K-band observations of boxy bulges. I. Morphology and surface brightness profiles

M. Bureau, G. Aronica, E. Athanassoula, R.-J. Dettmar, A. Bosma, K. C. Freeman

ABSTRACT
In this first paper of a series on the structure of boxy and peanut-shaped (B/PS) bulges, K-band observations of a sample of 30 edge-on spiral galaxies are described and discussed. K-band observations best trace the dominant luminous galactic mass and are minimally affected by dust. Images, unsharp-masked images, as well as major-axis and vertically-summed surface brightness profiles are presented and discussed. Galaxies with a B/PS bulge tend to have a more complex morphology than galaxies with other bulge types, more often showing centered or off-centered X structures, secondary maxima along the major-axis and spiral-like structures. While probably not uniquely related to bars, those features are observed in three-dimensional N-body simulations of barred discs and may trace the main bar orbit families. The surface brightness profiles of galaxies with a B/PS bulge are also more complex, typically containing 3 or more clearly separated regions, including a shallow or flat intermediate region (Freeman Type II profiles). The breaks in the profiles offer evidence for bar-driven transfer of angular momentum and radial redistribution of material. The profiles further suggest a rapid variation of the scaleheight of the disc material, contrary to conventional wisdom but again as expected from the vertical resonances and instabilities present in barred discs. Interestingly, the steep inner region of the surface brightness profiles is often shorter than the isophotally thick part of the galaxies, itself always shorter than the flat intermediate region of the profiles. The steep inner region is also much more prominent along the major-axis than in the vertically-summed profiles. Similarly to other recent work but contrary to the standard ‘bulge + disc’ model (where the bulge is both thick and steep), we thus propose that galaxies with a B/PS bulge are composed of a thin concentrated disc (a disc-like bulge) contained within a partially thick bar (the B/PS bulge), itself contained within a thin outer disc. The inner disc likely formed secularly through bar-driven processes and is responsible for the steep inner region of the surface brightness profiles, traditionally associated with a classic bulge, while the bar is responsible for the flat intermediate region of the surface brightness profiles and the thick complex morphological structures observed. Those components are strongly coupled dynamically and are formed mostly of the same (disc) material, shaped by the weak but relentless action of the bar resonances. Any competing formation scenario for galaxies with a B/PS bulge, which represent at least 45 per cent of the local disc galaxy population, must explain equally well and self-consistently the above morphological and photometric properties, the complex gas and stellar kinematics observed, and the correlations between them.

Key words: galaxies: bulges – galaxies: evolution – galaxies: formation – galaxies: photometry – galaxies: spiral – galaxies: structure

1 INTRODUCTION
Since the work of [Davies et al. 1983], bulges have largely been considered as low-luminosity elliptical galaxies, sug-
gesting a rapid formation dominated either by mergers and accretion of external material (e.g. Kauffmann 1993, Kauffmann, Charlot & White 1996; but see also Bekki & Chiba 2003; Aguerri, Balcells & Peletier 2001) or possibly by dissipative gravitational collapse (e.g. Eegens-Lynen-Bell & Sadler 1962; Carlberg 1984a,b; Steinmetz & Muller 1995; but see also Noguchi 1998, 1999).

Over the last decade, however, there has been much criticism of this idea. In particular, the realization that most bulges have an inner surface brightness profile shallower than the expected $r^{-4/3}$ law (e.g. Andreidakis, Peletier & Balcells 1993, de Jong 1996, MacArthur, Courteau & Holtzman 2003; Balcells et al. 2003) argues against both mechanisms (e.g. Horth & Madsen 1993, Aguerri et al. 2003). Alternative models where bulges grow secularly (i.e. over a long timescale and in relative isolation) have also been developed and studied, many of them bar-driven (e.g. Pfenniger & Norman 1990, Friedli & Benz 1993, Normand, Sellwood & Hasan 1992). and much observational data support them (e.g. Wozniak et al. 1995, Erwin & Sparkes 2002).

Of primary concern here, several pieces of evidence point to the identification of most boxy and peanut-shaped (B/PS) bulges in edge-on spiral galaxies with the bars of barred spirals. In N-body simulations, whenever a disc galaxy forms a bar, a B/PS bar/bulge develops soon after. This process was studied first by Combes & Sanders (1981), later on by Combes et al. (1990) and Raha et al. (1991), and more recently by Mihos et al. (1994, Athanassoula & Misiriotis 2002, Athanassoula 2002, 2003, 2005 and Martinez-Valpuesta, Shlosman & Heller 2005).

The observed incidence of B/PS bulges is consistent with expected if they are associated with relatively strong bars. Recent work by Lutticke, Dettmar & Pohlen (2000) demonstrates that 45 per cent of all bulges are B/PS, while amongst those the exact shape of the bulge depends mainly on the viewing angle to the bar. As shown by the numerical simulations, true peanuts are bars seen side-on, i.e. with the major-axis of the bar roughly perpendicular to the line-of-sight. For less favourable viewing angles, the bulge/bar looks boxy, and if the bar is seen end-on it looks almost spherical. Stronger bars also lead to more prominent peanut shapes, as demonstrated observationally (e.g. Lutticke et al. 2000b) and theoretically (e.g. Athanassoula & Misiriotis 2002, Bureau & Athanassoula 2003).

The kinematics of discs harboring a B/PS bulge, as measured from both ionized-gas emission lines and stellar absorption lines, show the behaviour expected of barred spirals viewed edge-on. This has been demonstrated by the observations of Kuijken & Merrifield (1993, Merrifield & Kuijken 1994, Bureau & Freeman 1995, 1996 and Chung & Bureau 2004), and by the modeling of Athanassoula & Bureau (1999) and Bureau & Athanassoula (1999, 2005). Similar tests are now also available for face-on bars (Debattista et al. 2003).

There is thus ample evidence that edge-on galaxies with a B/PS bulge are simply barred disc galaxies, and that the B/PS bulges themselves represent the thickest parts of the bars (see Athanassoula 2005 for a review of all arguments). Yet there has been little convincing evidence for this from surface photometry alone, the best work being that of Lutticke et al. (2000a,b). Early studies of non-spheroidal bulges used mostly optical images (e.g. Jarvis 1987, Shaw & Gilmore 1987; Shaw 1993), and the interpretation was often hampered by the large amount of extinction. As we will show in this paper, it truly takes the combination of N-body simulations and orbit studies with near-infrared (NIR) images to derive direct and convincing photometric evidence relating B/PS bulges and bars.

Skokos, Patsis & Athanassoula (2002, 2003) and Patsis, Skokos & Athanassoula (2003a) studied the orbital structure of three-dimensional (3D) bars exhaustively (but see also Pfenniger 1984, Pfenniger & Friedli 1990). They find families of orbits which can only provide the backbone of the boxy and peanut shapes, but can also cause local enhancements within the disc itself. Since extinction is far less important in the NIR, K-band images are the ideal tool to study the morphology of galaxies with a B/PS bulge, to look for similarities with barred orbital structures.

In this paper, the first of a series, we study a sample of 30 edge-on spiral galaxies, most of them with a B/PS bulge, for which we have obtained high-quality K-band images. Much complementary data exist for this sample (e.g. Bureau & Freeman 1997, 1999, Chung & Bureau 2004, Bureau & Chung 2006), but the primary goal here is to study the morphology of the B/PS structures and their host discs, similarly to Lutticke et al. (2000a,b). As we shall see, we find several features resembling closely those expected of 3D bars: bulges with X shapes, secondary disc enhancements, inner rings, etc. A second paper discusses scaleheight variations in the same galaxies (Athanassoula, Aronica & Bureau 2006, hereafter Paper II; but see also Aronica et al. 2003, 2004).

