STUDY OF MULTIPOLAR RF KICKS FROM THE MAIN DEFLECTING MODE IN COMPACT CRAB CAVITIES FOR LHC.

Barranco García, R. Calaga, R. De Maria, M. Giovannozzi, A. Grudiev, R. Tomás, CERN, Geneva, Switzerland

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Presented at the International Particle Accelerator Conference (IPAC’12) –
May 20-25, 2012, N. Orleans, USA

Geneva, Switzerland, May 2012
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A crab cavity (CC) system is under design in the framework of the High Luminosity LHC project. Due to transverse space constraints on one hand and the RF frequency requirements on the other hand, the design of the crab cavities has to be compact. This results in the crab cavity shape being far from axially symmetric and, as a consequence, higher order multipolar components of the main deflecting mode are non-zero. In this paper, multipolar RF-kicks from the main deflecting mode are calculated in the compact CC for LHC. They are compared to the multipolar error in LHC magnets. The possible influence of the RF-kicks on the beam dynamics has been investigated by means of analytical estimates.

MULTIPOLAR RF KICK CALCULATION

Geometry of the LHC CC prototypes
In Fig 1, geometry of the three prototypes of the LHC CC are shown: (left) Ridged Waveguide cavity (RWCAV) [1]; (middle) Quarter Wave resonator cavity (QWCAV) [2]; (right) 4 rod resonator cavity (4RCAV) [3]. They exploit different ways to make a 400 MHz deflecting cavity so compact that it fits in the space available between two beam pipes of the LHC interaction regions [4]. Nevertheless, in all three cases, geometry of the cavity is axially non-symmetric giving rise to higher order multipolar components of the main deflecting mode.

Electro-magnetic field representation of the main deflecting mode
In an axially symmetric deflecting cavity, the main dipole mode (m=1) has only dipolar variation of the electromagnetic field ~exp(inϕ), where n=m=1. It is not the case, however, for the LHC CC prototypes shown in Fig. 1 due to strong deviation of the cavity shape from the axially symmetric one. In this case, all multipolar components n ≥ 1 are present depending on the cavity symmetry. In order to simplify the math we assume that all cavities are symmetric with respect to the horizontal plane XZ and are symmetric with respect to the Z-axis with the speed of light c, the perpendicular Lorentz force acting on it can be expressed as follows:

\[ F_{\perp} = e\left[E_{\text{kick}} + Z_0 \hat{u}_z \times H_{\text{kick}} \right] \]

where \( Z_0 \) is vacuum impedance, \( u_z \) the unit vector in Z direction, \( E_{\text{kick}} = E_\perp \cdot e^{i\omega t} \) and \( H_{\text{kick}} = H_\perp \cdot e^{i\omega t} \) are the electric and magnetic fields in the particle frame, respectively, while \( E_\perp \) and \( H_\perp \) are the corresponding electric and magnetic fields in the cavity frame. These fields can be expressed as a sum of multipoles similarly to what is done for magnets [5]. Then, the perpendicular Lorentz force is expressed as a sum over its multipolar components \( F^{(n)}_{\perp} \):

\[ F_{\perp}(r, \varphi, z) = \sum_{n=1}^{\infty} F^{(n)}_{\perp}(z) [\hat{u}_r \cos(n\varphi) + \hat{u}_\varphi \sin(n\varphi)] \]

Furthermore the transverse momentum change of the particle when traversing the cavity, i.e., the kick is given by

\[ \Delta p_{\perp}(r, \varphi) = \frac{1}{c} \int_0^L F_{\perp} dz = \sum_{n=1}^{\infty} \Delta p^{(n)}_{\perp}(r, \varphi) \]

where

\[ \Delta p^{(n)}_{\perp} = \frac{1}{c} e^{i\omega t} [\hat{u}_r \cos(n\varphi) + \hat{u}_\varphi \sin(n\varphi)] \int_0^L F_{\perp}^{(n)} dz \]

are the multipolar components of the RF-kick expressed in terms of the transverse EM fields integrated over the cavity length \( L \). They can also be expressed in terms of multipolar components of the accelerating field, namely,

\[ E_{\text{acc}} = E_\perp \cdot e^{i\omega t} \]

used Panofsky-Wenzel (PW) theorem:

\[ p_{\perp}(r, \varphi) = \frac{1}{c} \int_0^L d z [\nabla \times E_{\text{acc}}(r, \varphi, z)] \]

