ENERGY FLOW BETWEEN JETS IN THE $k_T$ ALGORITHM.

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We consider the impact of the $k_t$ algorithm on energy flow into gaps between jets in any QCD hard process. While we confirm the observation that the $k_t$ clustering procedure considerably reduces the impact of non-global logarithms, we unearth yet new sources of logarithmic enhancement, that stem from using the $k_t$ algorithm to define the final state. We comment on the nature of the logarithms we find and discuss their all-orders treatment.

1. Introduction

The transverse energy ($E_t$) flow into gaps between hard jets is an observable that can offer important insights into different aspects of QCD. This includes information on the strong coupling, understanding of “all-order” behaviour as manifested in resummed predictions, non-perturbative power corrections and the underlying event at hadron colliders.

In order however to obtain information as accurately as possible from such an observable, one might expect that the minimum requirement is a solid (and correct) perturbative estimate, at least to the accuracy claimed. Failure to provide a correct estimate leads to attributing a potentially significant chunk of the model independent perturbative answer to model dependent pieces such as the underlying event or power corrections. How significant such a mis-attribution is, will naturally vary on a case-by-case basis but it suffices to say that the overall picture emerging from most studies of this kind, would be incomplete. This unfortunately is in fact the current situation, as described below.

*Work done in collaboration with Andrea Banfi.
2. Non global logs and the $k_t$ definition

We wish to consider the distribution in the $E_t$ flow into a gap $\Omega$, $1/\sigma d\sigma/dQ_\Omega$. We define the gap transverse energy as

$$Q_\Omega = \sum_i E_{t,i}$$

where the sum runs over either hadrons in $\Omega$ or as is more commonly the case, over soft jets in $\Omega$. These are obtained after a jet algorithm has been employed to cluster the final state into jets. This leaves aside from the high $E_t$ hard jets outside $\Omega$, soft jets that can populate the gap region.

The main problem in obtaining a resummed perturbative prediction for this observable is its non-global nature, that has itself been pointed out only relatively recently. Thus while for several observables such as many event shapes, one can obtain a next-to–leading log prediction by considering a veto on real emissions attached just to the hard emitting partons, this is not the case here even at leading logarithmic accuracy. For non global observables like the gap energy distribution the leading single-logarithmic resummed result can be expressed as (we consider first the definition involving a sum over hadrons in the gap)

$$\frac{1}{\sigma} \frac{d\sigma}{dQ_\Omega} \approx \frac{1}{\sigma} \frac{d}{dQ_\Omega} \left( e^{-R(Q/Q_\Omega)} S(Q/Q_\Omega) \right),$$

where the factor $e^{-R}$ represents uncanceled virtual emissions attached just to the hard jets, integrated over the gap region. The factor $S$ is the non-global term, where one has to consider a soft gluon emitted in $\Omega$ as being coherently emitted from an arbitrarily complex ensemble consisting not merely of the hard jets but additionally any number of soft gluons outside $\Omega$. Like the term $e^{-R}$, the factor $S$ also resums a class of single-logarithms, $\alpha_s^n \ln^n Q/Q_\Omega$. Till date the calculation of the non-global term $S$ has only been performed in the large $N_c$ approximation, making the non-global piece the dominant source of perturbative uncertainty at low $Q_\Omega$.

A partial solution to the problem was proposed by Appleby and Seymour, who pointed out that in several experimental studies one actually employs the definition based on summing over soft jets given by $k_t$ clustering. They assumed that the factor $e^{-R}$ was left intact by the clustering procedure, since it can be considered as exponentiating a result, $R$, obtained by considering just a single emission and its virtual counterpart. On the other hand the non-global piece involves multiple emission and has thus to be recomputed for the case of clustering. In particular gluons that fly...
into the gap can be pulled out of the gap by harder gluons outside which
reduces the non-global component significantly, but does not eliminate it
altogether. Thus the Appleby Seymour result assumed the form
\[
\frac{1}{\sigma} \frac{d\sigma}{dQ_\Omega} \approx \frac{1}{\sigma} \frac{d}{dQ_\Omega} \left( e^{-R(Q/Q_\Omega)} S_{kt}(Q/Q_\Omega) \right),
\]
where \( S_{kt} \) is the non-global contribution recomputed with \( k_t \) clustering.
This new non-global correction was found to be less than 20% of the
unclustered result.