We present our sample in § 2 discuss the observations and data reduction in § 3 and then describe and discuss the resulting images and surface brightness profiles in § 4 and 5 respectively. We examine the direct consequences of our results in § 6 and conclude briefly in § 7.

### 2 Sample Description

We study here the sample of Bureau & Freeman (1997, 1999) and Chung & Bureau (2004), which was drawn from the catalogs of galaxies with a B/PS bulge of Jarvis (1987), Shaw (1987), and de Souza & da Amoia (1987), as well as from galaxies with extreme axial ratios ($a/b \geq 7$; Karachentsev, Karachentseva & Parnovskiy 1993). We include galaxies with extended as well as no or confined ionized-gas emission, even though the latter were not specifically studied by Bureau & Freeman (1994). The sample thus consists of 30 edge-on spirals, all of which are visible from the south ($\delta \lesssim 15^\circ$) and have a bulge larger than 0.6 in diameter (to be able to identify bars kinematically in moderate seeing). 24 of the bulges were classified as having a B/PS bulge by Bureau & Freeman (1999) and 6 constitute a ‘control’ sample and have varied morphologies.

Basic galaxy properties are listed in Table 1, along with a classification of the bulges’ shape from Lutticke et al. (2000b). Although still relying on a visual inspection of optical contour plots, Lutticke et al. (2000a,b) homoge-
Table 1. Galaxy sample.

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**B/PS bulges**

NGC 128     | 00 29 15.1 | +02 51 50 | SB pec   | 1.0   | 12.7    | -21.4   | 2.81    | 4228    |
ESO 151-G004| 00 56 07.3 | -53 11 28  | SB 0     | 1.0   | 14.7    | -20.4   | 1.31    | 7456    |
NGC 1381    | 03 36 31.7 | -53 17 43  | SA0      | 2.0   | 12.7    | -19.2   | 2.63    | 1776    |
NGC 1596    | 04 27 37.8 | -55 01 37  | SA0      | 4.0   | 12.0    | -19.3   | 3.89    | 1509    |
NGC 1886    | 05 21 48.2 | -23 48 36  | Sab      | 1.0   | 13.8    | -19.2   | 3.23    | 1737    |
NGC 2310    | 06 53 53.8 | -40 51 46  | S0       | 2.0   | 12.6    | -18.6   | 4.16    | 1187    |
ESO 311-G012| 07 47 34.0 | -41 27 07  | S0/a?    | 2.0   | 12.4    | -20.0   | 3.71    | 1130    |
NGC 2788A   | 09 02 40.2 | -68 13 38  | Sb       | 1.0   | 13.6    | -21.5   | 2.88    | 4056    |
IC 2531     | 09 59 55.4 | -29 37 02  | Sb       | 1.0   | 12.9    | -21.6   | 6.76    | 2474    |
NGC 3203    | 10 19 34.4 | -26 41 53  | SA(r)0+  | 3.0   | 13.0    | -19.9   | 2.81    | 2410    |
NGC 3390    | 10 48 04.4 | -31 32 02  | Sb       | 2.0   | 12.8    | -21.5   | 3.46    | 3039    |
NGC 4469    | 12 29 28.0 | +08 44 59  | SB(s)0/a?| 1.0   | 12.4    | -17.8   | 3.46    | 576     |
NGC 4710    | 12 49 38.9 | +15 09 57  | SA(r)0   | 1.5   | 11.9    | -19.8   | 4.89    | 1324    |
PGC 4931    | 13 01 49.5 | -08 20 10  | Sbc      | 1.0   | 14.2    | -21.1   | 2.81    | 3804    |
ESO 443-G042| 13 03 29.9 | -29 49 36  | Sb       | 1.0   | 13.9    | -20.6   | 2.88    | 2912    |
NGC 5746    | 14 44 55.9 | +01 57 17  | SA(R)rs/b?| 1.0   | 11.4    | -21.8   | 6.91    | 1720    |
IC 4767     | 18 47 41.6 | -63 24 20  | S pec    | 1.0   | 14.3    | -19.5   | 1.51    | 3544    |
NGC 6722    | 19 03 39.6 | -64 53 41  | Sb       | 1.0   | 13.5    | -22.2   | 2.88    | 5749    |
NGC 6771    | 19 18 39.6 | -60 32 46  | SA(r)0+  | 1.0   | 13.6    | -20.5   | 2.34    | 4221    |
ESO 185-G053| 20 03 00.4 | -55 56 53  | SB pec   | 2.0   | 14.3    | -20.0   | 1.23    | 4475    |
IC 4937     | 20 05 17.9 | -56 19 20  | Sb       | 1.0   | 14.8    | -18.6   | 1.86    | 2337    |
ESO 597-G036| 20 48 15.0 | -19 50 58  | SB 0 pec | 1.0   | 15.2    | -20.7   | 0.87    | 8694    |
IC 5096     | 21 18 21.8 | -63 45 42  | Sb       | 4.0   | 13.6    | -20.7   | 3.16    | 3142    |
ESO 240-G011| 23 37 50.5 | -47 43 37  | Sb       | 4.0   | 13.4    | -21.0   | 4.89    | 2842    |

**Control sample**

NGC 1032    | 02 39 23.6 | +01 05 38  | S0/a     | 4.0   | 12.7    | -20.7   | 3.46    | 2722    |
NGC 3957    | 11 54 01.5 | -19 34 09  | SA0      | 3.0   | 13.0    | -19.0   | 3.09    | 1686    |
NGC 4703    | 12 49 18.9 | -09 06 30  | Sb       | 4.0   | 14.0    | -21.1   | 2.45    | 4458    |
NGC 5084    | 13 20 16.8 | -21 49 38  | S0       | 4.0   | 11.5    | -20.9   | 10.71   | 1725    |
NGC 7123    | 21 50 46.4 | -70 19 59  | Sa       | 4.0   | 13.6    | -20.3   | 2.51    | 3737    |
IC 5176     | 22 14 55.3 | -66 50 56  | SAB(s)bc?| 4.0   | 13.5    | -19.6   | 4.36    | 1746    |

Notes: All parameters are from LEDA (Lyon-Meudon Extragalactic Database), except the morphological type (Jarvis 1980; de Souza & dos Anjos 1987; Shaw 1987; Karachentsev et al. 1993) and bulge type (Lütticke et al. 2000a). Bulge types 1–3 represent B/PS bulges and bulge type 4 spheroidal bulges. ESO 151-G004’s redshift and corrected absolute $B$ magnitude are from Chung & Bureau (2004).

3 OBSERVATIONS AND DATA REDUCTION

The observations were carried out at the 2.3m telescope of Siding Spring Observatory over 8 runs from 1996 January to 1997 March, for a total of 24 nights. Mounted at the f/18 Cassegrain focus, the CASPIR1 (Cryogenic Array Spectrometer/Imager; McGregor et al. 1994) instrument was used in direct imaging mode with a 0.5 arcsec pix$^{-1}$ scale. Its 256 × 256 pix$^2$ SBRC InSb CCD (30μm pixels) then yields a 12 arcsec × 12 arcsec field-of-view. All observations were obtained with a narrow $K$-band filter, $K_0$, of central wavelength 2.165μm and width 0.33μm. The readout noise with the double sample readout method used (relative sampling) was about 60 e and the dark current $\lesssim 30$ e s$^{-1}$ over most of the array, although there are a number of hot pixels. For 5 s integrations (see below) and typical background brightness conditions (12.0–13.0 $K_0$ mag arcsec$^{-2}$), our observations are always background limited.