In this case:

\[ \Delta p^{(n)}_{\perp} = \frac{1}{c} \int_0^L m e^{i\omega t} \left[ \hat{u}_r \cos(n\varphi) + \hat{u}_\varphi \sin(n\varphi) \right] \int_0^L E_{\text{acc}}^{(n)}(z) dz \]

Comparison of \( \Delta p^{(n)}_{\perp} \) calculated from Eqs. (1) and (2) gives an estimate of numerical accuracy. Finally for ultra-relativistic particle the transverse RF-kick strength can be equated to a magnetic kick of the same strength:

\[ b_n = \int_0^L B^{(n)} dz \text{ where } B^{(n)} = \frac{1}{e c} F_{\perp}^{(n)} = \frac{n j}{e c} E_{\text{acc}}^{(n)} \]

This expression is useful for comparing the strength of RF and magnetic multipoles in the same units. Although RF and magnetic multipoles can be compared, there is an essential difference. Magnetic multipoles are static and the corresponding kick does not depend on the longitudinal position of particles, whereas RF multipoles are harmonically oscillating at \( f_{rf} \), i.e., at 400 MHz in the LHC case, and the corresponding kick depends on the RF frequency requirements on the other hand, the design of the crab cavities has to be compact. This results in the crab cavity shape being far from axially symmetric and, as a consequence, higher order multipolar components of the main deflecting mode are non-zero. In this paper, multipolar RF-kicks from the main deflecting mode are calculated in the compact CC for LHC. They are compared to the multipolar error in LHC magnets. The possible influence of the RF-kicks on the beam dynamics has been investigated by means of analytical estimates.

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\[ F_{\perp} = e\left[E_{\text{kick}} + Z_0 \hat{u}_z \times H_{\text{kick}} \right] \]

*The research leading to these results has received funding from the European Commission under the FP7 project HiLumi LHC, GA no. 284404, co-funded by the DoE, USA and KEK, Japan.
phase of particles. Consequently, RF multipoles $b_n$ are complex numbers, with the real part being the kick to a particle at the deflecting RF phase and the imaginary part being the kick to a particle at crabbing RF phase. The RF-kick to a particle at arbitrary RF phase difference from the crabbing RF phase $\delta \phi$ is given as:

$$\Delta Q_z^{(n)} = \Im\left\{ e^{i(n-1)\theta} \sin(n\phi) \right\} b_{n} e^{i\delta \phi}.$$ 

**Multipolar RF kicks in LHC CC prototypes**

EM field maps have been calculated for the three cavity geometries presented in Fig. 1 using frequency domain finite element code HFSS [6]. Then the integrated strength of the RF multipoles has been calculated in two different ways using Eqs. (1) and (2). Comparing these two values, the relative accuracy of the calculations is estimated to decrease from few per mille for $n=3$ down to few per cent for $n=7$. Finally, values of the multipoles calculated at the nominal deflecting voltage $V_{CC} = 10$ MV using Eq (1) are summarized in Table 1. There are a few things worth mentioning: due to symmetry all even multipoles are zero in RWCA V and 4RCA V; in QWCA V, which has no symmetry with respect to the plane YZ, all multipoles are present. Last but not least, the imaginary part of all calculated multipoles in all three cavities is zero within the accuracy of the calculation. In other words, higher order multipoles ($n>1$) are in phase with the main dipole component ($n=1$) and there is no transverse kick acting on a particle at the crabbing phase.

In order to give a scale to the RF multipolar strength for LHC CC prototypes, they are compared to the rms field quality of the D2 separation dipoles that are currently installed in the LHC ring in IR1 and IR2 [7]. It is interesting to note that both the magnetic and the RF multipoles are of the same order of magnitude. This means that a careful analysis of the impact of the CC field quality is mandatory, both in terms of analytical indicators of the beam dynamics (see next section) and in terms of numerical simulations of the long-term impact on the beam dynamics.

