As anticipated, the strength of a given C i line is unchanged along the temperature sequence from XX Cam at the high temperature end to GU Sgr at the low temperature end. Also anticipated is the very large width of all absorption lines in the spectra of RCB stars. Compare the widths with those of lines in γ Cyg, which itself is commonly referred to as showing broad lines. A surprise is the appearance of emission in the RY Sgr spectrum near the wavelength of the [C i] line. The illustrated spectrum is from 31 July 2002 in which the 8727 Å line is also in emission. A spectrum from 22 June 1997 also shows emission but to the red of the wavelength of the [C i] line. The possibility of emission contaminating the absorption line is an unfortunate complication.

The [C i] 8727 Å line (Figure 2) is slightly blended with a Si i line. Emission is very pronounced in the spectrum of RY Sgr. The central line depths of the [C i] line run from about 16 % for SU Tau to 36 % for RY Sgr with no obvious systematic trend with the effective temperature.
8710 8720 8730

Figure 2. The spectra of RCB stars in the region of \([\text{C} \text{I}]\) 8727.126˚A line, which is indicated by a vertical solid line, are plotted with their effective temperatures increasing from bottom to top. The spectrum of \(\gamma\) Cyg is also plotted. The positions of the key lines are also marked.

4 SPECTRUM SYNTHESIS - PROCEDURE

The fact that the \([\text{C} \text{I}]\) lines are blended with other lines requires that the carbon abundance be extracted by spectrum synthesis. In addition to synthesizing the \([\text{C} \text{I}]\) line and its blends, we use lines of C1, N1, Si1, Fe1 and Fe1II to obtain abundances and compare with values given by Asplund et al. (2000).

Synthetic spectra are generated using the H-deficient model atmospheres computed by Asplund et al. (1997) and the Uppsala spectrum synthesis code BSYNRUN. Models corresponding to a C/He ratio of 1% by number were chosen. This composition is equivalent to a carbon abundance \(\log e(C) = 9.54\) on the scale \(\Sigma m_i \xi_i = 12.15\), where \(m_i\) is the mean atomic weight of element \(i\). Stellar parameters derived by Asplund et al. (2000) were adopted as our initial values. Synthetic spectra were convolved with a Gaussian profile to account for the combined broadening of stellar lines by the atmospheric macroturbulent velocity, rotational broadening, and the instrumental profile.

4.1 The 9850 Å Region

The 9850 Å \([\text{C} \text{I}]\) line may be blended with Fe1II and CN lines. The Fe1II line at 9849.74 Å has a lower excitation potential of 6.729 eV. To obtain the \(gf\)-value of the line, we use the spectrum of \(\gamma\) Cyg where the Fe1II line is present. With a model atmosphere constructed for the stellar parameters and iron abundance given by Luck & Lambert (1981), we find \(\log gf = -2.60\). The basic data for the CN 1-0 lines were taken from Kurucz’s line list and checked against the SCAN tape (Jørgensen & Larsson 1990) and a recipe given by Bakker & Lambert (1998). A dissociation energy of 7.65 eV was adopted for the molecule (Bauschlicher, Langhoff & Taylor 1988). Unblended CN lines away from the 9850 Å feature were fitted by adjusting the N abundance, and this abundance used in computing the CN lines blended with the \([\text{C} \text{I}]\) line. An estimate of the N abundance was obtained not only from the CN lines but also from N1 lines near the \([\text{C} \text{I}]\) lines. The \(gf\)-values for the N1 lines were taken from Wiese, Fuhr & Deters (1996) compilation.

The \([\text{C} \text{I}]\) lines at 9852.27 Å, a blend of two lines, and 9859.15 Å (Figure 1) have also been used for estimating the C abundance. Their \(gf\)-values are determined by inversion of their equivalent widths in the \(\gamma\) Cyg spectrum using the C abundance derived from a set of weak \([\text{C} \text{I}]\) lines at shorter wavelengths.

4.2 The 8727 Å Region

The red wing of the \([\text{C} \text{I}]\) line is blended with a Si1 line (Figure 2). The \(gf\)-value of the latter line is derived from its equivalent width in the solar spectrum. The \(gf\)-value of the N1 line to the red of the Si1 line is taken from Wiese, Fuhr & Deters (1996). Spectrum synthesis of this region was confined to a 6 Å window around the \([\text{C} \text{I}]\) line.

5 SPECTRUM SYNTHESIS - RESULTS

The \([\text{C} \text{I}]\) lines were sought in order to shed new light on the ‘carbon problem’. Asplund et al. (2000) suggested that the 9850 Å \([\text{C} \text{I}]\) line was not subject to the carbon problem. Their analysis was flawed in that the contributions of the blending Fe1II and CN line were not considered.