3. Additional real-virtual mismatch induced by \( k_t \)
clustering

Now we reconsider the result (3) and show that it is incorrect in the sense
that it does not capture all the relevant single-logarithms even leaving aside
those suppressed by \( 1/N_c^2 \).

Let us concentrate on the factor \( e^{-R} \) where for the simple case of \( e^+ e^- \)
anihilation \( R \sim C_F \alpha_s(Q) \ln Q/Q_\Omega \). This term represents purely virtual
emissions above the scale \( Q_\Omega \), integrated in a phase space corresponding to
the gap region. Real emissions below the scale \( Q_\Omega \) have been assumed to
totally cancel while those above \( Q_\Omega \) are vetoed.

In fact real emissions attached to the hard jets (thus not pertaining
to the non-global term) do not completely cancel away, due to the use
of clustering. Consider two energy-ordered real emissions \( k_1 \) and \( k_2 \) for
which the probability of independent emission from a hard dipole \( ab \), can
be written as:

\[
\mathcal{P}_{\text{real, real}} = C_F^2 \alpha_s^2 W_{ab}(k_1) W_{ab}(k_2),
\]

where the \( W_{ab} \) are eikonal emission factors for gluons \( k_1 \) and \( k_2 \) from the
\( ab \) dipole. Likewise if the emitting (more energetic) gluon \( k_1 \) is virtual we have the one-real one–virtual independent emission probability.

\[
\mathcal{P}_{\text{real, virtual}} = -C_F^2 \alpha_s^2 W_{ab}(k_1) W_{ab}(k_2).
\]

For the pure virtual piece \( e^{-R} \) to be built up, one assumes these contribu-
tions to cancel, which is the case without clustering.

In the case one uses clustering on the final state, consider a situation
when the softer gluon \( k_2 \) is in \( \Omega \) and \( k_1 \) outside. If the distance between
the gluons \( (\Delta \eta_{12})^2 + (\Delta \phi_{12})^2 < R^2 \), where \( R \) is the jet radius, the gluon
\( k_2 \) is clustered out of the gap and hence in this region the double real
contribution to the energy distribution in \( \Omega \) is zero. However the mixed real-virtual term persists in this region, making a finite contribution to the distribution, since \( k_2^2 \) cannot be clustered away by a virtual gluon outside the gap. Thus instead of a cancellation we are left with a contribution that at order \( \alpha_s^2 \) has the colour factor \( C_F^2 \). This is clearly distinct from the non-global term at this order which has colour factor \( C_F C_A \) and is not accounted for by expanding \( e^{-R} \) either, confirming that it is a piece left out previously.

Specialising to the case of \( \Omega \) being a rapidity slice of width \( \Delta \eta \geq R \) we obtain the following additional single-logarithmic contribution to the integrated quantity \( \int_0^{Q_\Omega} \frac{1}{2 \pi} \frac{d\sigma}{dE_t} dE_t \):

\[
C_F^2 \text{primary} = \frac{16}{\pi} C_F^2 L^2 R^3,
\]

where \( L = \ln \frac{Q}{Q_\Omega} \) and “primary” refers to the fact that we have attached the gluons only to the primary hard partons produced in the process. We have confirmed the result above, valid for the simple case of \( e^+ e^- \rightarrow 2 \) jets with exact fixed-order computations.

4. All orders contribution

We have been able to numerically resum the terms we describe above, to all orders for the simple \( e^+ e^- \rightarrow 2 \) jets and DIS (1+1) jets cases. The effect we find is moderate over most of the phenomenological region of interest, changing the previous results by a maximum of 30%. However for the more complex cases of dijet photoproduction and hadron–hadron energy flow variables, further work is needed to estimate this effect at all orders. In these cases additional insight is also required to understand the potential role of superleading logarithms. A satisfactory understanding of the energy flow even to leading logarithmic accuracy, is thus some way off.

References