1 http://www.mso.anu.edu.au/observing/2.3m/CASPIR/
For all observations except biases and standard stars, the integration time for a single frame was 60 s, resulting from the average of twelve 5 s exposures (individual exposures were not saved). The total integrations were built-up using a script taking 5 frames of each field, each dithered by 4′′ (8 pixels; the biggest bad pixel patch) parallel or perpendicular to the galaxy major-axis, and interspersed by sky frames. The lot was bracketed by dark and bias frames and repeated typically 7 times, for a total on-source integration time of 35 min per field, slightly less on sky. For a galaxy requiring mosaicing of 2 fields, for example, the sequence was: bias, dark, sky, field_1 [dith 0], field_2 [dith 0], sky, field_1 [dith 1], field_2 [dith 1], sky, field_1 [dith 2], ... This ensures sufficiently fast sampling of the sky and allows to average out bad pixels when the frames are combined. Using the instrument rotator, the CCD rows were always aligned (and centered) on the galaxy major-axis, facilitating mosaicing. Mosaics of 1 to 5 fields (with 25 per cent overlap) were required depending on the galaxy, and all mosaics include plenty of sky area around the target.

The data reduction was mostly carried out within the CASPIR IRAF (Image Reduction and Analysis Facility) package, following as closely as possible the procedures recommended in the CASPIR manual. Dark, sky, and object frames were first bias-subtracted and linearized, and the skies and objects dark-subtracted. Sky and object frames were then flatfielded and the sky subtracted from the object frames. All the frames of a given object were then combined together with proper offsets to form the final image. The bias and dark frames used for subtraction were nightly (or fractions of a night if deemed necessary) medians of all such frames. For sky subtraction, running medians of a few sky frames were used, time-centered around the object frame being reduced. A flat was obtained every night from the difference of domes frames taken with an incandescent lamp. F or most objects, a mosaic was created for each field when necessary (imperfect sky subtraction).

A photometric zero-point on the Carter SAAO system (Carter & Meadows 1992) was derived nightly by observing IRIS (Infrared Imaging Spectrograph) photometric standards throughout the night. The K-band atmospheric extinction correction coefficient was also calculated each night by following a few standard stars over a range of airmasses (typical value ≈ 0.10 mag airmass^{-1}). The galaxy images were also corrected for Galactic extinction, although this is negligible for most objects. When frames were obtained under less than photometric conditions, their fluxes were first scaled to the level of photometric ones before being combined. The final images generally have a background level slightly different from zero, and this was again corrected by subtracting a constant.

4 IMAGES

4.1 Morphological features

Figures 1 and 2 show, respectively, our data for galaxies with a B/PS bulge and for those of the control sample. Each panel contains, from top to bottom, first a Digitized Sky Survey (DSS) image of the galaxy (for comparison), second our K-band image and third an unsharp-masked Kn-band image (see below), all spatially registered. The surface brightness profiles which follow will be discussed in the next section (§5). As our main interest here is the galaxy bulges, we only show the inner parts of our mosaics, which generally extend up to and beyond the large-scale discs.

Comparing the DSS and Kn-band images, the advantages of working in the NIR become obvious. Most importantly, dust extinction within the galaxies, when present, is drastically reduced (see, e.g., IC 2531 and NGC 7123), although it remains non-negligible in a few objects (e.g. NGC 1032). Generally, the B/PS and associated features are also sharper at Kn-band, although this is hard to quantify. For a similar reason, unsharp-mask, which enhances local extrema, is ideal to highlight readily apparent morphological features and reveal weaker ones.

The unsharp-masked images of Figures 1 and 2 were obtained by median-filtering the Kn-band images, that is by replacing the value of each pixel by the difference between it and that of the median within a centered circular aperture. The size of the aperture was held fixed across each image, but it was varied from galaxy to galaxy to best highlight the features of interest here: X shapes, secondary disc enhancements, ... Bigger apertures yield unsharp-masked images more similar to the original Kn-band image, while smaller apertures enhance smaller scale features such as central discs, spiral arms and remaining dust lanes. The apertures used in Figures 1 and 2 are thus tightly correlated to the characteristic size of the morphological features enhanced. We also note that the unsharp-masked images were obtained from Kn-band images with the foreground stars interpolated over. This process is not perfect, however, and can create artefacts in the unsharp-masked images, especially near the equatorial plane of the galaxies where the luminosity gradients are steep. One must thus be careful when interpreting the results (see, e.g., NGC 2788A and NGC 6722).

As the same morphological features are present in many galaxies, we provide a generic description of those features below and list the specific features observed in each object in columns 2–7 of Table 2. We note however that, for some objects, some features are best seen in images unsharp-masked on a spatial scale different from that shown in Figures 1 and 2. To create Table 2, we thus analysed images unsharp-masked with circular apertures of radii ranging from 2.5 arcsec (5 pixels) to 17.5 arcsec (35 pixels).

Off-centered X (OX). Many galaxies show an X shape structure in the bulge region, usually most easily visible in the unsharp-masked Kn-band image. In many cases, however, if one follows the branches of the X, the ridges do not cross in the centre of the galaxy but rather fall short of it. This leads to a ‘> <’ feature rather than a true ‘> <’, which we dub an off-centered X (OX). Excellent examples are NGC 128 and NGC 6771, but there are many.

Centered X (CX). In a comparable number of cases, the branches of the X do cross in the center of the galaxy,
Figure 1. Images and surface brightness profiles of the galaxies with a B/PS bulge. From top to bottom, each panel shows first a DSS image of the galaxy, second our $K_n$-band image, third an unsharp-masked $K_n$-band image and last major-axis (fainter) and summed (brighter) surface brightness profiles, all spatially registered. Contours for the $K_n$-band image are spaced by 0.5 mag arcsec$^{-2}$ and the faintest contour is indicated in the bottom-right corner of the panel.
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Figure 2. Same as Figure 1 but for the control sample galaxies.
and we dub this feature a centered X (CX). Although the transition between off-centered and centered X is somewhat ill-defined, good examples of centered X features include NGC 1381 and IC 2531.

**Minor-axis extremum (ME).** Many galaxies show a rather narrow and elongated local maximum along the minor-axis, which we dub minor-axis extremum (ME). They are generally only visible in $K_n$-band images unsharp-masked on a smaller scale than that presented in Figures 1–2, so we show the cases of NGC 1381 and NGC 3203 in Figure 3 for illustrative purposes.

**Secondary maxima (SM).** Many galaxies display local maxima away from the center along the major-axis, which we dub secondary maxima (SM). While sometimes readily visible in the original $K_n$-band image (e.g. NGC 4710 and ESO 597-G036), most are easily seen only in the unsharp-masked images (e.g. ESO 151-G004 and NGC 3203).

**Spiral arms (SA).** A number of galaxies display narrow and elongated local maxima which are not aligned with either axes but rather are typically slightly offset from the major-axis, and sometimes even partially wind around the bulge forming a pseudo-ring. Those features are generally bisymmetric and are again most easily visible in the unsharp-masked $K_n$-band images. We dub them spiral arms (SA) for obvious reasons. Good examples include NGC 4469 and NGC 5746 (but see also the much weaker example of NGC 3203 in Fig. 3).

In Table 2, we finally give an indication of the inclination of the galaxies, labeling them either inclined (I), when the galaxy is obviously more than a few degrees from exactly edge-on, quasi edge-on (Q), when the galaxy is apparently away from edge-on by at most a few degrees, and exactly edge-on (E), when there is no indication at all that...
Table 2. Galaxy features.

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<td>NGC 7123</td>
<td>CX</td>
<td>ME</td>
<td>E</td>
</tr>
<tr>
<td>IC 5176</td>
<td>ME</td>
<td>SA?</td>
<td></td>
</tr>
</tbody>
</table>

Notes: For the images, the features listed were identified on images unsharp-masked on scales ranging from 2.5 to 17.5 arcsec. OX: off-centered X feature, CX: centered X feature, ME: minor-axis extremum, SM: secondary maxima along major-axis, SA: spiral arms, I: inclination. For the inclination, the galaxies are classified as edge-on (E), quasi-edge-on (Q), or inclined (I). For the surface brightness profiles, we list the number of clearly distinct regions. MA: major-axis profile, SU: vertically-summed profile. The letter ‘F’ denotes a flat intermediate region.