In addition to fitting the \([\text{C} \text{I}]\) lines, we analysed other lines in order to check the abundances given by Asplund et al. (2000). To within the uncertainties of the analyses and with the exception of V482 Cyg (see below), we confirm the result for carbon using \([\text{C} \text{I}]\) lines in the 9850 Å region, the N abundance using N1 lines in the 9850 Å and 8727 Å regions, the Si abundance from Si1 lines in the 9850 Å and 8727 Å regions, and the Fe abundance from Fe1 and Fe1II lines in the 9850 Å region. The CN 1-0 lines are also used to check the N abundance. In these atmospheres, molecules are trace species whose formation does not affect the partial pressures of the constituent atoms. Then, consideration of the ratio of line to continuous opacity shows that the strength of a CN line should be dependent on the N abundance but independent of the C abundance.

Following remarks about \(\gamma\) Cyg, we discuss the RCB stars in order of decreasing effective temperature. It should be noted that the C abundances given here for the RCB stars are not self-consistent – the model atmospheres and their continuous absorption were calculated assuming a logarithmic C abundance of 9.54. If the same abundance is adopted for the calculations of the carbon lines – permitted and forbidden – the predicted and observed equivalent widths are in disagreement. Thus, the C abundance derived here for the RCB stars is basically a measure of the inconsistency, which we term the carbon problem.

5.1 \(\gamma\) Cyg

Our analysis of permitted and forbidden C1 lines uses a MARCS model atmosphere (Gustafsson et al. 1975) for
the atmospheric parameters derived by Luck & Lambert (1981). The 8727 Å [C\textsc{i}] line (Figure 2) yields the abundance \(\log \epsilon(C) = 7.88\). The 9850 Å [C\textsc{i}] line is not detectable and the upper limit to its equivalent width corresponds to \(\log \epsilon(C) \leq 8.18\). A selection of 14 C\textsc{i} lines from 5380 Å to 8873 Å with \(gf\)-values from Wiese, Fuhr & Deters (1996) give the abundance \(\log \epsilon(C) = 7.94 \pm 0.14\). These lines range in equivalent width from 6 mA to 82 mA. For \(\gamma\) Cyg, it is seen that the C\textsc{i} lines and the 8727 Å forbidden line return the same abundance within the uncertainties. This result is in sharp contrast to the results from the RCB stars.

### 5.2 XX Cam

The 9850 Å feature is dominated by the Fe\textsc{ii} line such that we can set only an upper limit to the [C\textsc{i}] contribution and the associated C abundance of \(\log \epsilon(C) \leq 8.8\) (Figure 3). The 8727 Å line (Figure 4) is fit with an abundance \(\log \epsilon(C) = 8.4\). The limit from the 9850 Å line to the C abundance is consistent with that from the 8727 Å line. The two C\textsc{i} lines near 9850 Å give an abundance consistent with the C abundance of \(\log \epsilon(C) = 9.0 \pm 0.4\) from Asplund et al. who used a collection of C\textsc{i} lines and Opacity Project \(gf\)-values (Lou & Pradhan 1989; Hibbert et al. 1993; Seaton et al. 1994).

Asplund et al. put the uncertainty in \(T_{\text{eff}}\) at about \(\pm 250\) K which corresponds to corrections to abundances derived from C\textsc{i} and [C\textsc{i}] lines of 0.05, and 0.2 dex, respectively. The uncertainty in \(\log g \simeq \pm 0.5\) provides only minor corrections to the abundances. An additional potential source of uncertainty is the adopted microturbulent velocity of \(\xi = 9.0\) km s\(^{-1}\) (Asplund et al. 2000). A change of \(\xi\) by \(\pm 2.0\) km s\(^{-1}\) changes the C abundance by 0.2 dex from the C\textsc{i} lines and by only 0.05 dex from the [C\textsc{i}] 8727 Å line. These estimates of uncertainty do not allow the 0.6 dex difference in the C\textsc{i} and [C\textsc{i}] abundances.

The carbon problem for the [C\textsc{i}] line is even more severe as seen for 8727 Å [C\textsc{i}] line which gives an abundance 0.6 less than the C\textsc{i} lines and 1.1 dex less than the input abundance.

### 5.3 RY Sgr

The 31 July 2002 spectrum shows emission in the blue wing of the 8727 Å forbidden line. Peak emission occurs at the velocity of \(-29\) km s\(^{-1}\). This is blue-shifted by about 17 km s\(^{-1}\) from the photospheric velocity of \(-12\) km s\(^{-1}\). The emission line has a width (FWHM) of about 19 km s\(^{-1}\). A spectrum taken the same night shows emission at 8580 Å (Figure 1). Emission at [C\textsc{i}] may be a common occurrence for this star. Our 1997 June 22 spectrum shows 9850 Å in emission at the radial velocity of \(-10.5\) km s\(^{-1}\) equivalent to a red-shift of 19.5 km s\(^{-1}\) relative to the photosphere. This spectrum does not include the region around the 8727 Å line.

RY Sgr exhibits a pulsation with an amplitude of about 35 km s\(^{-1}\) (Lawson, Cottrell & Clark 1991). The mean or systemic velocity was given as \(-21\) km s\(^{-1}\) by Lawson, Cottrell & Clark (1991). Relative to this velocity the emission at 8727 Å in the 2002 spectrum is blue-shifted by about 8 km s\(^{-1}\). Perhaps, the emission at maximum light is related to the fact that RY Sgr is a large-amplitude radial velocity pulsator. There is no hint of emission in our spectra of the other RCB stars, even in V482 Cyg observed below maximum light.