**BEAM DYNAMICS**

The presence of RF multipolar kicks will introduce different type of beam dynamics perturbations according to their order. The aberrations considered include up to the octupolar component:

- Linear tune shift
  $$\Delta Q_{x,y} = \frac{1}{4\pi} \frac{\beta_{x,y} B_{\rho}}{D_y} b_{2},$$

- Chromaticity shift
  $$\Delta \varepsilon_{x,y} = \frac{\pm 1}{4\pi} \frac{D_y B_{\rho}}{D_y} \frac{b_{2}}{b_{2}} \rho$$

- Coupling
  $$\Delta Q_{\min} = \frac{1}{2\pi} \sqrt{\beta_{x,y} B_{\rho} \frac{b_{2}}{b_{2}}}$$

- Amplitude detuning, generating a tune shift at $3\sigma$ (the coupling term is not considered here)
  $$\Delta Q_{x,y} = \frac{3}{8\pi} \beta_{x,y}^{2} \frac{b_{3}}{b_{2}} J_{x,y}$$

where $\beta_{x,y}, D_y$ are the beta and vertical dispersion functions at the location of the CC and $J_{x,y}, \delta, B_{\rho}$ are the action, relative momentum error and beam rigidity of the
The aberrations computed for the multipolar coefficients listed in Table 1 using the optical conditions of Table 2, where the maximum optical parameters between the left and right side of the insertions are used, are listed in Table 3. It is worth recalling that the CC installed in IR1 (vertical crossing) will feature also an $a_i$, generated by the rotation of the original $b_i$ component by 90 degrees, whose effect is not estimated here. The linear tune shift for the QWCAV type is found to be large. This is going to generate a tune modulation, whose amplitude is larger than in existing machines like SPS, HERA. The operational experience suggested that any tune modulation should be kept smaller than $\Delta Q_{\text{max}} \sim 10^{-4}$ [9]. Parenthetically, this criterion has been used to specify the performance of the LHC power converters.

### Table 3: Evaluation of the possible beam aberrations induced by the CC installed in one side of one IR in which we assume a horizontal crossing plane.

<table>
<thead>
<tr>
<th>RWCAV</th>
<th>QWCAV</th>
<th>4RCAV</th>
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<tbody>
<tr>
<td>$</td>
<td>\Delta Q_{x,y}</td>
<td>\times 10^3$</td>
</tr>
<tr>
<td>$</td>
<td>\Delta \xi_{x,y}</td>
<td>\times 10^3$</td>
</tr>
<tr>
<td>$</td>
<td>\Delta \phi</td>
<td>\times 10^6$</td>
</tr>
</tbody>
</table>

The coupling is zero by design for this particular optical configuration as the vertical dispersion at the location of the CC is zero. The values of the higher-order aberrations for the crab cavities are comparable with those of the D2 separation dipole.

**CONCLUSIONS**

A formalism characterizing the non-linearity of deflecting RF field has been applied to the LHC compact crab cavity prototypes in a similar way as it is done for the multipolar components of the EM field have been calculated for the LHC CC prototypes and found comparable to the field quality of the D2 separation dipole of the nominal LHC machine. First estimates of the effect of CC field quality on beam dynamics have been carried out. Assuming static perturbations, the linear tune shift, the chromaticity shift and the amplitude detuning have been estimated for the three CC layouts. The tune modulation from the QWCAV could be a source of concern. It is worth stressing that such analytical estimates are only indicators of what could be the situation. In reality, the long-term effect on the beam dynamics should be checked with detailed numerical simulations. Recent developments in the MAD-X [10] and the SixTrack codes have made it possible to perform such simulations in the near future.

**ACKNOWLEDGMENT**

We would like to thank L. Deniau, A. Latina, and F. Schmidt for their support in implementing the changes in MAD-X and SixTrack required for long-term CC simulations. Useful discussions and suggestions by S. Fartoukh are warmly acknowledged.

**REFERENCES**