4.2 Incidence of morphological features

The most striking result from Figures 1–2 is that the morphology of the galaxies with a B/PS bulge is much more complex than that of the control sample galaxies, in the sense that galaxies with a B/PS bulge contain more of the morphological features described in §4.1 (and best-revealed by unsharp-masking; see Table 2). Indeed, at least 50 per cent of galaxies with a B/PS bulge possess an off-centered X feature and at least 38 per cent a centered one, while only 33 per cent of the control galaxies have either. Similarly, 88 per cent of the galaxies with a B/PS bulge have secondary maxima along the major-axis, while only 33 per cent of the control galaxies do. For the presence of spiral arms, the ratios become, respectively, at least 38 per cent and perhaps nil. Only for the minor-axis extrema are those percentages inverted, with at most 58 per cent of galaxies with a B/PS bulge possessing one but all control galaxies. Those statistics thus suggest that both centered and off-centered X features, secondary major-axis maxima and spiral arms are preferentially associated with B/PS bulges, while only minor-axis extrema are found preferentially in other bulge types. In fact, the contrast between the main and control samples would be even greater if, as we will argue in §6.5, the control sample...
was not contaminated by weak B/PS bulges (NGC 3957 and NGC 4703).

We also note that, given the large number of secondary maxima observed at K-band (in fact, more than in the optical images), it is unlikely that they are due to the presence of obscuring dust filaments inclined with respect to the equatorial plane, as was suggested by various authors (e.g. Sandage 1961, Whitmore & Bell 1988).

4.3 Origin of morphological features and the orbital structure of bars

Although the accretion of external material can give rise to X-shaped features (e.g. Binney & Petros 1983; Hernquist & Quinn 1988; 1989; Whitmore & Bell 1988), these are generally centered and it is unlikely that any long-lasting off-centered X could be produced. Since we observe a majority of off-centered X in our sample, and since there is no obvious discontinuity between the two, it is more attractive to look for a unique mechanism which can simultaneously explain both centered and off-centered stable X features.

Athanassoula (2003) unsharp-masked a high-quality N-body simulation of a barred disc. The results provide an essentially perfect match to the secondary maxima discussed above and, depending on the viewing angle to the bar, to the off-centered and centered X features, suggesting a link between those features and edge-on bars. We discuss below the possible origin of this link through the orbital structure of barred discs and then discuss some caveats.

The most important orbit families in 3D bar models are those of the $x_1$ tree, all elongated parallel to the bar and located within corotation. They include the $x_1$ family itself (restricted to the equatorial plane) and many other families bifurcating from the $x_1$ at vertical resonances of increasing energies. Sellwood & Wilkinson (1980) provide a general introduction, but we adopt here the notation of Skokos et al. (2002) and call them $x_1v_1, x_1v_2, x_1v_3, \text{etc.}$ The morphological features identified in §4.1 can then be reproduced by superposing orbits of the appropriate shapes, as done by Patsis et al. (2002), Patsis & Xilouris (2006) recently explored very similar ideas, with identical conclusions.

The centered and off-centered X features, in particular, can both arise out of families extending out of the equatorial plane in the vertical direction. The $x_1v_2$ family (also called banana orbits; e.g. Pfenniger & Friedli 1991), bifurcating from the $x_1$ family at the $2:1$ vertical resonance, has the largest vertical extent and is shaped like a smile or a frown. As shown by the orbit superpositions of Patsis et al. (2002), $x_1v_1$ orbits seen side-on have a morphology entirely consistent with an off-centered X. Depending on the model (i.e. mass distribution and pattern speed) and viewing angle, centered X features can also be created.

The $x_1v_3$ and $x_1v_4$ orbit families (bifurcating from the $x_1$ family at the $3:1$ vertical resonance and shaped like a ‘c’) and the $x_1v_5$ orbit family (bifurcating at the $4:1$ resonance and shaped like a ‘w’) can also give rise to centered or off-centered X features, depending on the model and viewing angle (see, again, Patsis et al. 2002). This is also the case for the $z_1v_8$ family, which does not bifurcate from the $x_1$ orbits but is morphologically similar to the $x_1v_4$ family, although it was present in only one model.

The orbit superpositions of Patsis et al. (2002) further show that the orbit families described above (as well as higher order families such as $x_1v_7, x_1v_9$ and $x_1v_9$) can give rise to a number of maxima along the major-axis, similar to the secondary maxima identified in §4.1. Those maxima generally occur at larger radii than the X features and near (but within) the ends of the bar, often where orbits have loops. An alternative and more straightforward explanation is that the secondary maxima observed are simply the edge-on projections of inner rings (edge brightening), known to exist in a large fraction of barred spiral galaxies and located at the end of the bar (e.g. Kormendy 1974; Buta 1994). Inner rings naturally form in gas-rich discs under the influence of rotating bars (e.g. Schwarz 1984; Byrd et al. 1994), but they also develop in purely dissipationless simulations (e.g. Athanassoula & Misiriotis 2002; Bureau & Athanassoula 2003). Either way, both mechanisms support a relationship to bars, and both mechanisms may coexist. Accretion scenarios producing centered X features, on the other hand, will generally not produce secondary maxima along the major-axis (even less so systematically located outside of the X feature itself), making them even less attractive.

The orbit superpositions of Patsis et al. (2002) show that the barred orbit families considered above can also give rise to the minor-axis extrema described in §4.1, especially when seen end-on. The $z_1$ orbits are particularly interesting in this respect but were little studied by Skokos et al. (2002) and deserve more attention. This is encouraging, but the fact that all control sample galaxies also show a minor-axis extremum suggests that this feature is not related to B/PS bulges (and thus presumably to bars). We do not properly understand the origin of this feature.

As we mentioned in §4.1, because of the morphological similarities, we interpret the local maxima elongated parallel to but slightly offset from the major-axis of the galaxies as spiral arms. If this association is true, which we have no reason to doubt, then obviously spiral arms can only be detected in galaxies that are not perfectly edge-on. In fact, we have used the presence of spirals to quantify the inclination of the galaxies in Table 2. Since we have no other independent and reliable way to estimate the inclination of the galaxies precisely (as we do not know the intrinsic thickness of the discs), we can not guarantee that the inclination distribution is similar between the sample of galaxies with a B/PS bulge and the control sample. Thus, a reliable comparison of the incidence of spiral arms between the two samples is not possible, and we do not discuss this feature at length.

Suffice it to say that the apparently slightly higher incidence of spiral arms in galaxies with a B/PS bulge is consistent with the suggestion that B/PS bulges are simply thick bars viewed edge-on. Indeed, bars are very good at driving (grand design) spiral patterns (e.g. Athanassoula 1980; Thielemann & Wolf 1984; Elmegreen & Elmegreen 1989) and spiral arms are preferentially found in barred galaxies (e.g. Elmegreen & Elmegreen 1989; Buta et al. 2005). Furthermore, for the few galaxies where a (partial) inner ring with outer spiral arms is visible (e.g. NGC 5746), the ring is systematically located just at the end of the flat intermediate region of the major-axis surface brightness profiles (see §4.5 all such galaxies have an intermediate region, and it is flat in all cases but one). If, as we argue in §4.5 the end of the intermediate surface brightness profile region also marks the end of the
bar, then this trend is entirely consistent with observations of more face-on barred spiral galaxies (e.g. Kormendy 1979; Buta 1993) and the expectations from theory (e.g. Schwarz 1983, 1984; Byrd et al. 1994).