By inspection, the 9850 Å blend of Fe\textsc{ii} and [C\textsc{i}] lines
is very similar to that of XX Cam and, hence, the forbidden line is a minor contributor to the absorption feature. Synthesis of the 8727 Å line suggests that emission overlays the absorption line and one might suppose that an abundance log $\epsilon(C) \sim 9.2$ is a lower limit but this assumes that the emission comes from a region exterior to the regions in which the absorption line is formed. This is not necessarily a valid assumption for RY Sgr with its large pulsation and evidence for an internal shock. We decline to quote a carbon abundance. But, a clue might be derived from RY Sgr observations that the emission effects both [C I] lines. The estimations from the red absorption wing (assumed to be unaffected by emissions) of both these [C I] lines suggests the same carbon abundance of about 9.2 which agrees with that from C I lines.

With the adopted model, the CN lines return a N abundance about 0.6 dex higher than from the N I lines used here and by Asplund et al. This difference between N abundances can be eliminated by lowering the adopted $T_{\text{eff}}$ by about 200 K, or by a combination of a smaller $T_{\text{eff}}$ change accompanied by a reduction of log $g$. These adjustments are within the uncertainty limits considered reasonable by Asplund et al.

### 5.4 UV Cas

The 8727 Å line is reasonably well fit with the abundance log $\epsilon(C) = 8.3$ (Figure 4). Two observations show no convincing evidence for the 9850 Å [C I] line, indicating log $\epsilon(C) \leq 8.7$ (Figure 3). The two C I lines in the 9850 Å region suggest an abundance close to log $\epsilon = 9.0$, a value only 0.2 dex less than that obtained by Asplund et al. (2000). The CN lines are not detectable in our spectra. UV Cas joins XX Cam in showing a larger carbon problem for the 8727 Å line than for the C I lines.

### 5.5 VZ Sgr

This RCB star was classified as a ‘minority’ star by Lambert & Rao (1994), i.e., the metal lines are weak. This is evident from Figure 1 which shows that the Fe I and Fe II lines are absent from VZ Sgr’s spectrum. Unfortunately, the region around 8727 Å has not yet been observed. Synthesis (Figure 3) suggest an upper limit log $\epsilon(C) \leq 8.7$ for the [C I] 9850 Å line, a limit clearly below the input abundance of 9.5. The two C I lines give log $\epsilon(C) = 9.1$. Asplund et al. gave log $\epsilon(C) = 8.8 \pm 0.3$.

### 5.6 R CrB

At 8727 Å the synthesis corresponding to log $\epsilon(C) \approx 8.4$ provides a fair fit to the observed high-resolution profile of 2002 July 31 (Figure 4). A lower resolution ($R = 60000$) spectrum of 1995 June 17 provides the same abundance. In neither spectra is there a hint that emission has distorted the absorption profile.

The 9850 Å syntheses include the CN lines blending with the [C I] and Fe II feature. On the high-resolution spectrum of 2002 July 31, the CN 1-0 lines are blue-shifted by 3.6 km s$^{-1}$ relative to the high-excitation lines of C I and N I. Since the [C I] line and CN lines are all of low excitation, we assume that they have a common blueshift. Synthetic spectra (Figure 3) indicate an abundance log $\epsilon(C) = 9.1$. Examination of our library of 9850 Å spectra of R CrB (e.g., Rao & Lambert 1997; Rao et al. 1999) shows that the velocity shift between CN and the C I and N I varies during the star’s pulsation and this shift ranges from $-3.0$ to $+4.0$ km s$^{-1}$. Syntheses taking into the account the velocity shift return the same abundance from four different spectra. If the blue-shift is neglected, the fit to the 9850 Å feature is less satisfactory but the derived carbon abundance is little affected. Asplund et al.’s atmospheric parameters are used in all cases. The C I lines near 9850 Å are well fitted with a similar abundance (log $\epsilon(C) = 9.2$), an abundance equal to that determined by Asplund et al.

The 9850 Å line confirms the abundance obtained from C I lines, but the 8727 Å line gives an abundance 0.7 dex below that from the 9850 Å line or 1.1 dex below the input carbon abundance of the model.

The N I lines in the 8727 Å and 9850 Å regions give the abundance log $\epsilon(N) = 8.2$, a value consistent with the result of 8.4 ± 0.2 given by Asplund et al. The observed CN 1-0 lines are well matched with the abundance log $\epsilon(N) = 8.2$. This may be fortuitous agreement because the sensitivity of the CN line strengths to a change of the atmospheric parameters, especially to $T_{\text{eff}}$, is high. Note that a change of 250 K changes the required abundance by 0.5 dex.

### 5.7 SU Tau

The abundance log $\epsilon(C) = 8.2$ provides an excellent fit to the observed [C I] line at 8727 Å (Figure 4), whereas log $\epsilon(C) = 8.9$ is required for an equivalent fit to the observed 9850 Å feature (Figure 3). For this fit, the CN lines are blue-shifted by 6 km s$^{-1}$ from the high excitation atomic lines. This shift is assumed to apply to the [C I] line also, but the derived abundance is not critically influenced by the shift. The abundance from the 9850 Å [C I] line is consistent with that derived by us from the 9850 Å region’s C I lines and the C I lines used by Asplund et al. In contrast and consistent with the results for XX Cam and R CrB, the 8727 Å line gives a markedly lower carbon abundance.

The N abundance reported by Asplund et al. was based on just two N I lines. Examination of our superior spectra provides a more accurate equivalent width for one line. The other line is not present on our spectrum. We identify a third line. We adopt Wiese et al.’s $gf$-value and derive the N abundance log $\epsilon(N) = 7.9$. We suggest that this is a more reliable estimate than Asplund et al.’s value of 8.5. The N abundance estimate from the CN lines is 0.5 dex lower. The difference is erasable with minor adjustments to the adopted atmospheric parameters, such as an increase of $T_{\text{eff}}$ by only 140 K.

### 5.8 V482 Cyg

There is a hint of emission at the 9850 Å [C I] line. Examination of the three individual exposures shows that this ‘emission’ occurs in only one exposure that being the one with the lowest signal-to-noise ratio. We, therefore, consider the emission to be an artefact which is removed in Figure 3.
A fit to the absorption feature suggests an abundance upper limit \( \log \epsilon(C) \leq 8.9 \).

In conflict with this upper limit, the C abundance of \( \log \epsilon(C) = 9.5 \pm 0.3 \) is derived from C1 lines in our spectrum. Asplund et al. obtained \( \log \epsilon(C) = 8.9 \) from C1 lines. Examination of our spectrum shows that C1 lines across the spectrum are systematically stronger than reported by Asplund et al. This strengthening and the adoption of the same model as Asplund et al. necessarily results in a higher carbon abundance. Recall that the star was two magnitudes below maximum light. We suppose that the atmospheric structure differed from that prevailing at maximum light. The temperature gradient may have been steeper in decline than at maximum light. A very similar strengthening of C1 (and other) lines was noted by Rao et al. (1999) in spectra of R CrB taken early in a deep decline.

The N1 lines give an abundance 0.8 dex higher than that from the CN lines. Erasure of the N1 – CN discrepancy requires a different choice of atmospheric parameters. The possibilities range from raising \( T_{\text{eff}} \) by 400 K at the adopted \( \log g \) to lowering \( \log g \) by about 0.8 dex at the adopted \( T_{\text{eff}} \). These are not acceptable changes given that the atmosphere may have been perturbed.

### 5.9 GU Sgr

The blend containing the 9850 Å [C1] line is contaminated with CN lines which are strong in this spectrum. Our synthesis suggests an upper limit \( \log \epsilon(C) \leq 9.0 \), a value consistent with Asplund et al.’s value of 8.8.

Analysis of the CN lines gives a N abundance of \( \log \epsilon(N) = 8.1 \). Since the N1 lines in the 9850 Å region are blended, we compare the CN-based abundance with the result \( \log \epsilon(N) = 8.7 \pm 0.5 \) given by Asplund et al. The difference may be removed by a modest (250 K) increase of the adopted \( T_{\text{eff}} \) at the adopted \( \log g \).

### 5.10 Summary of the Carbon Abundances

Our analyses use the appropriate MARCS model for the atmospheric parameters recommended by Asplund et al. and C/He = 1% by number. This input abundance corresponds to a carbon abundance \( \log \epsilon(C) = 9.5 \). In Table 2, we summarize the carbon abundances derived from the two [C1] lines, and the C1-based abundance given by Asplund et al. illustrating the carbon problem.

Except for V482 Cyg, the carbon abundance derived from the two 9850 Å region C1 lines is in good agreement with that obtained by Asplund et al.

Table 2 shows that a carbon problem extends to the 9850 Å forbidden line. This conclusion is contrary to that reached by Asplund et al. from analysis of equivalent widths of the 9850 Å line. The explanation is that Asplund et al. did not recognize that the [C1] line was a blend. In the case of V482 Cyg, which was observed two magnitudes below maximum light, the carbon problem has vanished for the C1 lines but not for the 9850 Å [C1] line. This is the only case where the C1 lines and 9850 Å differ in the abundance they provide.

Surprisingly, the 8727 Å [C1] line offers further information on the carbon problem. For each of the four stars for which we have observed the 8727 Å line, the derived abundance is less than that from the C1 lines. In the case of R CrB and SU Tau, the 8727 Å line gives an abundance 0.7 dex less than that derived from the 9850 Å line. That the two forbidden lines give very different abundances is especially puzzling if these lines are formed, as expected, in or close to LTE and, as might be supposed, are a product of the stellar photosphere, i.e., a region with the temperature decreasing monotonically outward. An unidentified atomic line superimposed on the 8727 Å forbidden line would mean that we have overestimated the carbon abundance from that line. In this scenario, the carbon problem is more severe for both of the [C1] lines than for the C1 lines. However, it is difficult to find a carrier for a line which is strong in R CrB and SU Tau but weak in other RCB stars and γ Cyg.