On a different note, we must point out that while the (barred) picture proposed above is attractive and supported by much independent evidence (see § II), it may not be unique. Indeed, as simple tests show, surface brightness distributions created from nested rectangular isophotes can also give rise to centered and off-centered X-shapes when unsharp-masked, independently of their origin. As axisymmetric boxy distributions can in principle also exist (e.g. Binney & Petrucci 1982; May, van Albada & Norman 1983; Rowles 1988), it is possible that bars are not uniquely related to the complex features observed. However, as the N-body simulations and orbit properties discussed above show, barred models are consistent with the unsharp-masked features in much details, while this remains to be shown for other models.

5 SURFACE BRIGHTNESS PROFILES

5.1 Profile features

In addition to the DSS, Kν-band and unsharp-masked Kν-band images, the panels of Figures 12 also contain two surface brightness profiles. The fainter ones represent the major-axis surface brightness profiles, extracted using a constant position angle, an approximation that is valid in the inner parts of all galaxies and in the outer parts of all but a few (e.g. NGC 128, NGC 6722, NGC 6771 and ESO 597-G036). The brighter surface brightness profiles were obtained by summing the data in the vertical direction, as if the galaxies were infinitely thin, again assuming a constant position angle for the major-axis. In both cases, the Kν-band images with the foreground sources interpolated over were used, and the vertical profiles were summed until the noise level of the image was reached. While the noise level varies slightly from object to object, such a limit is sufficient for our purposes.

While our summed surface brightness profiles are not strictly equivalent to azimuthally-averaged surface brightness profiles in more face-on galaxies, or to cuts at a given position angle (i.e. they are also integrated along the line-of-sight through the edge-on discs), a comparison is nevertheless possible and useful. Indeed, because the luminosity of discs generally decreases rapidly with radius, the light at any given position in our profiles is still dominated by that near the same (cylindrical) radius. Our profiles can also be compared to those in Lattica et al. (2000), who showed profiles (cuts) parallel to but offset from the major-axis for a number of edge-on spiral galaxies.

In columns 8 and 9, Table 2 lists the number of distinct regions present in the surface brightness profiles of each galaxy, each region being typically separated by a clear radial break (although one must be cautious with badly interpolated foreground sources). From the surface brightness profiles of face-on galaxies, we would generally expect each profile to show only two distinct regions if the galaxy is axisymmetric. A first region at small radii, generally associated with the bulge, where the profile is steep (generally steeper than an exponential profile), and a second region at larger radii, generally associated with the disc, where the profile is exponential. However, such two-region profiles are rare in our data, especially along the major-axis (but see NGC 1596 and NGC 1032). Most profiles show an additional region at intermediate radii, where the profile is shallower than the outer disc, even often flat or slightly rising with radius (denoted by the letter ‘F’ in Table 2, e.g. ESO 151-G004 and NGC 2310). Such three-region surface brightness profiles are often referred to as Freeman Type II profiles in the literature (Freeman 1970). A few galaxies also show two distinct intermediate regions, one shallow following the steep central component and one steep preceding the outer exponential disc, for a total of four regions (e.g. NGC 2788A and NGC 4710). We do not discuss (or count) here apparent radial breaks in the outer discs, since they are presumably only weakly related to the bulge structure. Figure 4 shows a cartoon version of the various types of profiles identified in Table 2.

5.2 Incidence of profile features

Again, the most obvious and basic result from Figures 12 is that the surface brightness profiles of the galaxies with a B/PS bulge are much more complex than those of the control sample galaxies, especially along the major-axis, in the sense that the surface brightness profiles of galaxies with

2 We remind the reader that the surface brightness profile of an edge-on exponential disc is not exactly linear in a (logarithmic) surface brightness vs. projected radius plot, but slightly curved (see, e.g. Pohlen et al. 2001).
a B/PS bulge typically contain a larger number of distinct regions separated by clear radial breaks (see Table 2). Indeed, along the major-axis, 96 per cent of the galaxies with a B/PS bulge have a surface brightness profile with 3 or more regions (Freeman Type II profile), 78 per cent of which have a flat intermediate region, while the corresponding fractions for the control sample are, respectively, only 50 and 66 per cent. The statistics are not as sharply contrasted when the profiles summed along the vertical axis are considered, but the differences remain significant. For example, 58 per cent of the galaxies with a B/PS bulge have a summed surface brightness profile with a least 3 regions, while only 17 per cent of the control galaxies do. Those fractions are in rough agreement with those of Lüticke et al. (2000), who carried out a similar analysis. As for the morphological features, it thus appears that complex surface brightness profiles are preferentially associated with B/PS bulges. And again, the contrast between the main and control samples would be sharper if the control sample had been better selected (i.e. if NGC 3957 and NGC 4703 has been excluded; see § 5.3).

5.3 Origin of profile features and relationship to bars

The fact that practically all galaxies with a B/PS bulge have a major-axis surface brightness profile with an intermediate, often flat region is particularly important, as it is inconsistent with an axisymmetric bulge and exponential disc model, no matter how steep the bulge profile is (i.e. no matter what the Sersic n index of the bulge is). There is thus apparently a third photometric or morphological 'component' dominating at intermediate radii. We argue below, however, that in such Freeman Type II profiles, both the steep inner region and the flat intermediate region are ultimately caused by a single bar viewed edge-on, without the necessity of a classic bulge.

Bureau & Athanassoula (2003) followed the temporal evolution of the major-axis surface brightness profile in a large number of barred N-body simulations viewed edge-on. They showed that bar formation and evolution within an (initially) exponential disc is associated with the buildup and continued growth of a dense central region, resulting in a central peak (which would normally be identified with a bulge) and with the formation and gradual flattening of an intermediate region, in addition to the outer exponential disc (see also Athanassoula & Misiriotis 2002, Athanassoula 2003). The central peak extends to 1–1.5 original disc scalelengths while the intermediate region extends to the end of the bar, well beyond the central peak and the thickest part of the bar. The existence of three surface brightness profile regions is thus totally normal (and expected) in barred galaxies, but it remains unexplained in classic bulge models, unless there is an additional bar. In fact, four-region surface brightness profiles are often produced in barred models (e.g. Fig. 3 in Bureau & Athanassoula 2003).

As expected from the elongated boxy/peanut shape of the bar, the ratio of the length of the thickest part of the bar (or that of the central peak) to that of the flat intermediate region increases in N-body simulations as the viewing angle increases (from end-on to side-on). This trend is also present in our data, where the ratio of the length of the thickest part of the bulges (or that of the central peak) to that of the intermediate region is generally larger in peanut-shaped bulges than in boxy ones (compare, e.g., NGC 5746 and NGC 6722 to NGC 1381 and IC 5096). There is however much variety, as many strong peanut-shaped bulges have a small ratio (see, e.g., NGC 1886 and ESO 443-G042), and likely many causes for it. First, as far as the surface brightness profile and bar shape are concerned, there is a certain degeneracy between bar strength and viewing angle (a strong bar seen at intermediate viewing angle may appear similar to a weaker bar seen exactly side-on; e.g. Bureau & Athanassoula 2003). Second, the ratio described evolves in N-body simulations, generally decreasing with time as the bar strengthens and lengthens but the central component changes little (see, again, Bureau & Athanassoula 2003). Third, gas inflow and subsequent star formation (absent from the pure N-body models) may substantially modify the light distribution in the central regions (e.g. Heller & Shlosman 1994; Friedli & Benz 1993). Fourth, the presence (or absence) of rings and lenses will complicate the comparison. We also note that as the profiles of Bureau & Athanassoula (2003) were integrated vertically over 0.5 original disc scalelengths, they are somewhat intermediate between our major-axis and summed profiles.