The sensitivities of the permitted and forbidden lines to the choice of model atmosphere are different. To illustrate these sensitivities, we give in Table 3 the abundances derived from R CrB’s lines for a series of model atmospheres centred on Asplund et al.’s choice of \((T_{\text{eff}}, \log g, \xi) = (6750, 0.5, 7)\). The carbon abundance from C1 lines is, as expected from the insensitivity of the ratio of line to continuous opacity to physical conditions, almost independent of the choice of model. The abundance is also insensitive to the choice of the microturbulent velocity \( \xi \) for a weak line but not for a strong line. In Table 3, the mean C abundance derived from weak and strong permitted lines is given. The abundance from the forbidden lines is insensitive to the choice of the surface gravity but dependent on \( T_{\text{eff}} \). We note that the change ± 250 K in \( T_{\text{eff}} \) leads to a change in the carbon abundance of ± 0.2 dex. There is no dependence on \( \xi \). The \( T_{\text{eff}} \) and \( \log g \) sensitivities are not very different for the hottest stars (XX Cam and RY Sgr) and the coolest (GU Sgr).

To increase the carbon abundance from the 9850 Å [C1] line to the input abundance \( \log \epsilon(C) = 9.5 \), requires that the \( T_{\text{eff}} \) of R CrB and SU Tau be raised about 500 K, and 750 K, respectively. These increases are not only outside the bounds considered acceptable by Asplund et al., but they do not remove the carbon problem for the permitted carbon lines. In addition, they introduce a discrepancy between the N abundance from the N1 and CN lines. Also note that the higher temperatures do not eliminate the abundance difference of 0.7 dex from the 9850 Å and 8727 Å forbidden lines of R CrB and SU Tau. In short, the [C1] lines are part of a now enlarged carbon problem.

### 6 DISCUSSION

Possible solutions to the carbon problem presented by the C1 lines were discussed by Asplund et al. (2000 –see also Gustafsson & Asplund 1996). Two proposed solutions were suggested by Asplund et al. as worthy of further consideration: the \( gf \)-values of the C1 lines are in error, or theoretical model atmospheres are a misrepresentation of the temperature structure of real RCB stars. In addition to commenting on these solutions as they affect the forbidden lines, we examine departures from LTE as they affect line formation. Finally, we discuss hand-crafted model atmospheres including models with a chromosphere (a temperature rise in the outer layers).
Table 2. Photospheric C abundances from C I and [C I] lines

<table>
<thead>
<tr>
<th>Star</th>
<th>Model</th>
<th>$T_{\text{eff}}$</th>
<th>log $g$</th>
<th>$\xi$</th>
<th>C I</th>
<th>C I</th>
<th>[C I] $9850\text{Å}$</th>
<th>[C I] $8727\text{Å}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>RCB stars:</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>XX Cam</td>
<td>7250</td>
<td>0.75</td>
<td>9.0</td>
<td></td>
<td>9.0</td>
<td>8.9</td>
<td>&lt;8.8</td>
<td>8.4</td>
</tr>
<tr>
<td>RY Sgr</td>
<td>7250</td>
<td>0.75</td>
<td>6.0</td>
<td></td>
<td>8.9</td>
<td>9.3</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>UV Cas</td>
<td>7250</td>
<td>0.50</td>
<td>7.0</td>
<td></td>
<td>9.2</td>
<td>9.0</td>
<td>&lt;8.7</td>
<td>8.3</td>
</tr>
<tr>
<td>VZ Sgr</td>
<td>7000</td>
<td>0.50</td>
<td>8.0</td>
<td></td>
<td>9.1</td>
<td>9.1</td>
<td>&lt;8.7</td>
<td>...</td>
</tr>
<tr>
<td>R CrB</td>
<td>6750</td>
<td>0.50</td>
<td>7.0</td>
<td></td>
<td>9.2</td>
<td>9.2</td>
<td>9.1</td>
<td>8.4</td>
</tr>
<tr>
<td>SU Tau</td>
<td>6500</td>
<td>0.50</td>
<td>7.0</td>
<td></td>
<td>8.8</td>
<td>9.0</td>
<td>8.9</td>
<td>8.2</td>
</tr>
<tr>
<td>V482 Cyg</td>
<td>6500</td>
<td>0.50</td>
<td>4.0</td>
<td></td>
<td>8.9</td>
<td>9.6</td>
<td>&lt;8.9</td>
<td>...</td>
</tr>
<tr>
<td>GU Sgr</td>
<td>6250</td>
<td>0.50</td>
<td>7.0</td>
<td></td>
<td>8.8</td>
<td>9.3</td>
<td>&lt;9.0</td>
<td>...</td>
</tr>
<tr>
<td>Standard star:</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\gamma$ Cyg</td>
<td>6100</td>
<td>0.55</td>
<td>3.5</td>
<td></td>
<td>7.9</td>
<td>...</td>
<td>&lt;8.2</td>
<td>7.9</td>
</tr>
</tbody>
</table>