In fact, the variety observed in the ratios of the length of the thickest part of the bulge (or that of the central peak) to that of the intermediate region may well be dominated by the range of bar strengths present in the sample, rather than by the range of viewing angles. This is also suggested by the data of Lüticke et al. (2000), and it would similarly imply that the observed boxy-peanut ‘sequence’ is not mainly one of viewing angle, but is instead driven by a varying (i.e. increasing) bar strength. The small ratios observed in strong peanut-shaped bulges such as NGC 1886 and ESO 443-G042 can then be explained only if the central surface brightness peak and the thick part of the bulges are shorter in stronger bars (in a relative sense). This is natural if the inner part of the surface brightness profiles is dominated by a disc-like bulge, as will be argued in § 6.4. Indeed, Athanassoula (1992, 2002) showed that the size of (bar-driven) nuclear gas discs is limited by the outer inner Lindblad resonance (i.e. by the outermost $x_2$ orbit, elongated perpendicular to the bar), which itself inversely depends on the bar strength. Then, the stronger the bar, the smaller the nuclear disc.

6 DISCUSSION

6.1 Radial redistribution of material

Athanassoula (2002, 2003) argued that much of the bar-driven evolution in discs is due to a transfer of angular momentum from the inner (barred) disc to the outer disc (beyond corotation) and the dark halo, leading to a strengthening, lengthening and slowing down of the bar with time. To isolate the effects of angular momentum exchange, which leads to a radial redistribution of matter, we must preferentially consider the vertically-summed surface brightness profiles of Figures 1 and 2. While fewer galaxies with a B/PS bulge exhibit 3 or more regions in the summed profile than along the major-axis, a majority still do, while only one of the control sample galaxies does (see Table 2).

In scenarios where the formation and evolution of discs is dominated either by collapse or the gradual accretion of discrete components (e.g. cannibalized dwarf or satellite...
galaxies), there is no obvious mechanism to create the radial breaks in the disc density distribution observed in our sample galaxies, even less to explain the spatial correlation of those breaks with the ionized-gas and stellar kinematics (see Bureau & Freeman 1999; Chung & Bureau 2004). Similarly, the traditional interpretation of Freeman Type II profiles as inwardly truncated discs neither proposes an origin for the truncation nor explains why the truncation is only partial and the surface brightness profiles systematically have the specific shape observed. In bar-driven evolution scenarios, however, a break naturally occurs at the end of the bar (e.g. Athanassoula & Misiriotis 2002; Athanassoula 2002, 2003; Bureau & Athanassoula 2003). If this is indeed the case here, then the break in a surface brightness profile at the end of the intermediate region should systematically mark the bar's end.

Our major-axis surface brightness profiles match those of Athanassoula & Misiriotis (2002) and Bureau & Athanassoula (2003), and as we mentioned already the kinematics of the sample galaxies agree very well with barred galaxy models (see Bureau & Freeman 1999; Athanassoula & Bureau 1994; Chung & Bureau 2004; Bureau & Athanassoula 2005). A quick comparison further shows that the position where the rotation curve flattens, generally associated with the end of the bar (e.g. Bureau & Athanassoula 2003), always occurs very near (although generally slightly within) the break in the K-band surface brightness profile. Thus, our surface brightness profiles are consistent with those expected from bar-driven evolution scenarios, and the complex surface brightness profiles observed appear to be a direct consequence of the transfer of angular momentum and the radial (and vertical) rearrangement of material by the bar.

Most barred galaxies appear to have fast bars, whereby corotation, defined as the radius where the stars and bar pattern rotate at the same speed (and marking the transition between material loosing and gaining angular momentum), is located just beyond the end of the bar (see, e.g., Gerssen, Kuijken & Merrifield 2003; Aguerri, Debattista & Corsini 2003). Since the pattern speeds of the bars in our sample are unknown, a direct test of this is impossible. However, for the few galaxies with an apparent inner ring (see Table 2), the ring radius is systematically (roughly) equal to that of the break at the end of the intermediate surface brightness region. Given that inner rings are generally thought to occur near the inner 4:1 (ultra-harmonic) and corotation resonances (e.g. Schwarzschild 1981, 1982; Athanassoula et al. 1982; Buta 1993; Patsis, Skokos & Athanassoula 2003b), our galaxies are consistent with harbouring fast bars.

It is also interesting to note that the two latest galaxies in our sample, ESO 240-G011 and IC 5176 (with very small bulges and thus presumably very weak bars, if any) both have completely featureless (single) exponential outer discs. They thus offer useful benchmarks against which more complex systems can be compared.

6.2 Vertical redistribution of material

To isolate the effects of the vertical redistribution of material within bars, as predicted by models where the bar buckles and thickens (e.g. Combes & Sanders 1981; Combes et al. 1990; Raha et al. 1991), we compare here the major-axis and vertically-summed surface brightness profiles of our sample galaxies.

A cursory examination of our major-axis and summed surface brightness profiles shows that, for most galaxies with a B/PS bulge, they differ markedly from one another. As already pointed out, the vertically-summed profiles generally show fewer of the features described in §5 and they approach more closely a simple bulge plus exponential disc description, while the major-axis profiles generally show a steeper inner peak and a flatter intermediate region (the outer regions have a comparable slope in both profiles). Excellent examples of this are ESO151-G 004 and PGC 44931, but there are many.

If the scaleheight of the stars in a galaxy is constant with radius, one expects the major-axis and vertically-summed profiles to have the same functional form but different zero-points, i.e. be offset but parallel to one another (although not exactly, since the profiles are vertically-summed to a given surface brightness level rather than to a fixed number of scaleheights). This is not the case for most galaxies, however, even when the major-axis surface brightness profile is flat (when summing to a given surface brightness level is equivalent to summing to a fixed number of scaleheights), suggesting that the stellar scaleheights are not constant. See, again, ESO 151-G004 and PGC 44931, although any galaxy with a flat intermediate profile will do.

By discussing a single scaleheight, even in regions where the inner peak of the surface brightness profiles is non-negligible, we have however implicitly amalgamated the bulge and disc (no matter how the former is defined). In the classic bulge and disc model, where the bulge and disc represent structurally and kinematically distinct components, this will necessarily lead to a variation of the scaleheight, although one would naively expect this variation to be monotonic. The major-axis and vertically-summed surface brightness profiles of Figure 14 however, show that the variations are not monotonic. Furthermore, the functional difference between the two profiles is systematically greatest in the flat intermediate region, which would never be associated with a bulge (even in classic models) and is clearly dominated by disc material. Our major-axis and vertically-summed surface brightness profiles thus show that the radial scaleheight variations are real and that they occur in the discs (i.e. they are not due to bulge material), in direct contradiction to the long accepted wisdom that disc scaleheights are radially constant (see, e.g., van der Kruit & Searle 1981; 1982a; B. Shaw & Gilmore 1996; de Grijp & van der Kruit 1996). In fact, our results support a picture where the bulge and disc are not intrinsically distinct (either structurally or kinematically), but both instead emerge from the rapid radial variation of the scaleheight of the disc material (due to bar-related vertical resonances). This will be discussed in more depth in Paper II where it will also be shown that the variations are as expected from barred N-body models. Viewed simplistically, our observations argue for the reassignment of (most of) the bulge material to the disc (a B/PS bulge), but this is an important conceptual change.

To our knowledge, a variable stellar scaleheight has only been identified and discussed in a few galaxies. In the outer parts of NGC 3115, Capaccioli et al. (1983) assign the variation to the end of the disc self-gravity, while in the inner parts of the Milky Way Kent, Dame & Fazio (1991) assign
it to a variable ratio of young to old stars (their data in fact already show strong hints of the peanut shape of the Galactic bulge and a Freeman Type II profile). De Grijs & Peletier (1997) do note an increase of the scaleheight with (projected) radius in a large number of edge-on spiral galaxies, but they assign the increase to the presence of a thick disc with scale-length and scaleheight larger than those of the dominant (thin) disc component. Upon closer inspection, a number of other studies arguing for a constant scaleheight as a function of (projected) radius actually do detect variations, but those are generally monotonic and are argued to be consistent with the influence of a structurally distinct bulge (in the center) or a thickening of the outer disc (near the edge of the optical disc). The scaleheight variations are thus rejected as artefacts unrelated to the dominant disc, presumed thin (e.g. de Grijs & van der Kruit 1996). While this is defendable in some cases, it is doubtful in others, and the non-monotonic variations observed here and in Paper II in the inner and intermediate regions of our sample galaxies can not be explained away in such manners. Possible variations in the outer parts of our galaxies (e.g. flaring) are not of direct interest to the bulges’ structure.

6.3 Classic versus ‘pseudo’-bulges

As discussed also in Athanassoula (2005), galaxy bulges have traditionally been defined either as 1) the steep central component of the surface brightness profile (generally steeper and brighter than the inward extrapolation of the outer exponential disc; e.g. Carollo, Ferguson & Wyse 1999), 2) the thick galactic component (in isophotal terms, clearly sticking out of the equatorial plane in nearly edge-on objects) or 3) the kinematically hot component (the central peak in the velocity dispersion profile). Those three definitions have been used interchangeably and have generally been considered equivalent and consistent. But is this really the case?

Many diverse results show those views to be grossly oversimplified (see, e.g. Wyse, Gilmore & Frans 1995 for an early review). For example, there is now much evidence that some bulges are really concentrated (thin) discs, or equivalently that the central parts of discs can have a density profile steeper than that of the outer disc (see, e.g., Kormendy 1992 and Kormendy & Kennicutt 2004 for reviews). The central surface brightness profiles of most bulges also rarely approach the classic R^{1/4} law (e.g. Andrelikis et al. 1995; Balcells et al. 2005). Although our own surface brightness profiles are not ideal to tackle this last issue, because of the line-of-sight integration through the outer disc, many objects are nearly exponential in the center and we specifically study the minor-axis profiles in Paper II. Most importantly, our sample is ideal to test the equivalence of thick structures with steep light profiles and high velocity dispersions.

The images and surface brightness profiles of Figure 1 reveal that, in some galaxies, the thick central structure (defined isophotally) extends past the steep inner region of the surface brightness profile (both along the major-axis and vertically summed). Those two definitions of a bulge are thus neither equivalent nor consistent, at least in those cases. Our observations thus show that the steep inner region of the surface brightness profile is often contained within the thick central component. This was already hinted at by the poorer optical profiles of Chung & Bureau (2004), although Lutticke et al. (2001) found them to be generally equal. Examples include ESO 151-G004, NGC 2788A and ESO 443-G042, but there are many. In fact, our observations suggest that those two bulge definitions only seem to agree systematically in galaxies with a boxy-shaped (or round) bulge, such as NGC 1381 and IC 5096, while the disagreement is worst in some of the the strongest peanut-shape bulges. There is no straightforward explanation for this in classic bulge formation scenarios, but it is a natural consequence of bar-buckling/thickening mechanisms, where it is related to the bar viewing angle (seen respectively end-on and side-on).

Having argued that the steep inner region of the surface brightness profiles is generally contained within the thick part of the galaxies, it is also true that the intermediate and often flat region of the surface brightness profiles always extends beyond the thick component. Some of the best examples are found in the galaxies listed above, but again there are many (see, e.g., NGC 1886 and NGC 4469). This behaviour is hard to reconcile with classic bulge formation scenarios, but it arises naturally in bar-buckling/thickening ones. Indeed, while the (projected) surface brightness profile at intermediate radii remains shallow or flat within most of the barred region of a galaxy (i.e. over the entire deprojected bar length, no matter what the viewing angle is), only part of the bar is thick. So, while the fact that most thick components are shorter than the flat intermediate region of the surface brightness profile is simply a consequence that most bars are not seen exactly side-on, the fact that all thick components are shorter than the flat intermediate region reflects a fact much less widely appreciated, i.e. that only a fraction of the bar is actually thick. This point has not been stressed in the literature, but it is nicely illustrated and discussed in some depth in Athanassoula (2005). It is explained by the orbital structure of barred discs (e.g. Skokos et al. 2002a; Patsis et al. 2002) and is in fact visible in most published N-body models (e.g. Combes et al. 1990; Athanassoula & Misiriotis 2002; Athanassoula 2005).

6.4 Three-dimensional structure of B/PS bulges

A comparison of the major-axis and vertically-summed surface brightness profiles reveals another interesting fact, i.e. that the steep inner region of the surface brightness profiles is systematically more pronounced and dominant in the major-axis profiles than in the summed ones (see also Lutticke et al. 2001). This is a key indication of the 3D structure of galaxies with a B/PS bulge. It shows that most of the material found at high z belongs to the flat intermediate region of the surface brightness profiles rather than to the inner steep component, and that the latter has a smaller characteristic height than the former (this will be shown more transparently in Paper II). Thus, the steep inner region of the surface brightness profiles truly appears to be a thin concentrated disc, while the shallow intermediate region appears to be thick (and, according to the aforementioned arguments, barred). Figure 5 presents a cartoon version of this model, contrasted to the classic one. Athanassoula (2005) recently proposed a new nomenclature to differentiate different types of bulges based on their formation mechanism. In addition to classic bulges, which presumably form rapidly through either collapse or merging, she identifies B/PS bulges, which are simply (part of) thick
bars, and disc-like bulges, formed by (bar-driven) gaseous inflow and subsequent star formation. As more than one type of bulge can coexist in a single object, the latter two types are well suited to a description of the observations presented here.

As noted above, Kormendy (1993) and Kormendy & Kennicutt (2004) review much of the evidence for the interpretation of some bulges as concentrated discs, which they include in a broader category of ‘pseudo’-bulges, generally formed secularly (i.e. slowly). This vision of disc-like bulges does not explicitly include the extended bulges discussed in this paper, nor the many bulges which are isophotally thick and kinematically hot. The thick bar model advocated in Athanassoula (2003) and here addresses all three points simultaneously. In this respect, we know that the radial component of the velocity dispersion of the Milky Way bulge is similar to that of its inner disc ($\approx 100 \text{ km s}^{-1}$; e.g. Spaenhauer, Jones & Whitford 1992; Lewis & Freeman 1989), and that the velocity dispersion of the galactic bar/bulge is nearly isotropic. So although little is known about either, a difference between the vertical velocity dispersion of the inner disc and that of the bar/bulge must underlie their large structural differences. While it is unlikely that concentrated discs are hot vertically, (thick) bars appear to be (e.g. Gadotti & de Souza 2005; Debattista et al. 2005).
One could argue that our model comprises a number of distinct building blocks, like the classic models, but those blocks are very different and tightly intertwined dynamically. First, the large-scale bar leads to the formation of a concentrated disc, as in Kormendy’s (1992) picture, but this disc is thin, largely decoupled from the large-scale bar, and only addresses the first of the bulge definitions (steep surface brightness profile). Second, the large-scale bar itself is thick over most but not all its length, and it addresses the second definition of bulges (isophotally thick structure). As for the third definition of bulges (kinematically hot), it is unclear which of the inner disc or large-scale bar dominates the large velocity dispersions observed by Chung & Bureau (2004) in the equatorial plane, but the models suggest that the culprit is slightly elongated (Bureau & Athanassoula 2004). As mentioned above, face-on bars do appear to have large vertical velocity dispersions (e.g. Gadotti & de Souza 2003).