$^a$Asplund et al. (2000) for R CrB stars, the expected C abundance is 9.54 when C/He of 1% models are used; Luck & Lambert (1981) for $\gamma$ Cyg

$b$based on two C I lines near 9850Å

$c$see Section 5.1 of the text

Table 3. Carbon abundances for R CrB for various model atmospheres

<table>
<thead>
<tr>
<th>Model $T_{\text{eff}}, \log g, \xi$</th>
<th>C I</th>
<th>[C I] $9850\text{Å}$</th>
<th>[C I] $8727\text{Å}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>6750,0.5,7</td>
<td>9.2</td>
<td>9.1</td>
<td>8.4</td>
</tr>
<tr>
<td>6500,0.5,7</td>
<td>9.2</td>
<td>8.9</td>
<td>8.2</td>
</tr>
<tr>
<td>7000,0.5,7</td>
<td>9.3</td>
<td>9.3</td>
<td>8.6</td>
</tr>
<tr>
<td>6750,0.0,7</td>
<td>9.3</td>
<td>9.1</td>
<td>8.4</td>
</tr>
<tr>
<td>6750,1.0,7</td>
<td>9.1</td>
<td>9.1</td>
<td>8.4</td>
</tr>
<tr>
<td>6750,1.0,9</td>
<td>8.9</td>
<td>9.1</td>
<td>8.4</td>
</tr>
<tr>
<td>6750,1.0,5</td>
<td>9.3</td>
<td>9.1</td>
<td>8.4</td>
</tr>
</tbody>
</table>

6.1 The gf-values

The gf-values of the [C I] lines are known accurately from theoretical calculations. Confirmation and extension of the carbon problem from the forbidden lines strongly suggests that the gf-values of the permitted lines can not be primarily responsible for the problem.

6.2 Non-LTE Effects

Non-LTE effects in neutral carbon atoms in RCB star atmospheres were evaluated by Asplund & Ryde (1996). Their results indicate that departures from LTE are confined to shallow optical depths. Effects on the C I lines are slight because a line and the continuum are formed deep in the atmosphere between levels in the carbon atom with small and very similar departures from LTE. Gustafsson & Asplund (1996) give the correction to LTE abundance from C I lines as less than 0.1 dex and suggest that +0.02 dex is a typical value. Thus, non-LTE effects can not account for the 0.6 dex typical carbon problem arising from the permitted lines. Departure coefficients given by Asplund & Ryde show that a non-LTE correction to the LTE abundance from a forbidden line is also very small and cannot be the source of the line’s carbon problem. These non-LTE calculations adopt a model computed in LTE.

Carbon problems may be reduced or even eliminated by invoking a chromosphere. A temperature rise at the top of the photosphere is possible even in the absence of mechanical energy deposition. This non-LTE effect sometimes called the Cayrel mechanism (Cayrel 1963) is discussed by Mihalas (1978). To assess the temperature rise for a RCB star due to the Cayrel mechanism, non-LTE effects on the populations of the carbon atom were included in the computation of the model atmosphere. These non-LTE effects on the carbon level-populations were calculated using MULTI (Carlsson 1986), with the modelled carbon atom described in Asplund & Ryde (1996). Realistic background fluxes, computed with the MARCS model atmosphere code, were included. Typical departure coefficients (b=n/n(LTE), where n is the occupation number density) of about b = 1 for the lower C I levels, b = 1.4 for intermediate ones and b = 0.6 for the upper ones were obtained. These figures were found to be characteristic of the depth interval $-4 < \log \tau_{\text{Ross}} < -3$, with much smaller and opposite effects at larger optical depths. These departures from LTE should mainly affect the opacities. C I is also an important electron contributor, but our calculations show only an underionization of less than 10% at shallow depths.

We then brought these results into the MARCS program in order to include their effect on the continuous opacities. This procedure is not fully self-consistent - the models
are still LTE models - but the main effect of the departures on the opacities are modelled in a reasonable way. They essentially diminish the C1 opacities longwards of 6000 Å, to 0.6 times the LTE value, and increase them shortwards of that wavelength to 1.4 times the LTE value. This was applied for the uppermost atmosphere with a decrease corresponding to our non-LTE results at greater depth. A Cayrel effect resulted, as expected, with a heating of the upper layers, but only by about 1-6 K, the maximum occurring at log $\tau_{Ross} = -2$.

Thus, the resultant Cayrel effect is small for C1 in RCB stars with a temperature increase less than 10 K at the top of the photosphere. This micro-chromosphere does not alleviate the carbon problem! The small Cayrel effect is partly linked to the fact that carbon atoms are not totally dominating the scene. In fact at these relevant depths photoionization of carbon atoms only contributes typically 50% of the opacity. So, even drastic changes (at least reductions) to the carbon opacity only lead to moderate changes of fluxes and of the energy balance. Another circumstance is that the ultraviolet flux in the surface layers is so weak that it does not seem to matter very much and is unlike the situation for hot stars and the H\textsc{i} opacity, only a tiny fraction of the flux of typical RCB stars comes at energies higher than the ionization threshold, or even at energies able to photoionize from lower excited C\textsc{i} levels.

6.3 Hand-crafted Models

The fact that the carbon problems defined by the C\textsc{i} and the [C\textsc{i}] lines are very similar from star-to-star across the sample of analysed RCB stars implies that the problems’ resolution cannot depend very sensitively on a star’s individual characteristics. Departures of the atmospheric structure from that predicted by a standard (MARCS) model must be very similar across the sample of stars. If departures are attributable to deposition of mechanical energy, the flux of mechanical energy cannot vary widely across the sample.

That the atmospheric structure is partly the culprit is suggested by the observation that an RCB star in a decline shows a reduced carbon problem. The case of V482 Cyg was discussed in Section 5.8. There, the carbon problem for the C\textsc{i} lines vanished. Note that the C\textsc{i} lines are noticeably stronger in our spectrum than in spectra taken at maximum light. The problem remains for the 9850 Å forbidden line, but this might be the result of overlying emission, as occurred in the case of RY Sgr. (Unfortunately, the 8727 Å line was not on the recorded spectrum.) A similar strengthening of C\textsc{i} lines leading to elimination of the carbon problem but with retention of the problem for the [C\textsc{i}] 9850 Å line was seen by Rao et al. (1999) for R CrB in decline.

6.3.1 The photospheric temperature gradient

The carbon problem may be related to an incorrect representation by the MARCS RCB star models of the temperature gradient through the region of line formation; a steeper/shallower gradient produces a stronger/weaker line. In this connection, it may be noted that Lambert & Rao’s (1994) abundance analysis used unblanketed model atmospheres (Schönbener 1975) and found fair agreement between the observed and calculated equivalent widths of the C\textsc{i} lines. Note that the small carbon problem of 0.2 dex might be attributed to a combination of identifiable uncertainties. The primary factor responsible for the negligible carbon problem is that Schönbener’s unblanketed models have a shallower temperature gradient in the line forming region than the MARCS (line blanketed) RCB star models (Gustafsson & Asplund 1996).

Unblanketed models were calculated with the MARCS code. Our abundance analysis for R CrB using unblanketed MARCS models shows that the carbon problems presented by permitted and forbidden C\textsc{i} lines are similar for unblanketed and blanketed MARCS models. The carbon abundance derived from [C\textsc{i}] 9850 Å line using unblanketed models again agrees that from the C\textsc{i} lines. The large difference between the carbon abundance derived from the [C\textsc{i}] 9850 Å line and the [C\textsc{i}] 8727 Å line using the blanketed models is unaltered by using the unblanketed models. The difference between the unblanketed MARCS and Schönbener’s models reflects differing temperature gradients of the models.

Asplund et al. describe hand-crafted adjustments to the line blanketed models that reduce the photospheric temperature gradient and alleviate the carbon problem presented by the C\textsc{i} lines. A shallower gradient (with respect to a blanketed MARCS model) was qualitatively justified as arising from deposition of mechanical energy whose source may be the sub-photospheric convection zone. Examples of temperature modified models, which alleviate the carbon problem from C\textsc{i} lines, were kindly provided by Martin Asplund (private communication). Our analyses using these models show that the carbon problem vanishes for the 9850 Å [C\textsc{i}] line as it does for the C\textsc{i} lines, but the 8727 Å [C\textsc{i}] line continues to present a problem; the carbon abundance from 8727 Å [C\textsc{i}] line is about 0.7 dex less than that from 9850 Å [C\textsc{i}] line, a difference found also for unmodified line blanketed and unblanketed models.

6.3.2 A chromosphere

Addition of a temperature rise (‘a chromosphere’) at the top of the photosphere may also modify the predicted strengths of the absorption lines in the emergent spectrum. The chromosphere must be crafted by hand without theoretical guidance, but mindful of the discussion in the previous section.

Experiments were made by adding a chromosphere to a MARCS line-blanketed model. Figure 5 shows one experiment appropriate for R CrB itself. A temperature rise is introduced at log $\tau_{Ross} = -1.6$. A flat temperature profile of 7000 K in the interval log $\tau_{Ross} \leq -1.6$ and a LTE spectrum synthesis reproduce the profiles of the observed [C\textsc{i}] 9850 Å and [C\textsc{i}] 8727 Å absorption lines. The two lines now return the same C abundance which is equal to the input abundance, or, the carbon problem vanishes for these lines. The weaker observed C\textsc{i} lines are also reproduced with the input abundance. Very strong C\textsc{i} lines are predicted, as expected for an LTE calculation, to have emission cores. The emission cores might be reduced if the non-LTE effects were taken into consideration. Quite obviously, the chosen chromosphere is not a unique solution. If the onset of the sharp temperature rise is placed as in Figure 5, a chromospheric temperature of 6750 to 7000 K provides acceptable solutions to the carbon problems. Similar chromospheres remove the carbon problems for the other RCB stars. The chromo-
spheric temperature may scale with a star’s effective temperature, but the range of sampled effective temperatures is small.