The 3D structure proposed for B/PS bulges in Athanassoula (2005) and here is thus diametrically opposed to the classic one. Instead of a thick and steep inner (classic bulge) component within a thin exponential disc, we propose a thin and steep inner disc component (a disc-like bulge) within a thick and shallow bar (a B/PS bulge), followed by an outer exponential disc (see Fig. 5). Furthermore, while the two components of classic bulges are generally thought to have formed separately through vastly different mechanisms and to interact little with each other, the three components of galaxies with B/PS bulges are largely made of the same material, which has evolved and been shaped by the single dominant influence of the bar, through the weak but relentless action of both radial and vertical resonances. We do not appear to need additional classic bulges to explain our observations, although it is difficult to rule them out completely at the very center of our galaxies (i.e. on a scale much smaller than the isophotally thick part).

We note that while this may seem to imply that no stellar population gradient should exist in barred galaxies, this is not true. The radial resonances, once established, will clearly affect the behaviour of the gaseous (dissipative) material, which in turns feeds star formation. It is thus expected that a certain segregation of the stellar populations (mainly due to influence of the radial resonances) will develop over time. While pure N-body models do not capture this richness, others models including gas dynamics and star formation do (e.g. Friedli, Benz & Kennicutt 1994, Heller & Shlosman 1994, Friedli & Benz 1997, but see also Nakasato & Nomoto 2003). However, the current simplicity of those models and the ignorance of much of the physics involved (primarily regarding star formation) prevent a detailed comparison with current observations, which are in any case limited (but see Martin & Roy 1992, Zaritsky, Kennicutt & Huchra 1994, Fisher, Franx & Illingworth 1996). There is no doubt, however, that the next step for bar-driven bulge formation models is to include realistic star formation and stellar population information (chemodynamics; see, e.g., Samland, Hensler & Theis 1997, Spurzem et al 2001, Michel-Dansac & Wozniak 2004).

The fact that the maximal radial extent of the inner disc appears to be systematically equal to or slightly smaller than the radius where the bar is thickest (see Fig. 4) may reflect some form of ‘self-regulating’ dynamics, i.e. that the inner disc is limited by a radial resonance (perhaps the radial inner Lindblad resonance, ILR) coinciding with a vertical resonance (perhaps the vertical ILR), as advocated e.g. by Combes et al. (1990). This should be explored further.

6.5 Control sample

Due to its limited size compared to the number of galaxies with a B/PS bulge in our sample, the control sample was always meant more for qualitative checks than for statistically robust comparisons. However, it is now clear that some galaxies do not belong in it. For example, Lütticke et al. (2000) classified the bulge of NGC 3957 as close to boxy-shaped, Bureau & Freeman (1999) and Chung & Bureau (2004) both found kinematic indications of weak bars in NGC 3957 and NGC 4703, and Chung & Bureau (2004) also found indication of non-axisymmetric motions in NGC 7123. The current K-band imaging also shows centered X features in NGC 3957 and NGC 7123 and secondary maxima in NGC 3957 and NGC 4703. All three galaxies show at least 3 regions in their surface brightness profile. The contrast between the B/PS bulge and control samples would thus be greater if NGC 3957 and NGC 4703 (and perhaps NGC 7123) were appropriately classified as B/PS.

Even if bars are peanut-shaped, a fraction of round bulges should still be expected to be bars, as the latter appear round when seen end-on. This fraction is however expected to be small, since already for an angle of 10–30° between the bar major-axis and the line-of-sight a boxy shape is observed (e.g. Combes et al. 1994, Lütticke et al. 2000, Athanassoula 2005). The large contamination of our control sample thus probably arises from the way it was selected. Indeed, no catalog of edge-on galaxies with a round bulge existed when it was constructed, so the control sample galaxies were selected from apparently misclassified objects in catalogs of galaxies with a B/PS bulge (Jarvis 1984, de Souza & dos Anjos 1987, Shaw 1987) and from the Flat Galaxy Catalog of Karachentsev et al. (1993). Carrying a study along the lines of the current one and those of Bureau & Freeman (1994) and Chung & Bureau (2004), but for an enlarged and properly selected sample of round bulges (for example using the more recent catalog of Lütticke et al. 2006), would thus be valuable. While this represents a significant amount of work, showing with more certainty that round bulges behave differently from B/PS bulges would greatly strengthen our conclusions.

7 SUMMARY AND CONCLUSIONS

We have presented K-band imaging observations of a sample of 30 edge-on spiral galaxies, most of which harbour a boxy or peanut-shaped (B/PS) bulge. Those data are minimally affected by dust and best trace population II stars, where most of the luminous mass resides. Our multi-faceted analysis suggests that B/PS bulges are simply the thick part of bars viewed edge-on (see Figures 1 and 2).

Galaxies with a B/PS bulge tend to have a more complex morphology than galaxies with other bulge types, more often showing centered or off-centered X structures, secondary maxima along the major-axis, and spiral-like structures. Best revealed by unsharp-masking, those features are also observed in three-dimensional N-body simulations of
barred discs (Athanassoula 2003), and can be explained by the orbital structure of bars (see, e.g., Patsis et al. 2002), although they need not be uniquely related to them. Only minor-axis extrema may be preferentially related to other bulge types. Whether taken along the major-axis or summed vertically (to simulate a flat galaxy), the surface brightness profiles of galaxies with a B/PS bulge are also more complex, more often showing 3 or more clearly separated regions, including a rather shallow or flat intermediate region (see Figure 4). Such Freeman Type II profiles are expected from barred galaxies (e.g. Bureau & Athanassoula 2005), but they do not have a natural self-consistent explanation in classic bulge formation scenarios.

The radial breaks observed in the vertically-summed profiles of our objects provide further evidence of the transfer of angular momentum and radial redistribution of disc material mediated by the (presumed) bars (see, e.g., Athanassoula 2002, 2003). Furthermore, the spatial correlations of the radial breaks with the ionized-gas and stellar kinematics (Bureau & Freeman 1999, Chung & Bureau 2004) are as expected for fast bars, currently favoured by observations. The differences between the major-axis and vertically-summed profiles provide evidence for abrupt variations of the scale-height of the disc material. This is, again, as expected from the diverse orbital families and vertical resonances and instabilities present in barred discs, but contrary to conventional wisdom. A quantitative and robust analysis of those scale-height variations and a comparison with N-body simulations will appear in future papers of this series.

Three other facts stand out. First, the steep inner region of the surface brightness profiles is systematically equal to or shorter than the isophotally thick part of the galaxies. Second, the isophotally thick part is itself systematically contained within the flat intermediate region of the surface brightness profiles. Third, the steep inner region of the surface brightness profiles is much more prominent along the major-axis than in the vertically-summed profiles.

We are thus led to radically alter the classic ‘bulge + disc’ model, composed of a thick and steep spheroidal bulge largely decoupled from a thin (possibly barred) exponential disc. Analogously to Athanassoula (2003), we propose here that galaxies with a B/PS bulge are composed of a thin concentrated disc (a disc-like bulge), formed secularly by the bar and responsible for the steep inner region of the surface brightness profiles, contained within a (partially) thick bar (the B/PS bulge), responsible for the flat intermediate region of the surface brightness profiles and the complex morphological structures, itself contained within a thin outer exponential disc (see Figure 3). Those components are closely intertwined dynamically and are largely made of the same (disc) material, shaped over long timescales by the bar.

The challenge to any competing formation scenario for galaxies with a B/PS bulge, which represent at least 45 per cent of the local galaxy population (Lüttdicke et al. 2000), is thus to simultaneously and self-consistently explain, equally well or better, their numerous morphological, photometric, and kinematic properties, as well as the correlations between them.

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