Addition of a chromosphere not only changes the predicted strengths of the C\(_i\) and [C\(_i\)] lines but also the predicted strengths of other lines and, therefore, the derived abundances. High-excitation lines such as the N\(_i\) line at 8729 \(\text{Å}\) are so weakened by chromospheric emission that the observed absorption line in XX Cam and R CrB cannot be fit whatever the adopted nitrogen abundance. Although careful crafting of the chromospheric temperature profile may alleviate this difficulty, it seems likely that a non-LTE synthesis may be necessary. One may need to adopt spherical rather than a plane-parallel geometry. The chromosphere, if it exists, is likely to be highly-structured and a far cry from the homogeneous plane-parallel (or spherical) layers assumed for the models.

Although the artifice of a chromosphere serves to reduce or even eliminate the carbon problems, its presence presents another puzzle. The energy flux needed to produce the extra energy, the flux of deposited mechanical energy cannot vary widely across the sample.

Uncovering the carbon problem was a surprise. Confidence in model atmospheres was shaken. Extension of the problem to the [C\(_i\)] lines is also disconcerting. In spite of the fact that the RCB stars in several ways are very complex, they have the one property that should make the analysis of C\(_i\) lines in principle simpler than for almost all other stars. This as the fact that the lines are due to the same element is the dominant opacity source and lines and opacity originate from levels of similar character and excitation. This provides a test, that is not offered for other stars. This test, however, fails.

Perhaps, the problems are restricted to very peculiar stars such as the rare RCB stars. But this supposition deserves to be tested. Luminous normal (i.e., H-rich) supergiants of the temperature of the RCB stars are the stars from which chemical compositions of external galaxies are now being derived. Are their atmospheric structures reliably with the hand-crafted chromosphere. Emission from a shell around the star would dilute the photospheric absorption lines. Emission in the 9850 \(\text{Å}\) and 8727 \(\text{Å}\) might weaken the absorption lines such that the emission is not seen but the weakened absorption lines return different carbon abundances. The [C\(_i\)] lines are seen in emission in spectra taken in decline (Rao & Lambert 1993; Rao et al. 1999). The measured fluxes in decline are possibly lower limits to the fluxes at maximum light as the dust cloud responsible for the decline presumably obscures parts of the emitting shell. At maximum light, the cores of the strongest low excitation absorption lines of singly-ionized metals appear distorted by emission (Lambert, Rao & Giridhar 1990; Rao et al. 1999).

It is just such lines which are prominent in the emission line spectrum. The emission line spectrum is of low excitation and does not include the C\(_i\) lines. Therefore, the emitting shell which may contribute to the [C\(_i\)] problem is not a player in the C\(_i\) problem. In this context, it is relevant to recall the case of V482 Cyg observed below maximum light where the carbon problem is absent for the C\(_i\) lines but not for the [C\(_i\)] 9850 \(\text{Å}\) line. Our speculation is that the emitting shell fills in the forbidden line and the photospheric temperature gradient is steeper than at maximum light.

7 CONCLUDING REMARKS

The intriguing carbon problem presented by the RCB stars was discovered by Asplund et al. (2000) in an application of MARCS RCB star models to analysis of C\(_i\) lines. Our contribution to the problem has been to present and analyse observations of the [C\(_i\)] 8727 \(\text{Å}\) and 9850 \(\text{Å}\) lines, also with the MARCS models. The 9850 \(\text{Å}\) line presents a problem of similar magnitude (about 0.6 dex) as C\(_i\) lines, but the carbon problem presented by the 8727 \(\text{Å}\) line is more severe by about 0.7 dex than the original problem (about 0.6 dex).

The fact that the carbon problems defined by the C\(_i\) and the [C\(_i\)] lines are similar from star-to-star across the sample of analysed RCB stars implies that the solution cannot depend on a star’s individual characteristics. Departures of the atmospheric structure from that predicted by a standard (MARCS) model must be similar across the sample of stars. If departures are attributable to deposition of mechanical energy, the flux of deposited mechanical energy cannot vary widely across the sample.

Figure 5. Three \(T - \tau_{\text{Ross}}\) structures are shown - see key at the top of the figure. The solid line shows a MARCS model for \(T_{\text{eff}} = 6750\) and \(\log g = 0.5\). The modified \(T - \tau\) structure with a heated plateau devised by Asplund et al. (2000) to eliminate the carbon problem for C\(_i\) lines is shown by the dotted line. Our model with a chromosphere is shown by the dashed line.
simulated by our codes? Or is there a problem analogous to the carbon problem awaiting discovery?

This research has been supported in part by The Robert A. Welch Foundation through a grant to DLL.

REFERENCES

Lambert, D. L., Rao, N. K., 1994, JAA, 15, 47
Mihalas, D. 1978, Stellar Atmospheres, Freeman & Co.: San Francisco
Moore, Ch. E.: 1972, A Multiplet Table of Astrophysical Interest, NSRDS - NBS 40, Washington